

Hot Quarks and Gluons at an Electron-Ion Collider

Matthew A. C. Lamont
Brookhaven National Lab

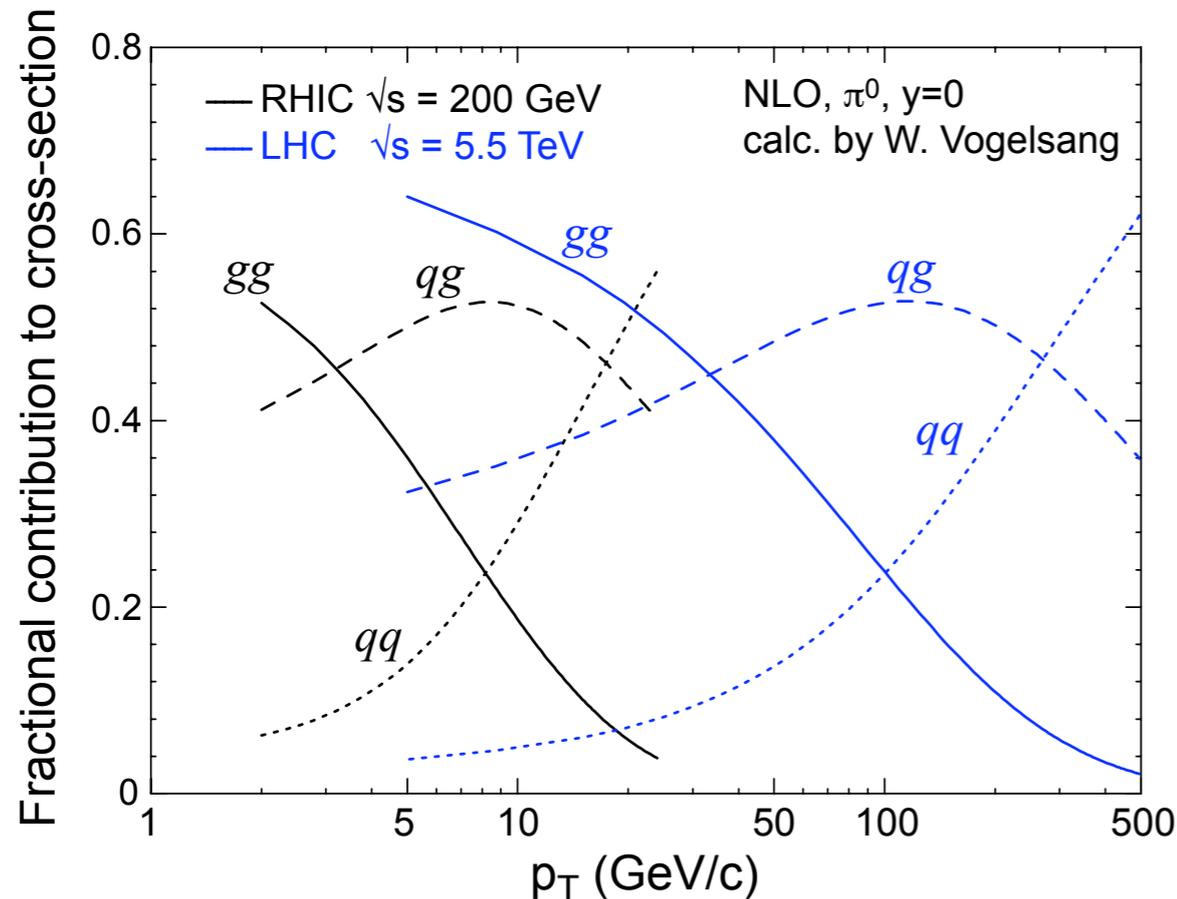
Talk Outline

- The role of glue in HI collisions
- How to measure the gluon distributions
- eA vs ep and the “Nuclear Oomph” factor
- The EIC machine and detector concepts
- Where we are and where we’re going

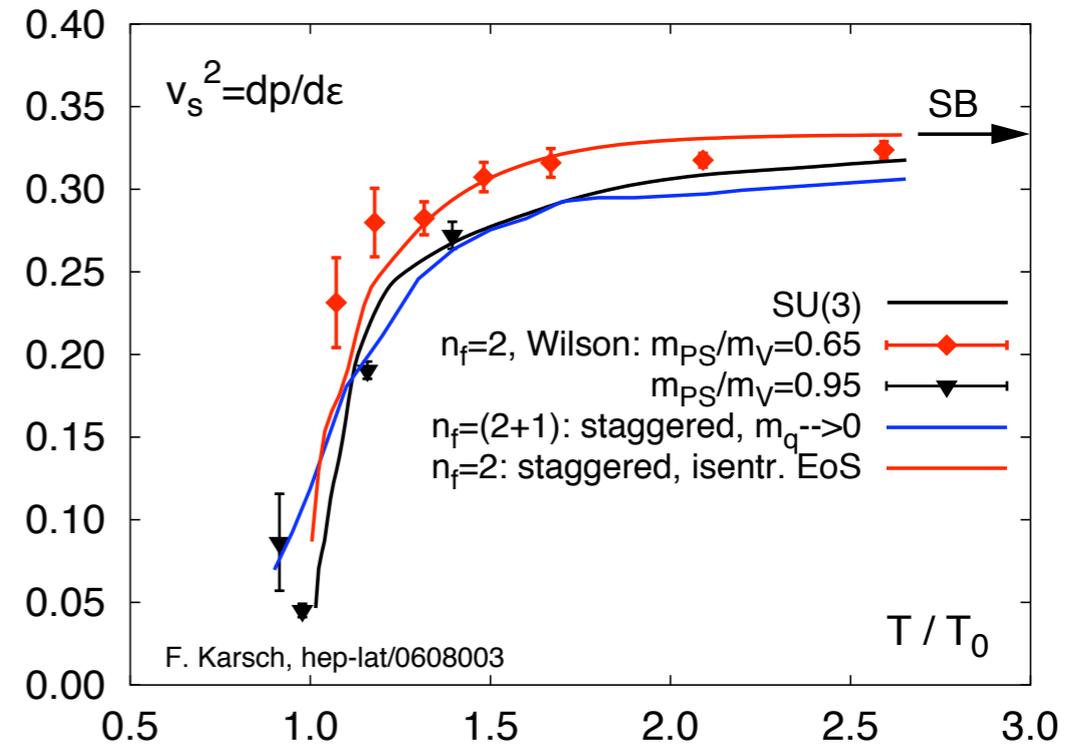
EIC on the web: <http://web.mit.edu/eicc>
e+A working group: <http://www.eic.bnl.gov>

The role of Glue in Heavy-Ion collisions

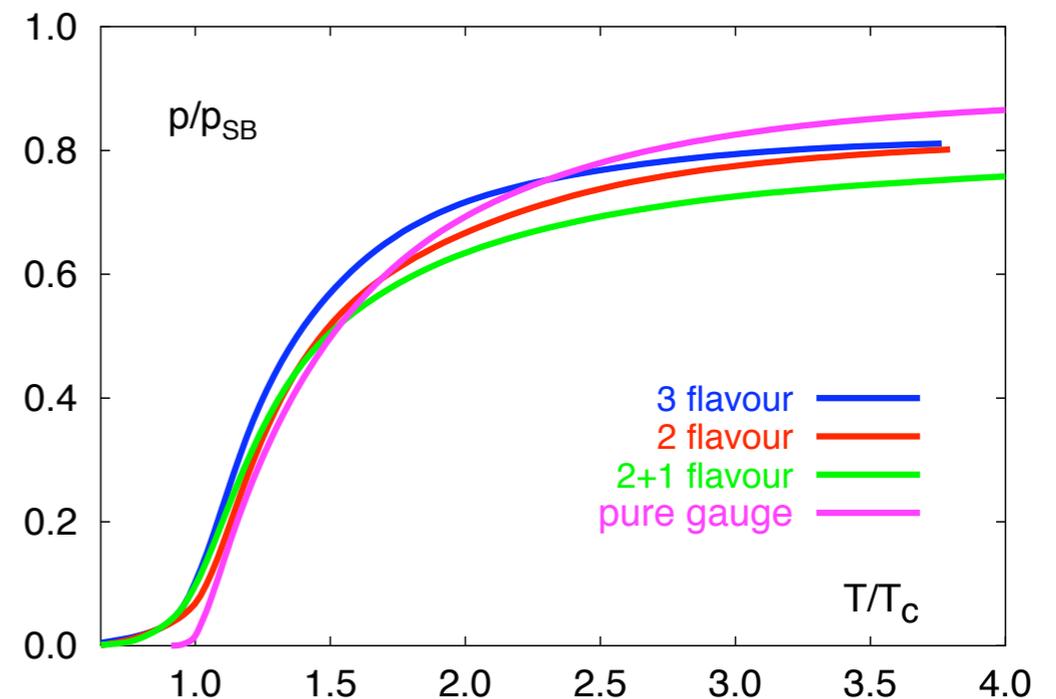
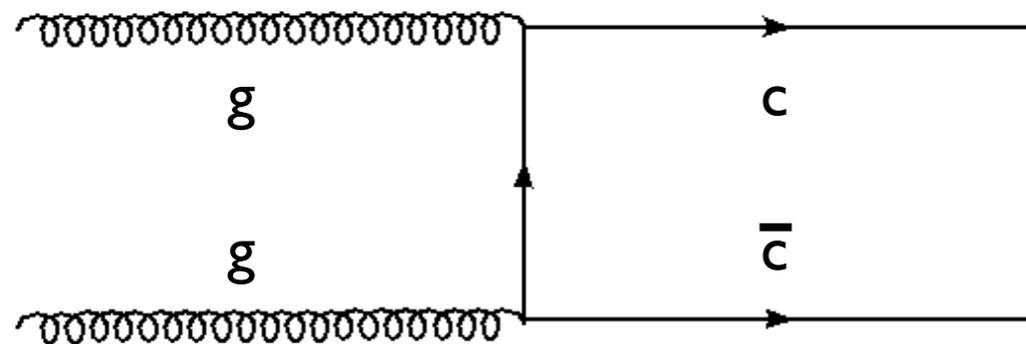
Jets (π^0 production)



Lattice Gauge Theory:



Heavy Flavour Production



The role of Glue in Heavy-Ion collisions

Jets (π^0 production)

Lattice Gauge Theory:

To move towards
understanding HI physics
quantitatively, we need to
understand the role of glue
in HI Collisions !!

What do we know about gluons?

What do we know about gluons?

Glue and the QCD Lagrangian:

$$L_{QCD} = \bar{q}(i\gamma^\mu \partial_\mu - m)q - g(\bar{q}\gamma^\mu T_a q)A_\mu^a - \frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu}$$

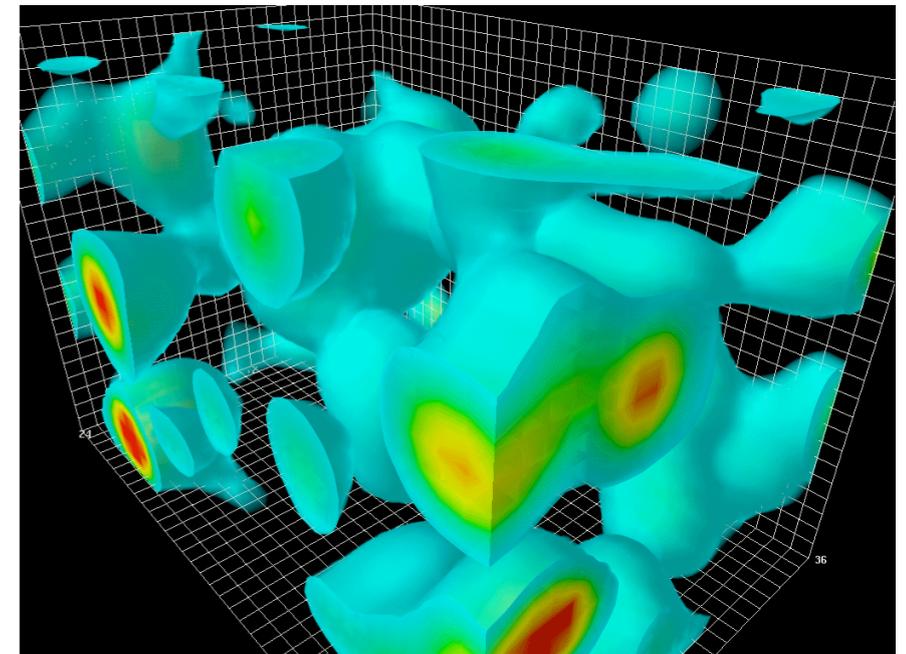
- **>98% of all visible mass due to “emergent” phenomena not evident from Lagrangian**
 - χ SB & Colour Confinement

What do we know about gluons?

Glue and the QCD Lagrangian:

$$L_{QCD} = \bar{q}(i\gamma^\mu \partial_\mu - m)q - g(\bar{q}\gamma^\mu T_a q)A_\mu^a - \frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu}$$

- >98% of all visible mass due to “emergent” phenomena not evident from Lagrangian
 - χ SB & Colour Confinement



Action (\sim energy) density fluctuations of gluon-fields in QCD vacuum (2.4 x 2.4 x 3.6 fm) (Derek Leinweber)

What do we know about gluons?

Glue and the QCD Lagrangian:

$$L_{QCD} = \bar{q}(i\gamma^\mu \partial_\mu - m)q - g(\bar{q}\gamma^\mu T_a q)A_\mu^a - \frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu}$$

- >98% of all visible mass due to “emergent” phenomena not evident from Lagrangian

- χ SB & Colour Confinement

- **Gluons**

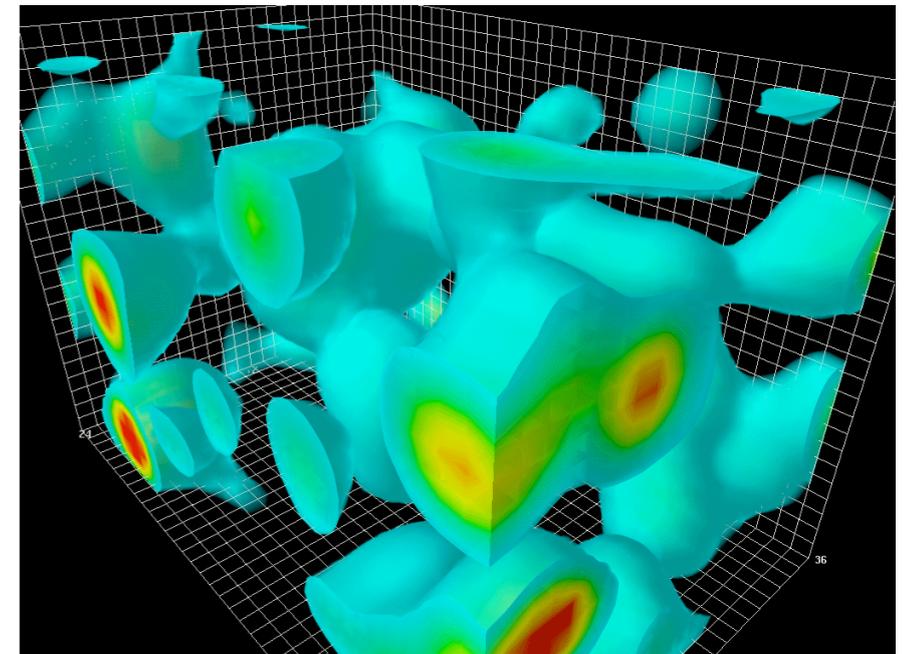
- ➔ Mediators of the strong interaction

- ➔ Determine essential features of QCD

- ▶ Asymptotic freedom from gluon loops

- ➔ Dominate structure of QCD vacuum (χ SB)

- ➔ Quenched L_{QCD} gets hadron masses correct to $\sim 10\%$

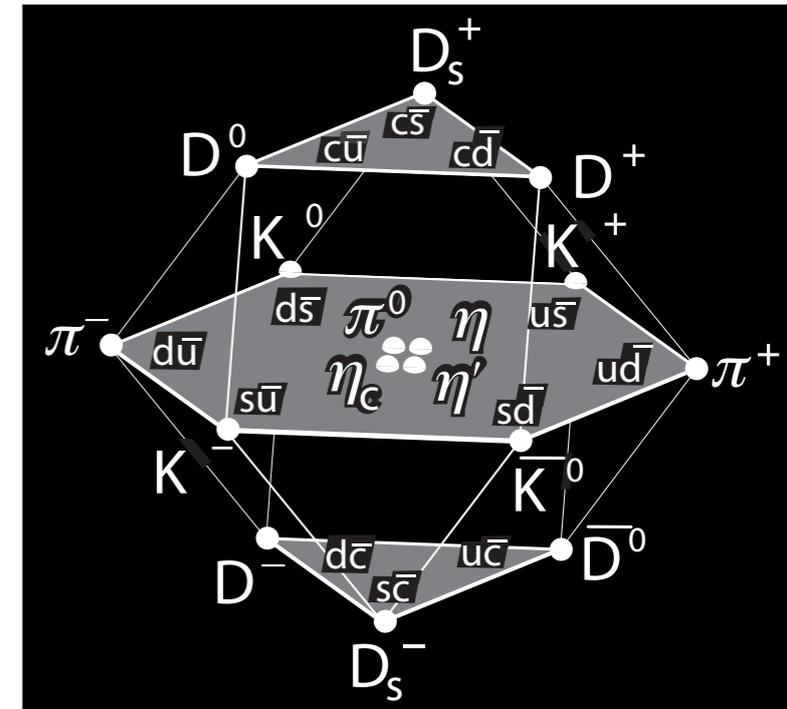


Action (\sim energy) density fluctuations of gluon-fields in QCD vacuum (2.4 x 2.4 x 3.6 fm) (Derek Leinweber)

Glue and the Lagrangian

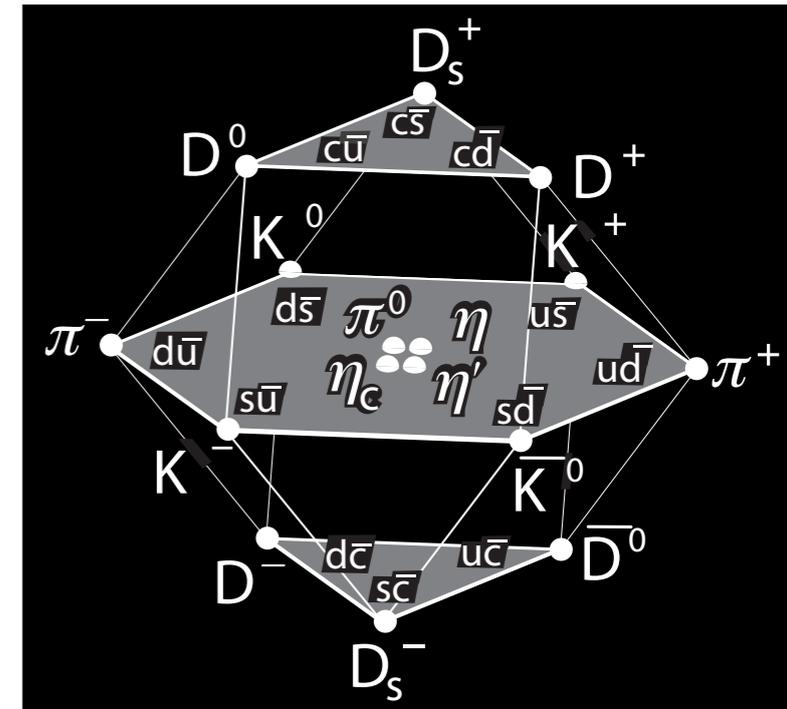
Glue and the Lagrangian

- **Hard to “see” glue in the low-energy world**
 - ➔ Gluon degrees of freedom “missing” in hadronic spectrum
 - ➔ Constituent Quark Picture?
- From DIS:
 - ➔ Drive the structure of baryonic matter already at medium-x
- Crucial players at RHIC and LHC
 - ➔ Drive the entropy
 - ➔ High energy cross-section suggests Pomeron (2 gluon) exchange important



Glue and the Lagrangian

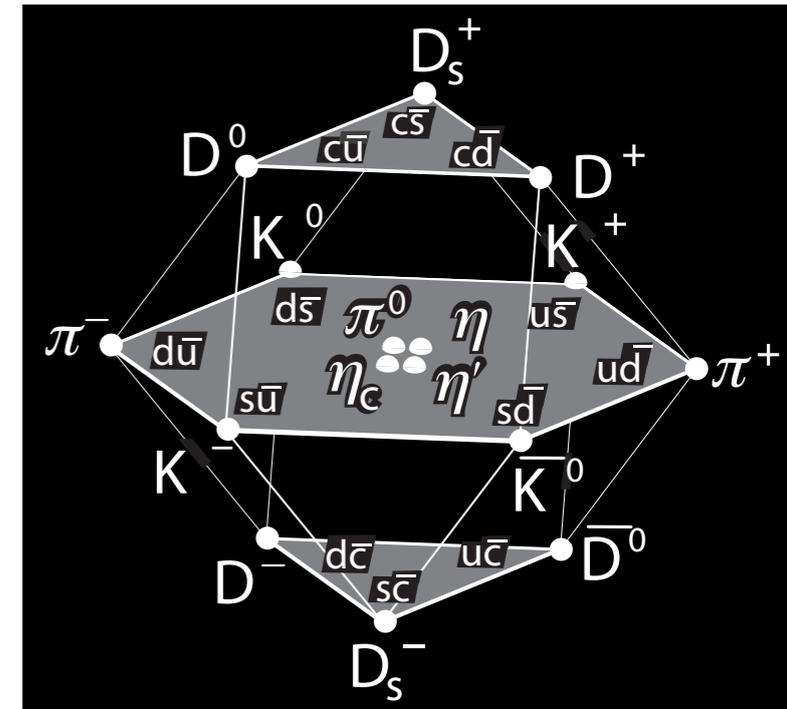
- **Hard to “see” glue in the low-energy world**
 - ➔ Gluon degrees of freedom “missing” in hadronic spectrum
 - ➔ Constituent Quark Picture?
- From DIS:
 - ➔ Drive the structure of baryonic matter already at medium-x
- Crucial players at RHIC and LHC
 - ➔ Drive the entropy
 - ➔ High energy cross-section suggests Pomeron (2 gluon) exchange important



- What is the **spatial** and **momentum** distribution of gluons in nuclei/nucleons?
- What are the **properties** of **high-density gluon matter**?
- How do **quarks** and **gluons** interact as they traverse matter?
- What role do the **gluons** play in the **spin structure** of the nucleon?

Glue and the Lagrangian

- **Hard to “see” glue in the low-energy world**
 - ➔ Gluon degrees of freedom “missing” in hadronic spectrum
 - ➔ Constituent Quark Picture?
- From DIS:
 - ➔ Drive the structure of baryonic matter already at medium-x
- Crucial players at RHIC and LHC
 - ➔ Drive the entropy
 - ➔ High energy cross-section suggests Pomeron (2 gluon) exchange important



- What is the **spatial** and **momentum** distribution of gluons in nuclei/nucleons?
- What are the **properties** of **high-density gluon matter**?
- How do **quarks** and **gluons** interact as they traverse matter?
- What role do the **gluons** play in the **spin structure** of the nucleon?

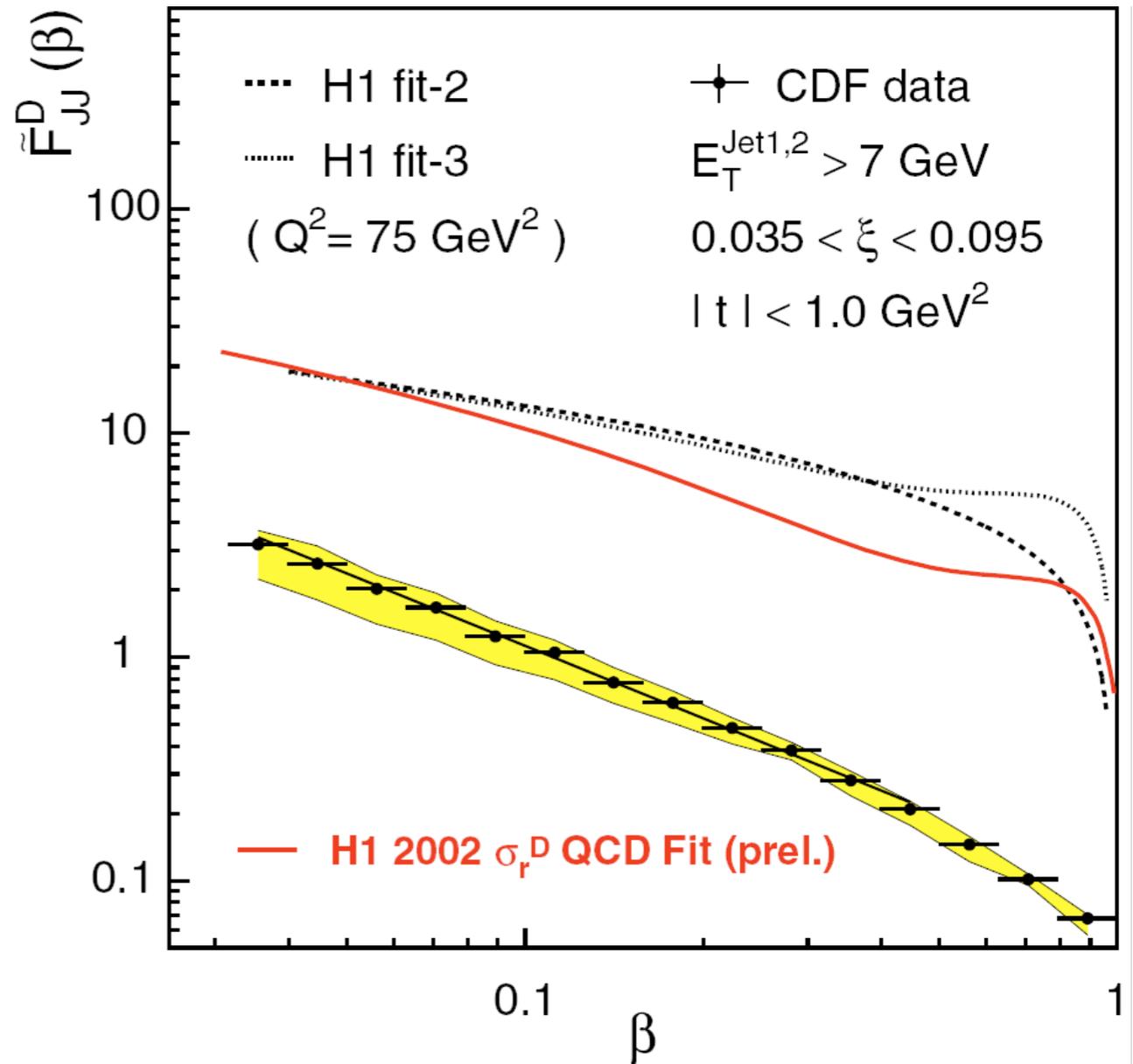
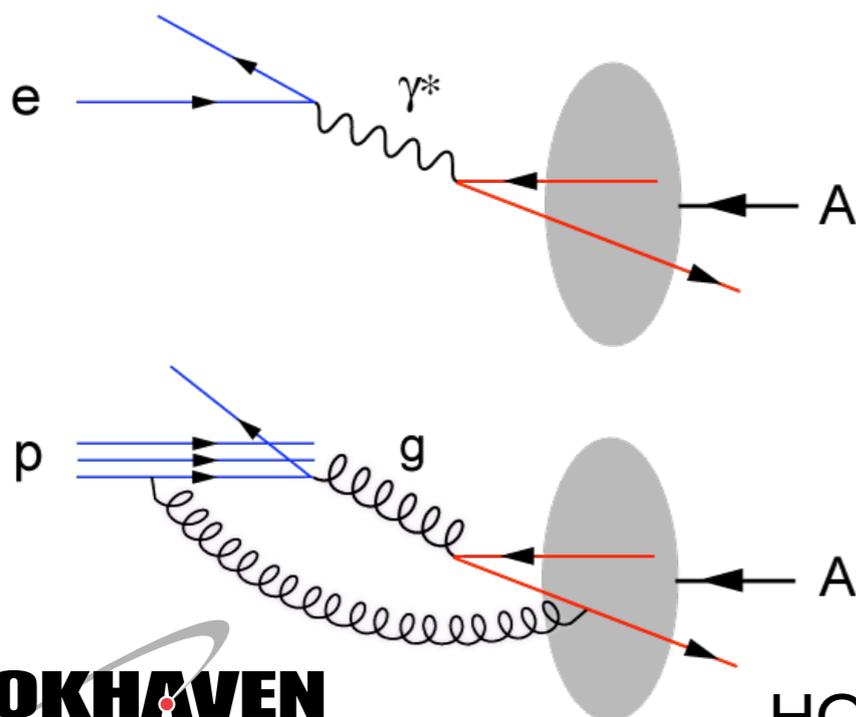
How do we get to the answers?

Accessing the Glue - $p+A$ vs $e+A$

F. Schilling, hex-ex/0209001

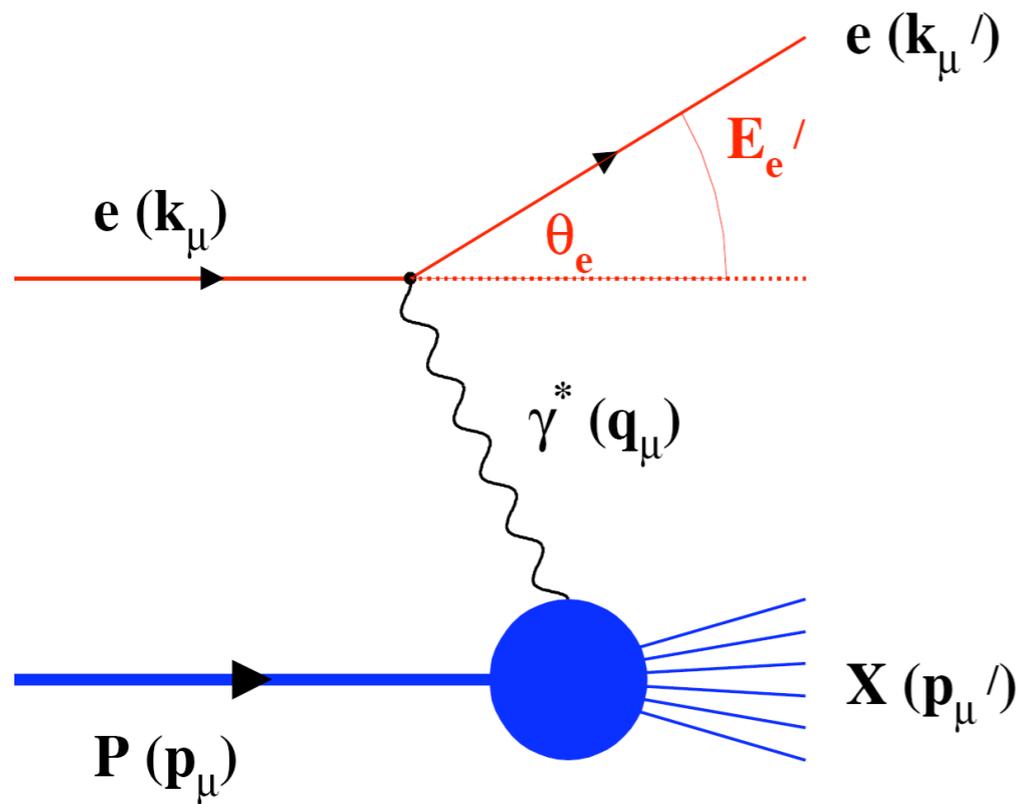
- Both $e+A$ and $p+A$ provide excellent information on properties of gluons in the nuclear wave functions
- Both are complementary and offer the opportunity to perform stringent checks of factorization/universality \Rightarrow
- But:

\rightarrow soft colour interactions between p and A before and after the primary interaction

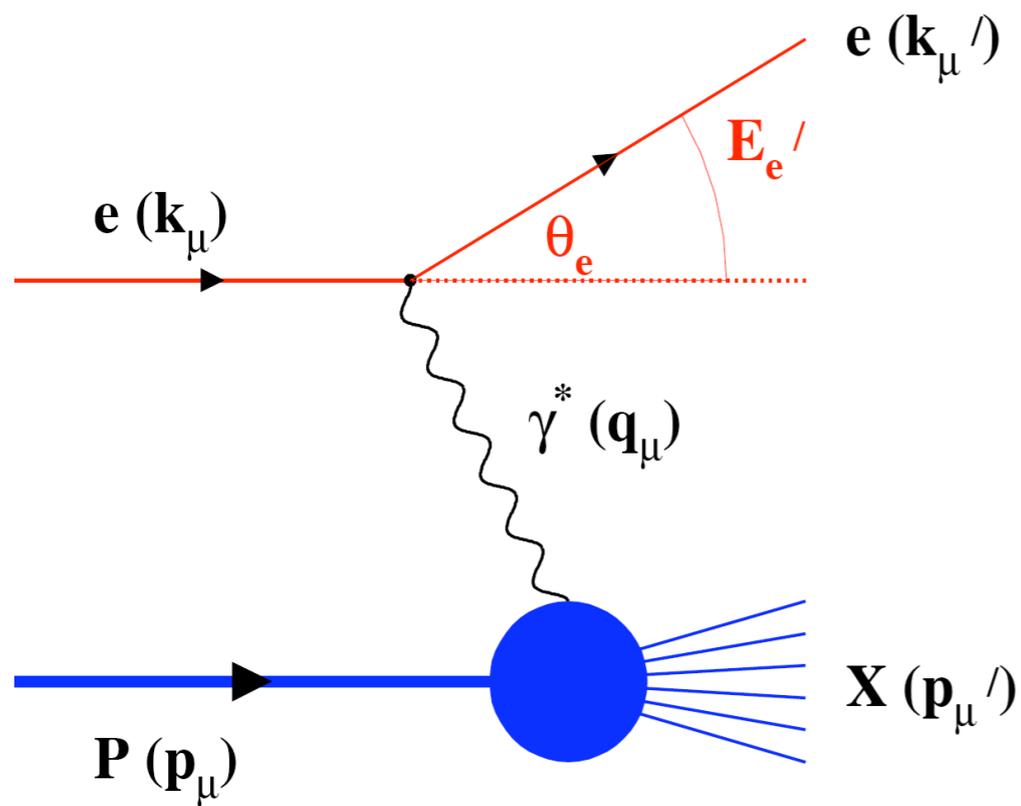


Breakdown of factorization ($e+p$ HERA versus $p+p$ Tevatron) seen for diffractive final states.

DIS Kinematics



DIS Kinematics

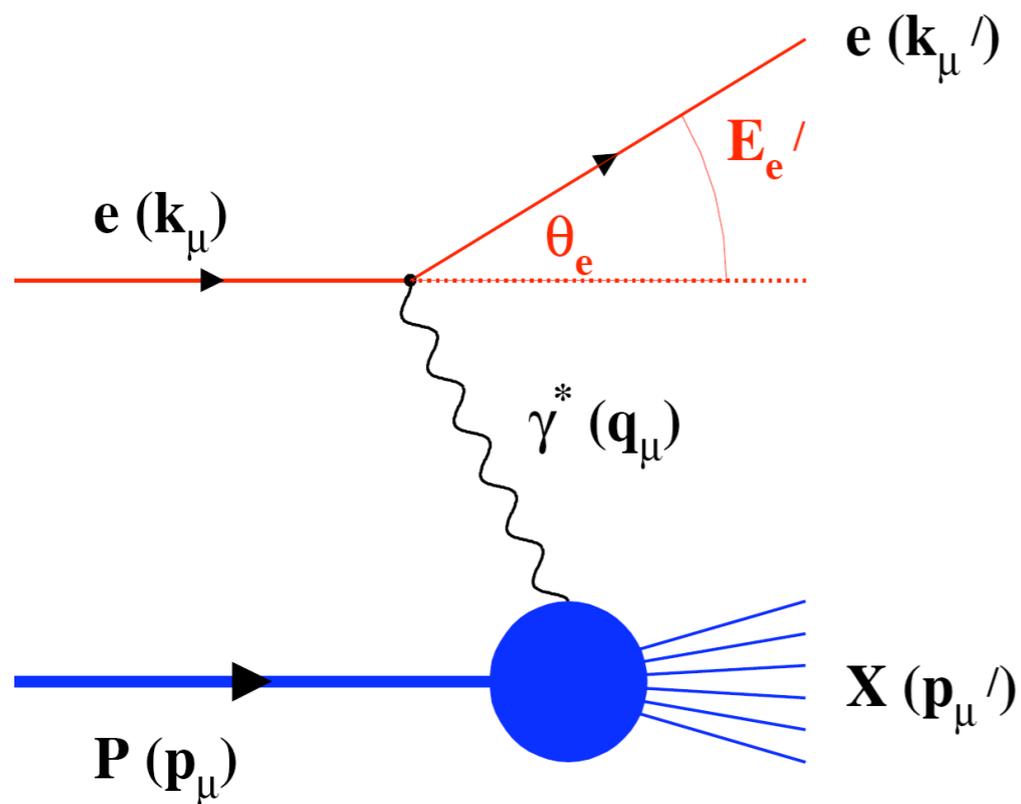


$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

$$Q^2 = 4E_e E'_e \sin^2\left(\frac{\theta'_e}{2}\right)$$

Measure of resolution power or "Virtuality"

DIS Kinematics



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

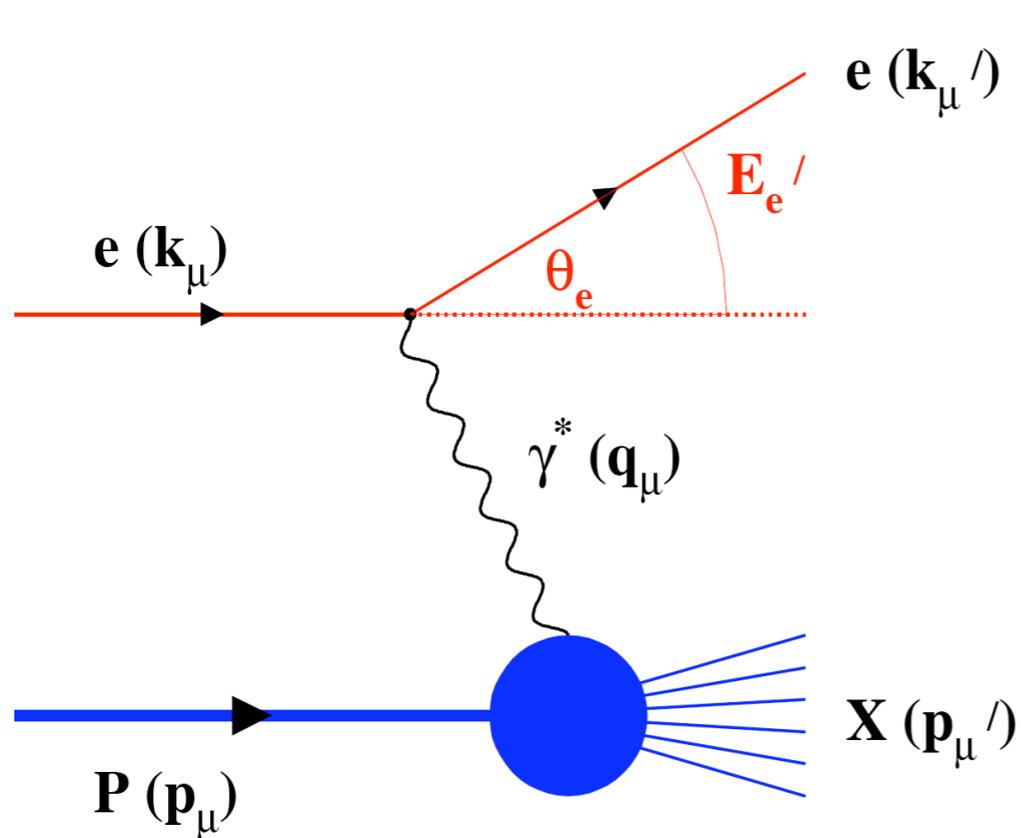
Measure of resolution power or "Virtuality"

$$Q^2 = 4E_e E_{e'} \sin^2\left(\frac{\theta'_e}{2}\right)$$

$$y = \frac{pq}{pk} = 1 - \frac{E_{e'}}{E_e} \cos^2\left(\frac{\theta'_e}{2}\right)$$

Measure of inelasticity

DIS Kinematics



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

Measure of resolution power or "Virtuality"

$$Q^2 = 4E_e E_{e'} \sin^2\left(\frac{\theta'_e}{2}\right)$$

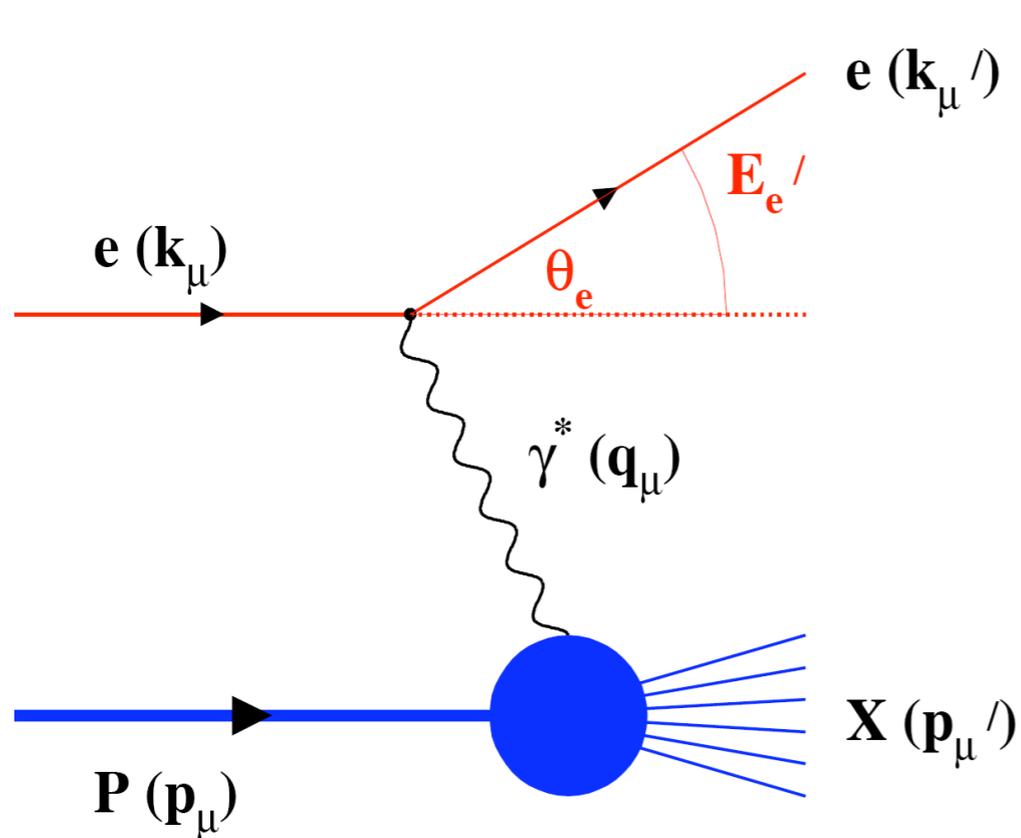
$$y = \frac{pq}{pk} = 1 - \frac{E_{e'}}{E_e} \cos^2\left(\frac{\theta'_e}{2}\right)$$

Measure of inelasticity

$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of momentum fraction of struck quark

DIS Kinematics



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

Measure of resolution power or "Virtuality"

$$Q^2 = 4E_e E_{e'} \sin^2\left(\frac{\theta'_e}{2}\right)$$

$$y = \frac{pq}{pk} = 1 - \frac{E_{e'}}{E_e} \cos^2\left(\frac{\theta'_e}{2}\right)$$

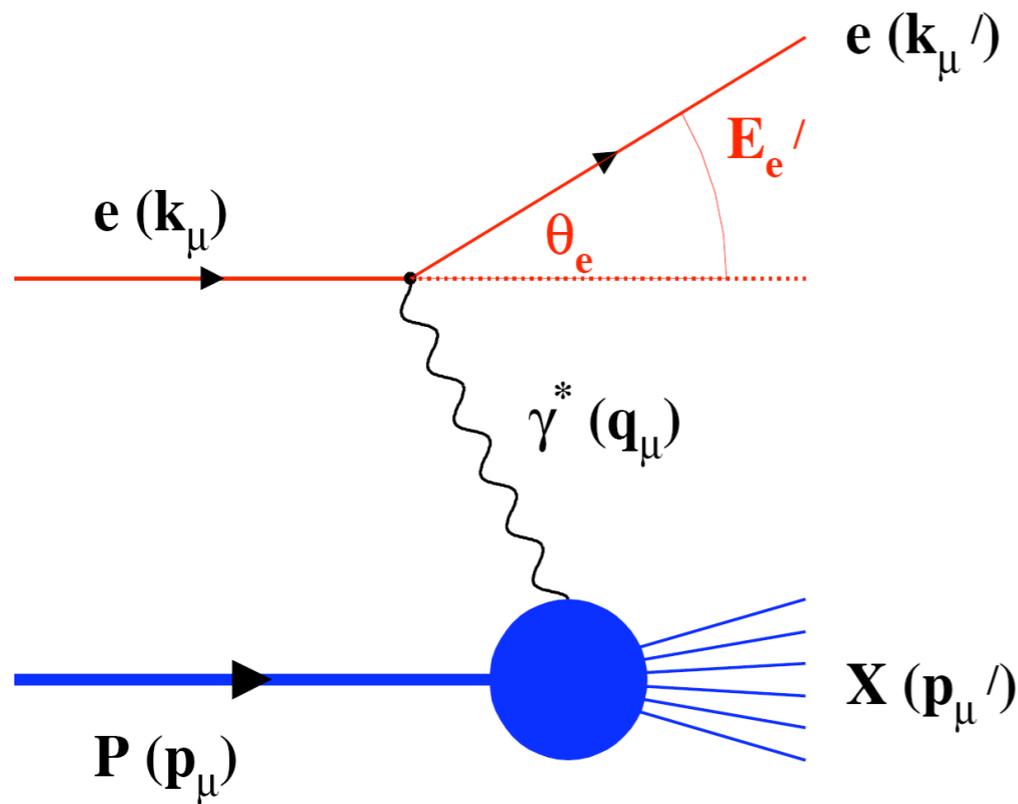
Measure of inelasticity

$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of momentum fraction of struck quark

$$\frac{d^2\sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

DIS Kinematics



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

Measure of resolution power or "Virtuality"

$$Q^2 = 4E_e E_{e'} \sin^2\left(\frac{\theta_e}{2}\right)$$

$$y = \frac{pq}{pk} = 1 - \frac{E_{e'}}{E_e} \cos^2\left(\frac{\theta_e}{2}\right)$$

Measure of inelasticity

$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of momentum fraction of struck quark

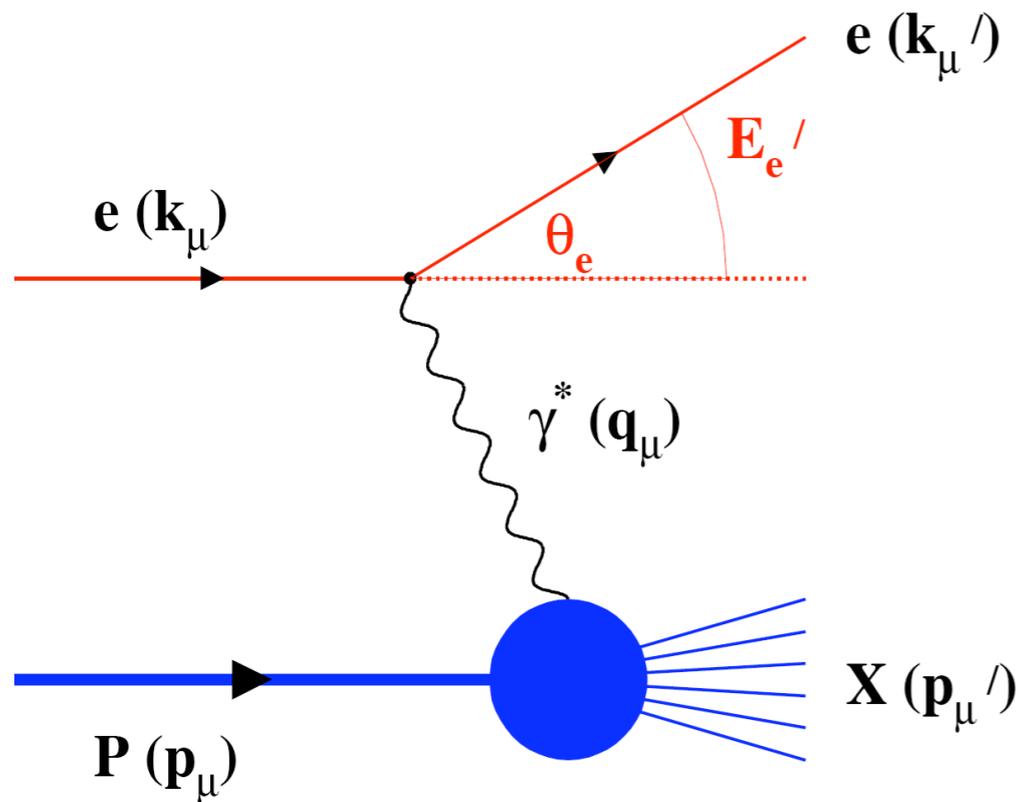
$$\frac{d^2\sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

- *Structure functions:*

→ $F_2(x, Q^2) \Rightarrow$ q and \bar{q} momentum distributions

→ $F_L(x, Q^2) \Rightarrow$ gluon momentum distribution

DIS Kinematics



$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

Measure of resolution power or "Virtuality"

$$Q^2 = 4E_e E_{e'} \sin^2\left(\frac{\theta'_e}{2}\right)$$

$$y = \frac{pq}{pk} = 1 - \frac{E_{e'}}{E_e} \cos^2\left(\frac{\theta'_e}{2}\right)$$

Measure of inelasticity

$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

Measure of momentum fraction of struck quark

$$\frac{d^2\sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

- *Structure functions:*

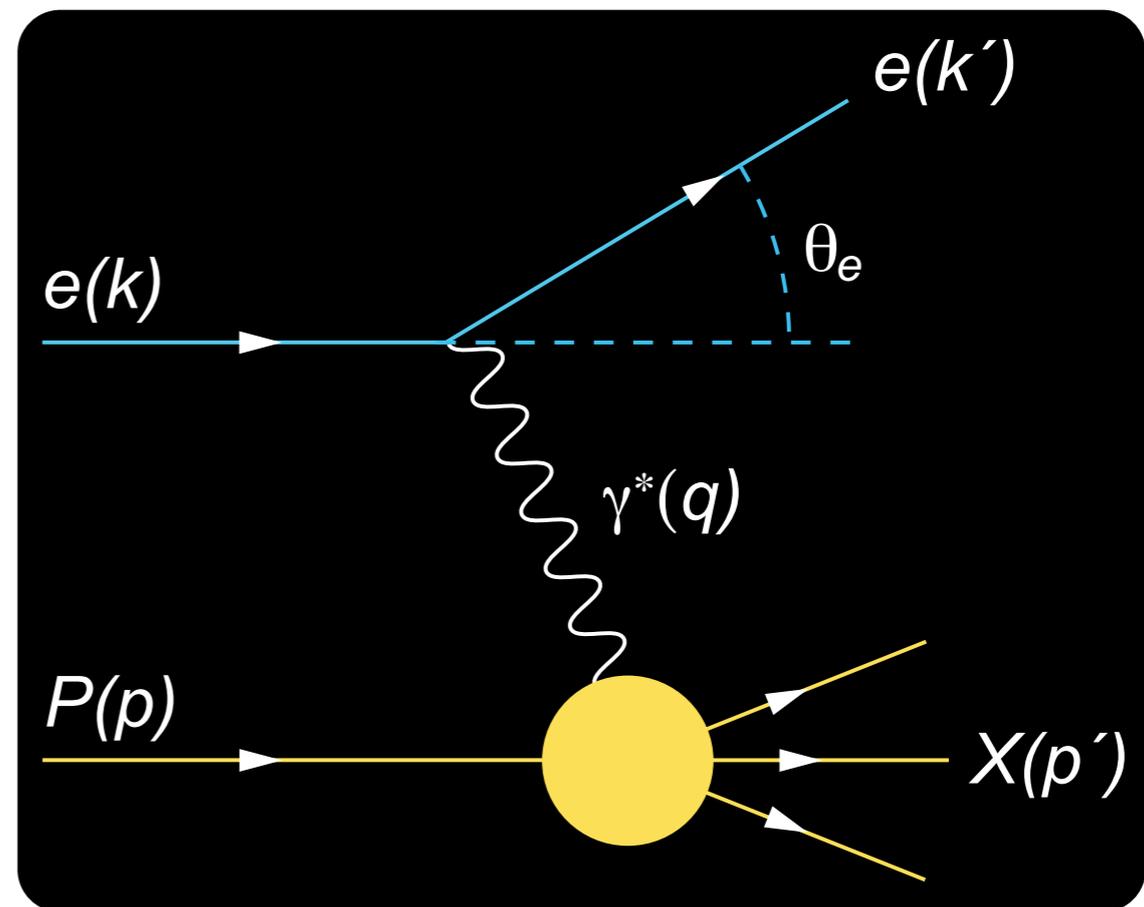
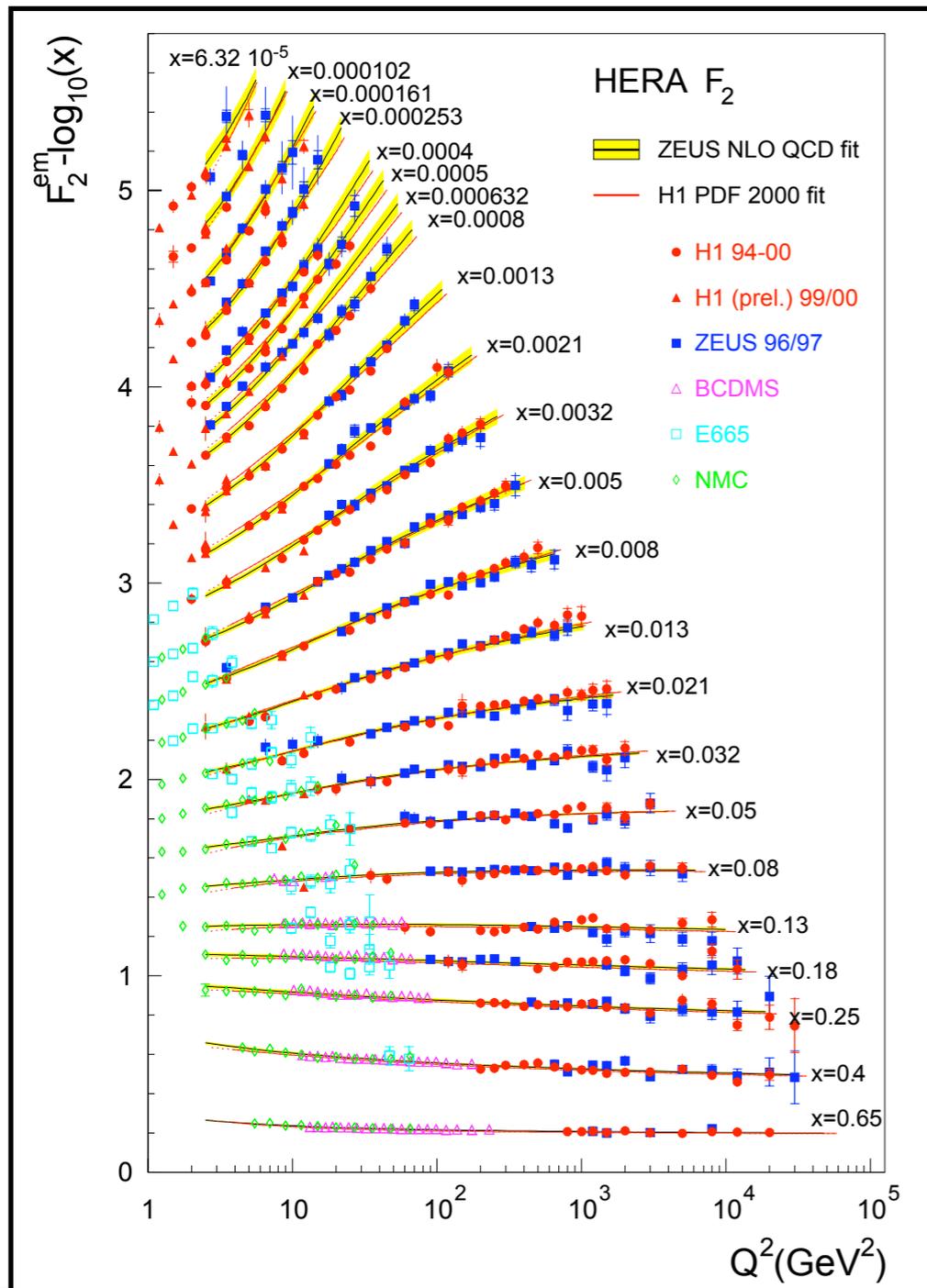
➔ $F_2(x, Q^2) \Rightarrow$ q and \bar{q} momentum distributions

➔ $F_L(x, Q^2) \Rightarrow$ gluon momentum distribution

No direct information on x, Q^2 from $p+A$ collisions !!

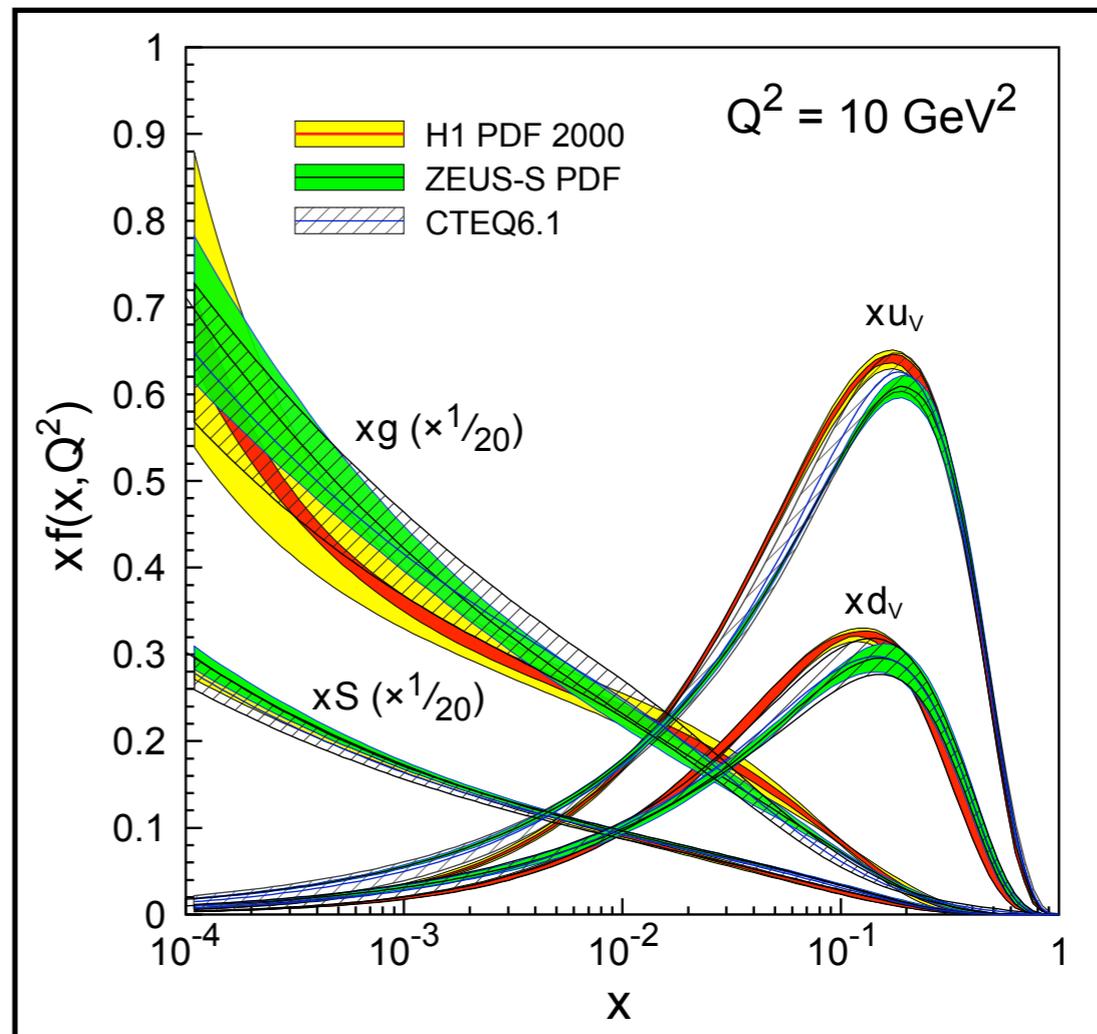
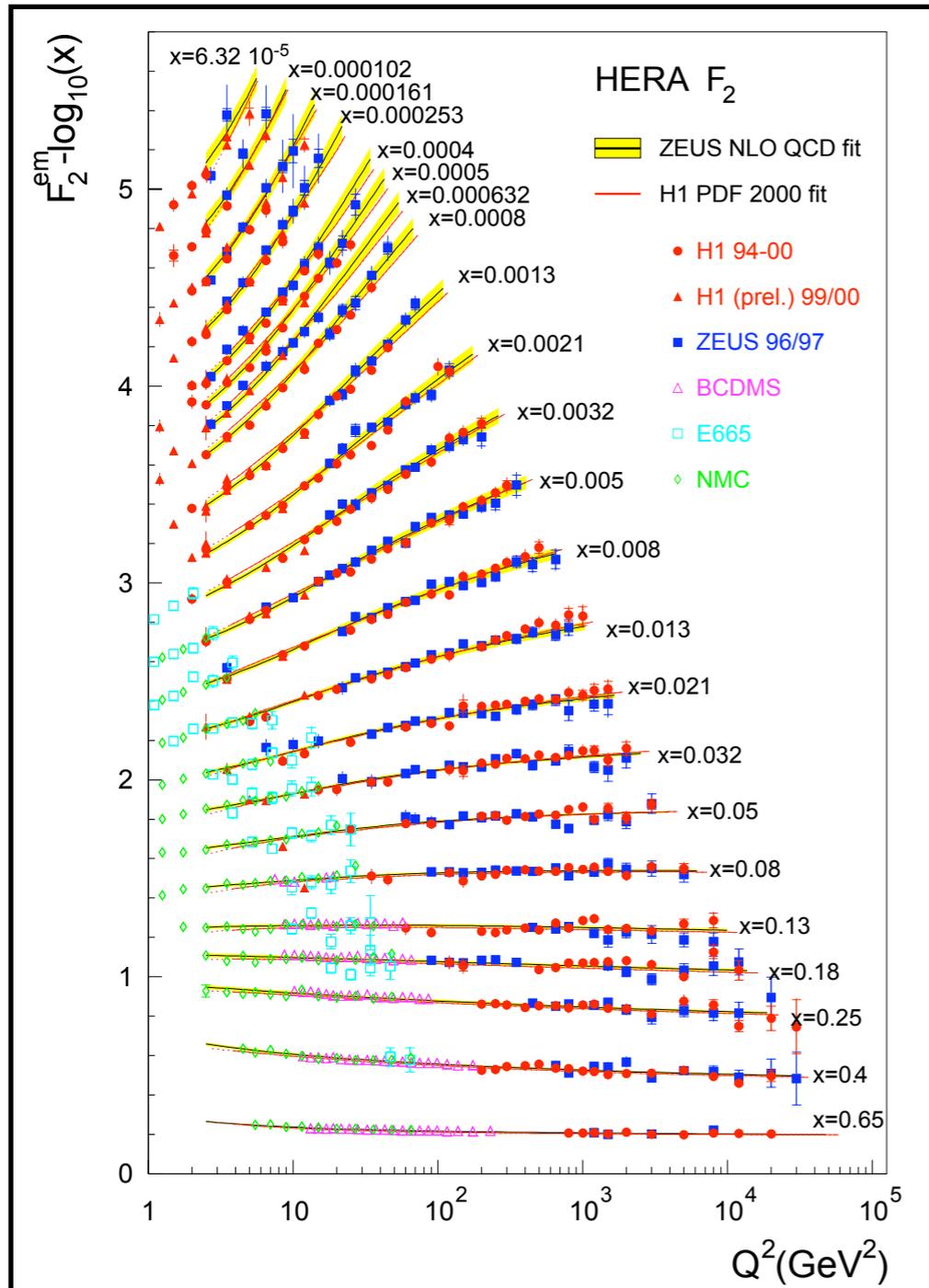
How to Measure the Glue ?

$$\frac{d^2\sigma^{ep\rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$



How to Measure the Glue ?

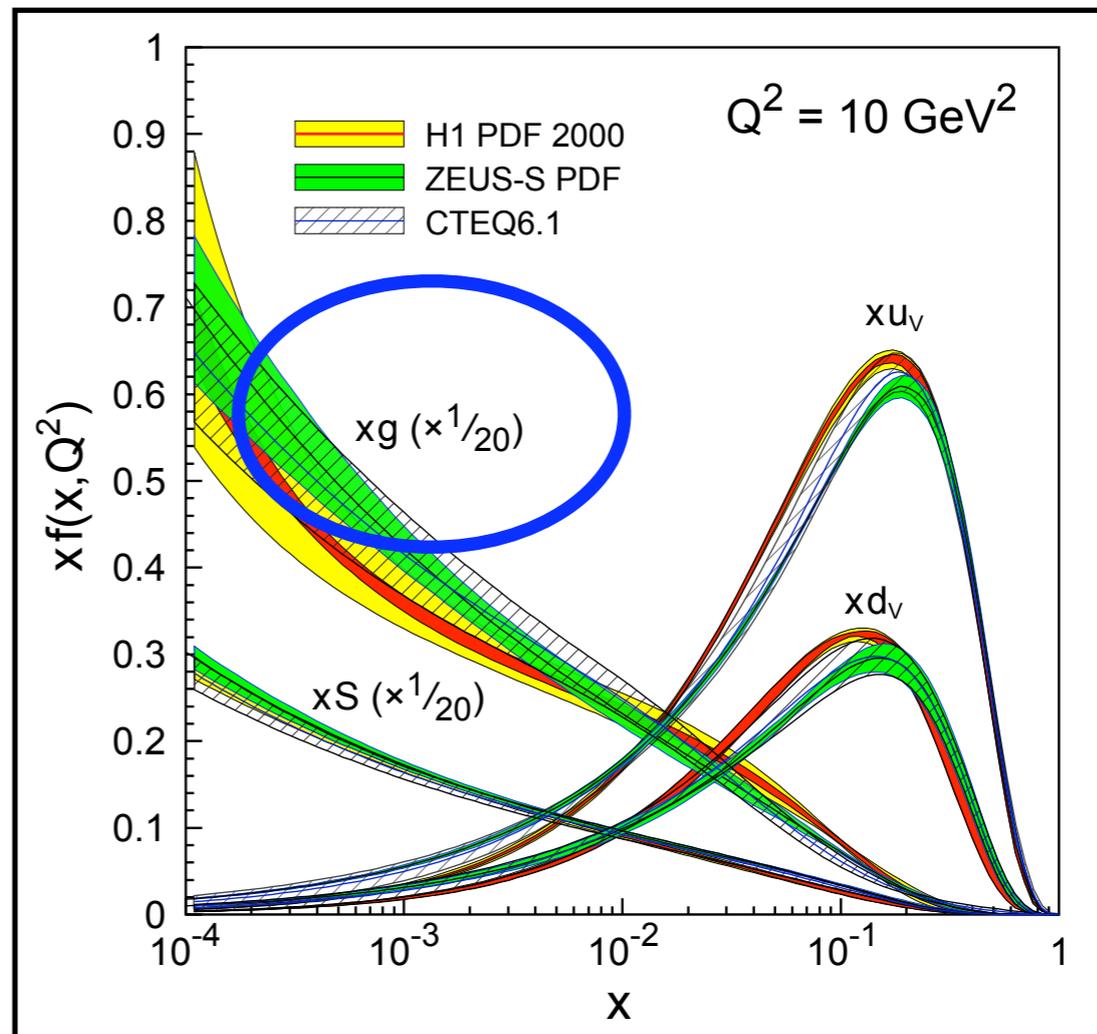
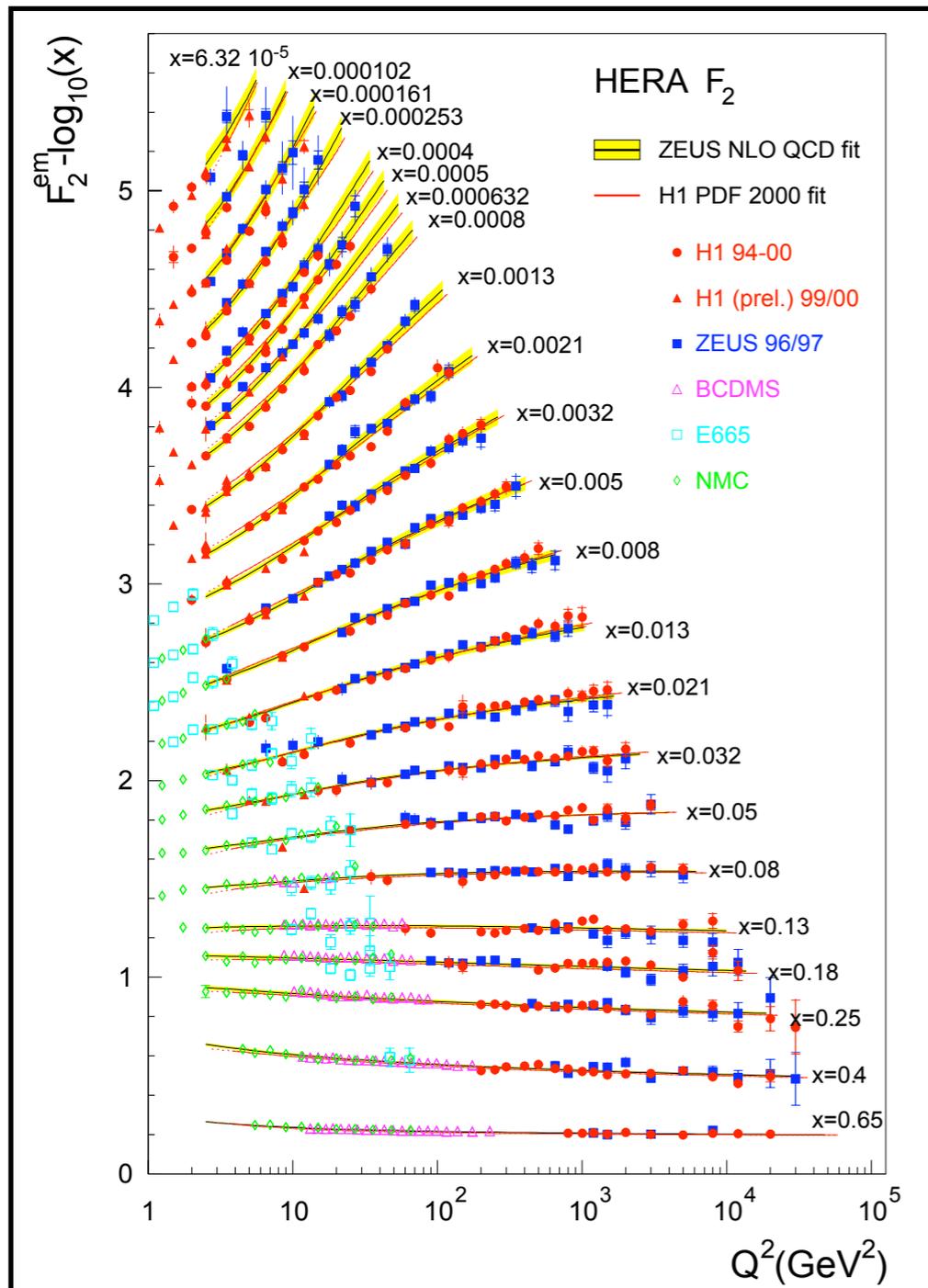
$$\frac{d^2 \sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi \alpha_{e.m.}^2}{x Q^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$



Scaling violation of $dF_2/d \ln Q^2$ and linear DGLAP evolution $\Rightarrow xG(x, Q^2)$

How to Measure the Glue ?

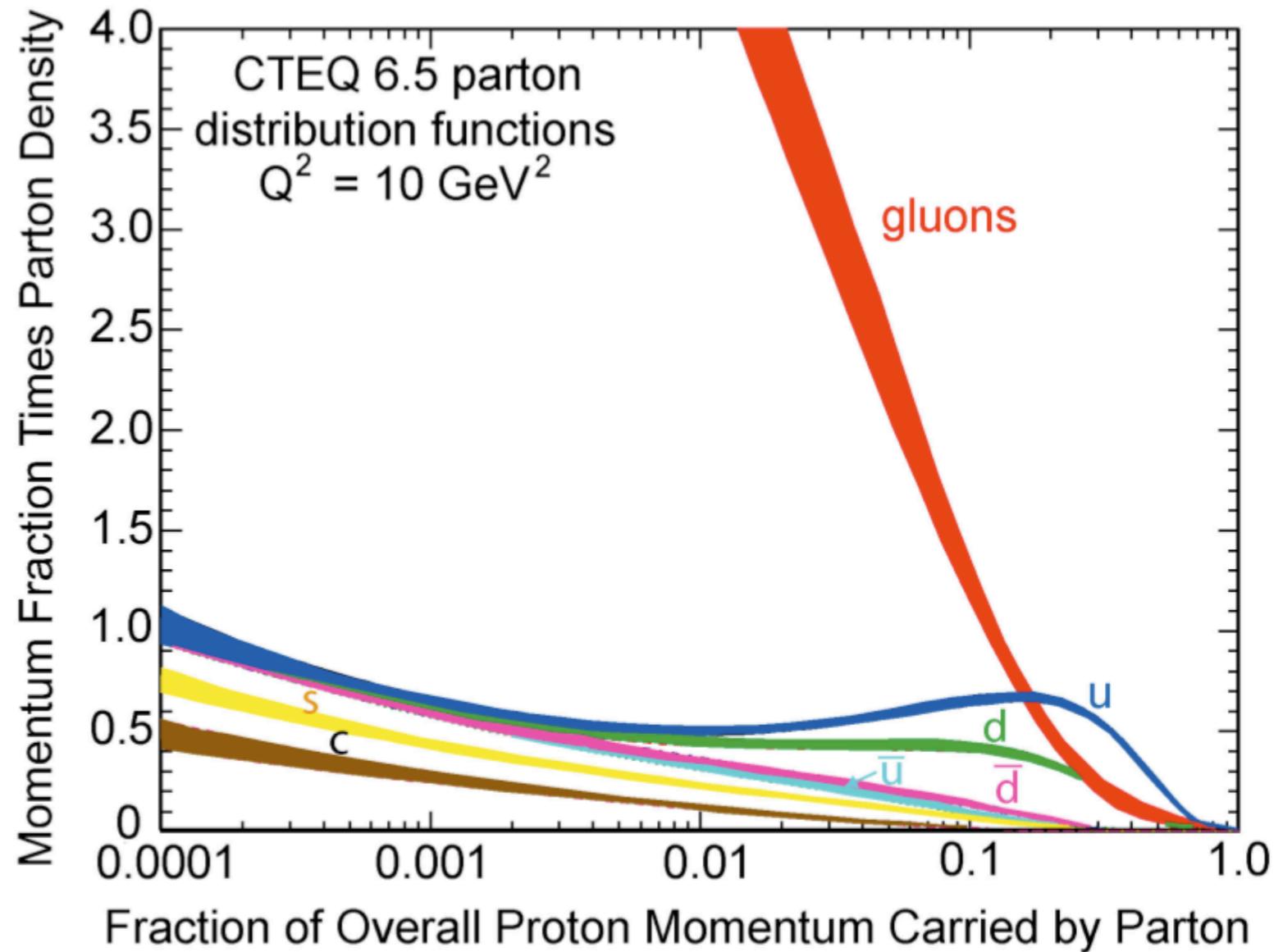
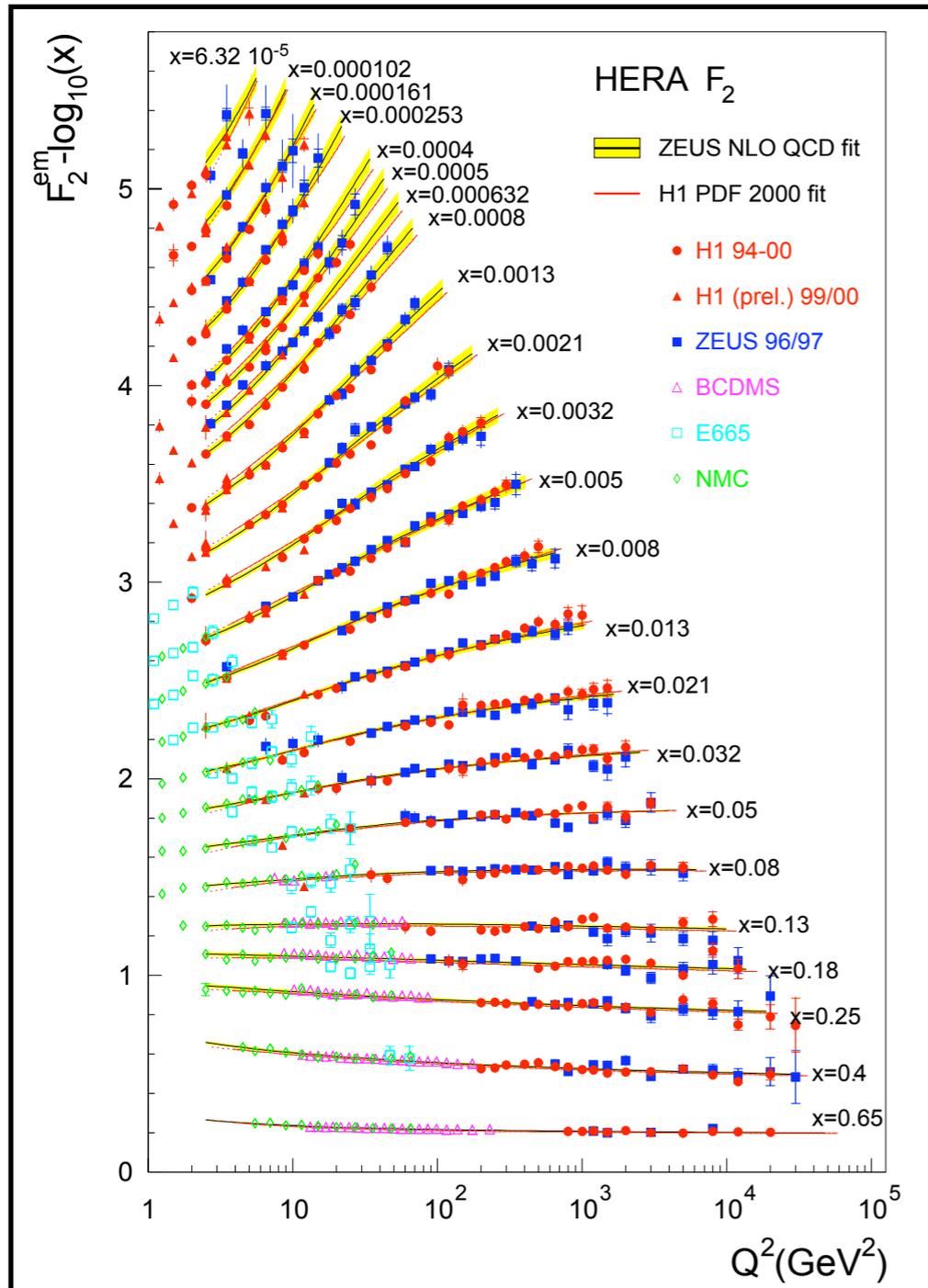
$$\frac{d^2\sigma^{ep\rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$



Scaling violation of $dF_2/d\ln Q^2$ and linear DGLAP evolution $\Rightarrow xG(x, Q^2)$

How to Measure the Glue ?

$$\frac{d^2\sigma^{ep\rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

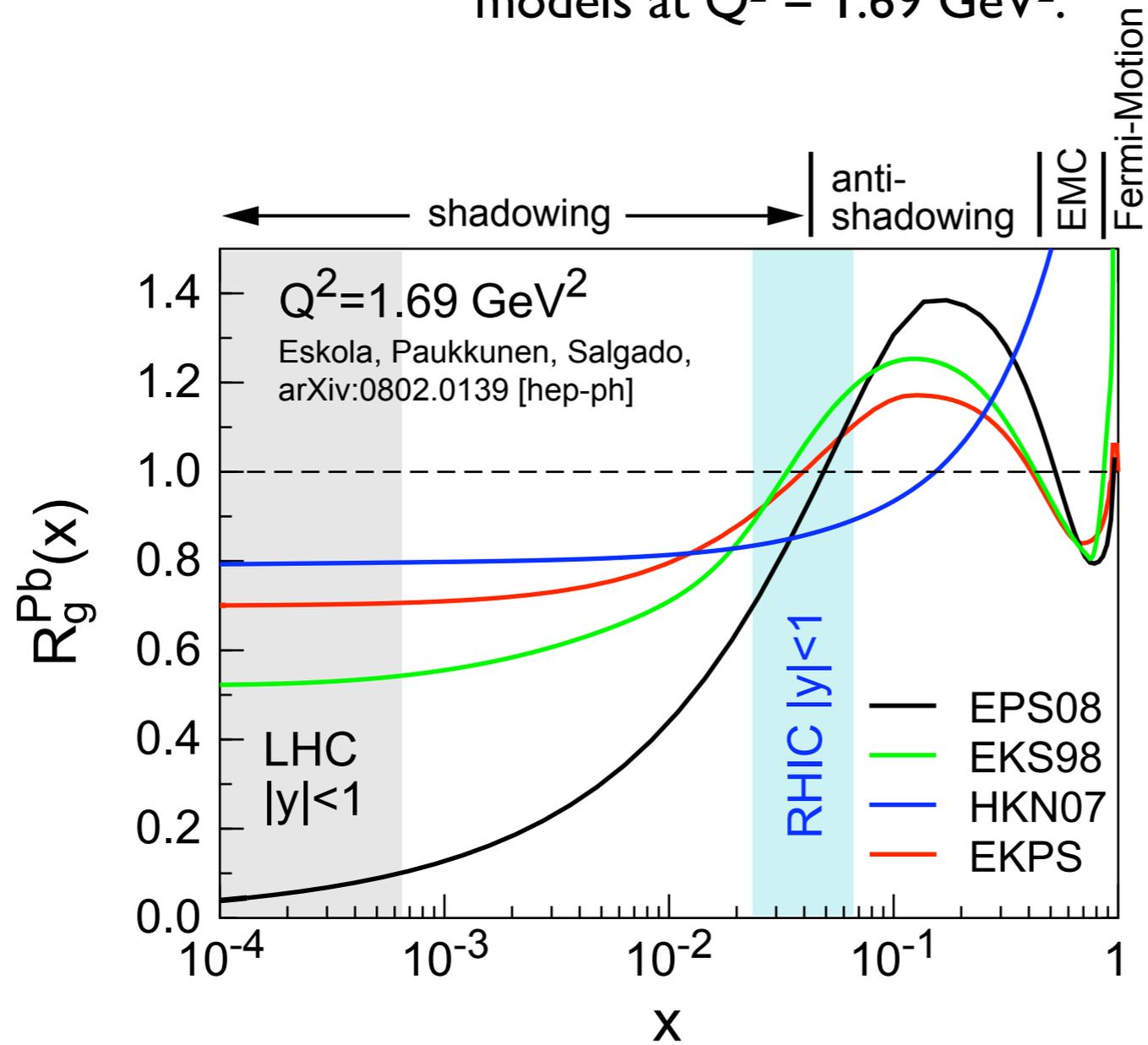


Scaling violation of $dF_2/d\ln Q^2$ and linear DGLAP evolution $\Rightarrow xG(x, Q^2)$

How to Measure the Glue ?

Important for RHIC and LHC:

Ratios of gluon distribution functions for Pb/p versus x from different models at $Q^2 = 1.69 \text{ GeV}^2$:



Models agree well for mid-rapidity RHIC, but discrepancies are there for forward RHIC rapidities as well as mid-rapidity at the LHC

$$R_i^A(x, Q^2) = \frac{f_i^A(x, Q^2)}{A f_i^{\text{nucleon}}(x, Q^2)}, \quad f_i = q, \bar{q}, g$$

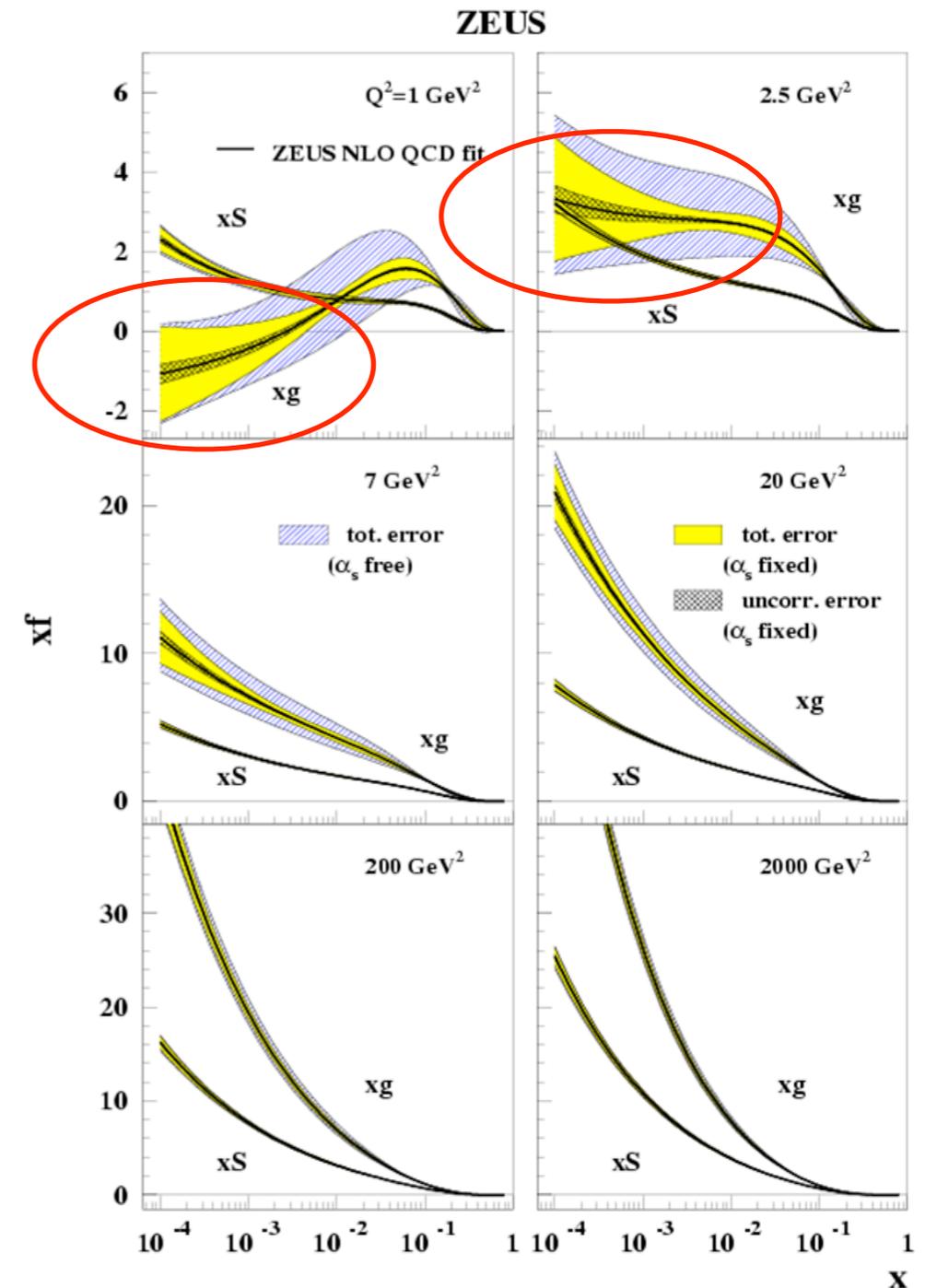
The problem with our current understanding

- Using the Linear DGLAP evolution model:

➔ Weird behaviour of xG at low- x and low Q^2 in HERA data

▶ xG goes negative !!

▶ $xS > xG$, though sea quarks come from gluon splitting ...



The problem with our current understanding

- Using the Linear DGLAP evolution model:

- ➔ Weird behaviour of xG at low- x and low Q^2 in HERA data

- ▶ xG goes negative !!

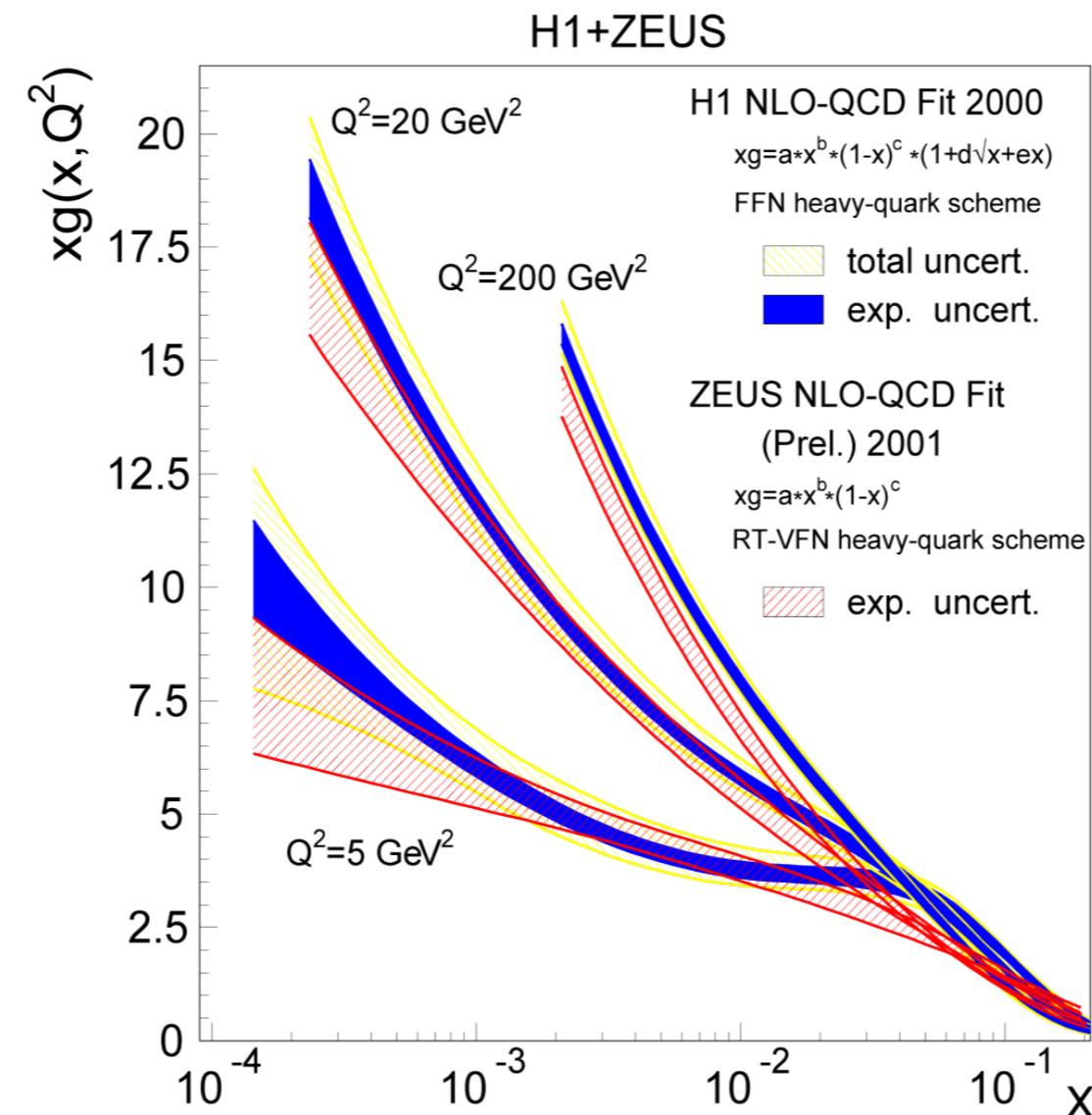
- ▶ $xS > xG$, though sea quarks come from gluon splitting ...

- More severe

- ➔ Linear evolution has a built-in high-energy “catastrophe”

- ➔ xG has rapid rise with decreasing x (and increasing Q^2) \Rightarrow violation of Froissart unitarity bound

- ▶ Must have saturation



The problem with our current understanding

- Using the Linear DGLAP evolution model:

- ➔ Weird behaviour of xG at low- x and low Q^2 in HERA data

- ▶ xG goes negative !!

- ▶ $xS > xG$, though sea quarks come from gluon splitting ...

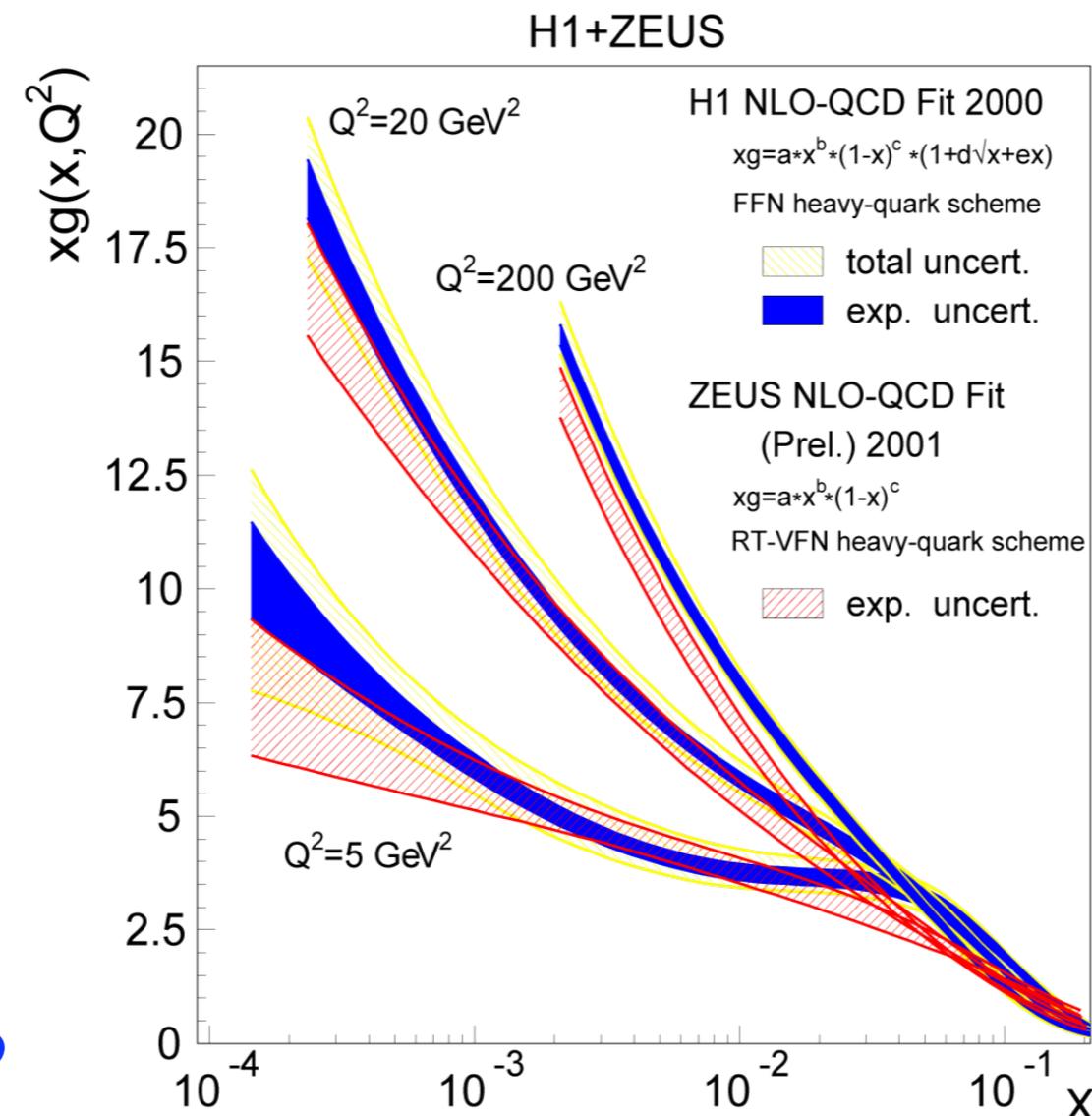
- More severe

- ➔ Linear evolution has a built-in high-energy “catastrophe”

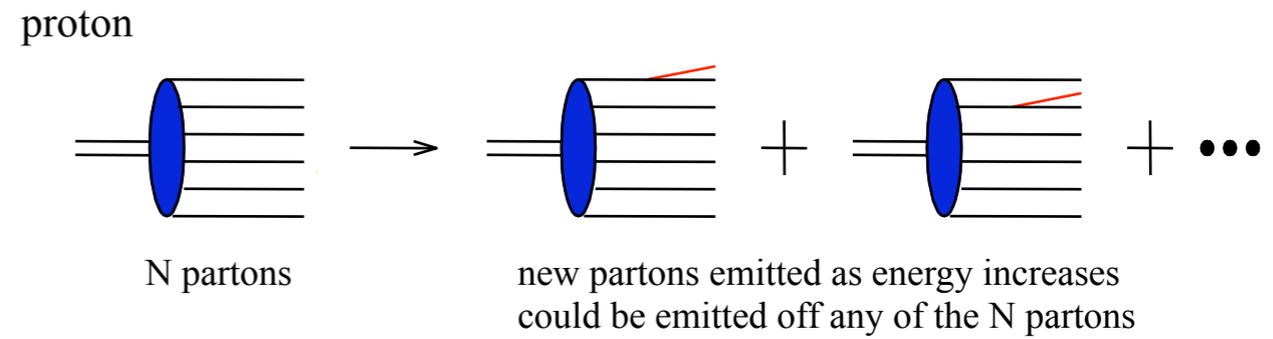
- ➔ xG has rapid rise with decreasing x (and increasing Q^2) \Rightarrow violation of Froissart unitarity bound

- ▶ Must have saturation

What's the underlying dynamics?



Non-linear QCD - Saturation

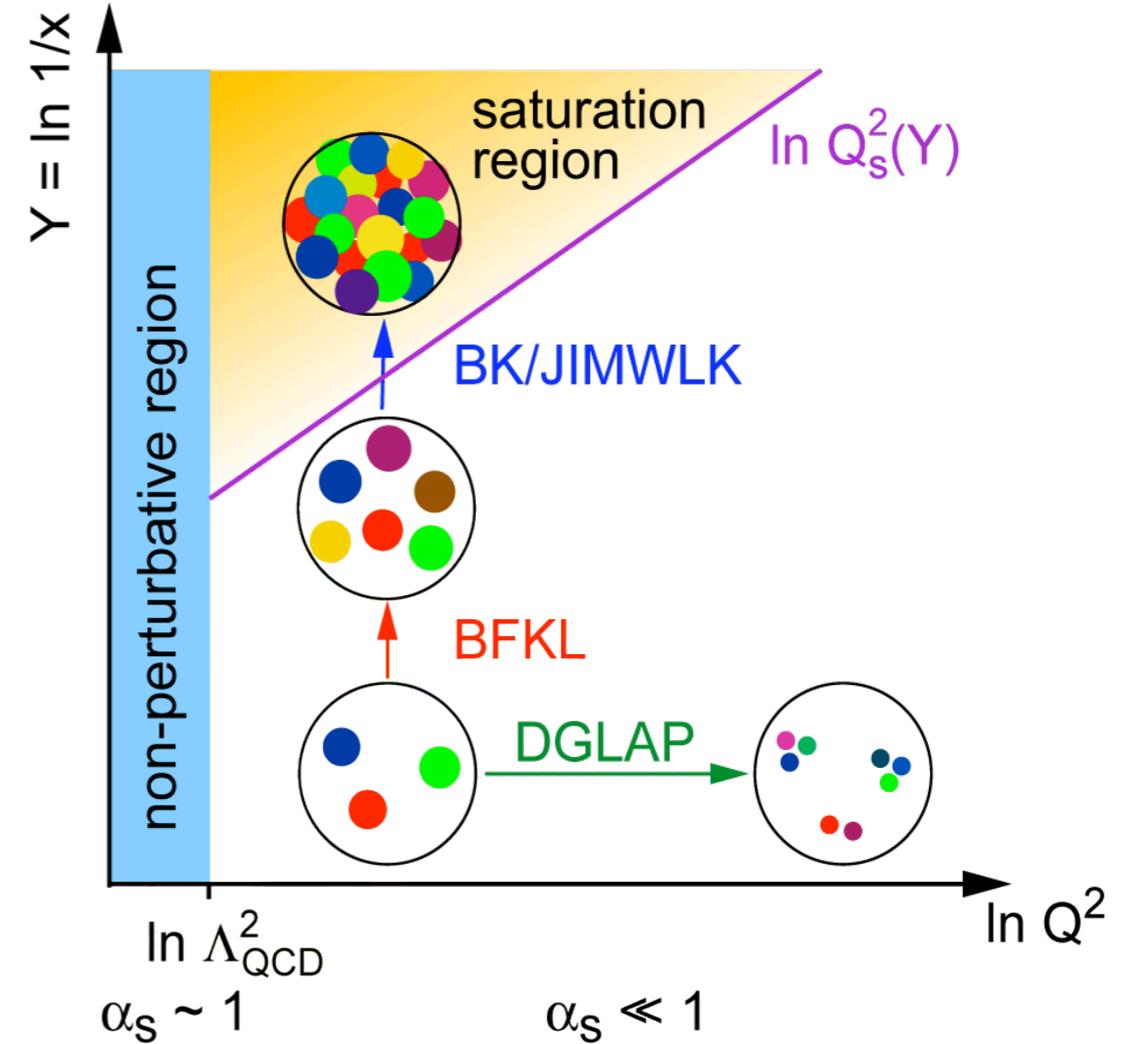
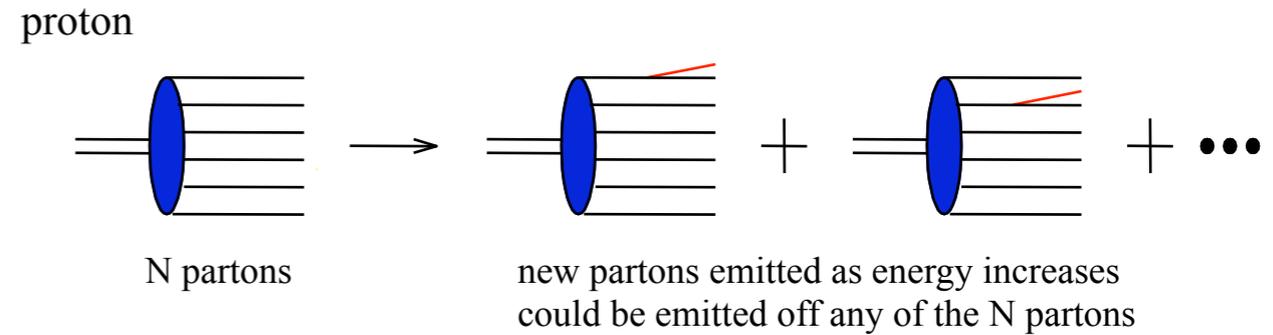


Non-linear QCD - Saturation

- **BFKL**: evolution in x

➔ linear

▶ explosion in colour field at low- x



Non-linear QCD - Saturation

- **BFKL**: evolution in x

➔ linear

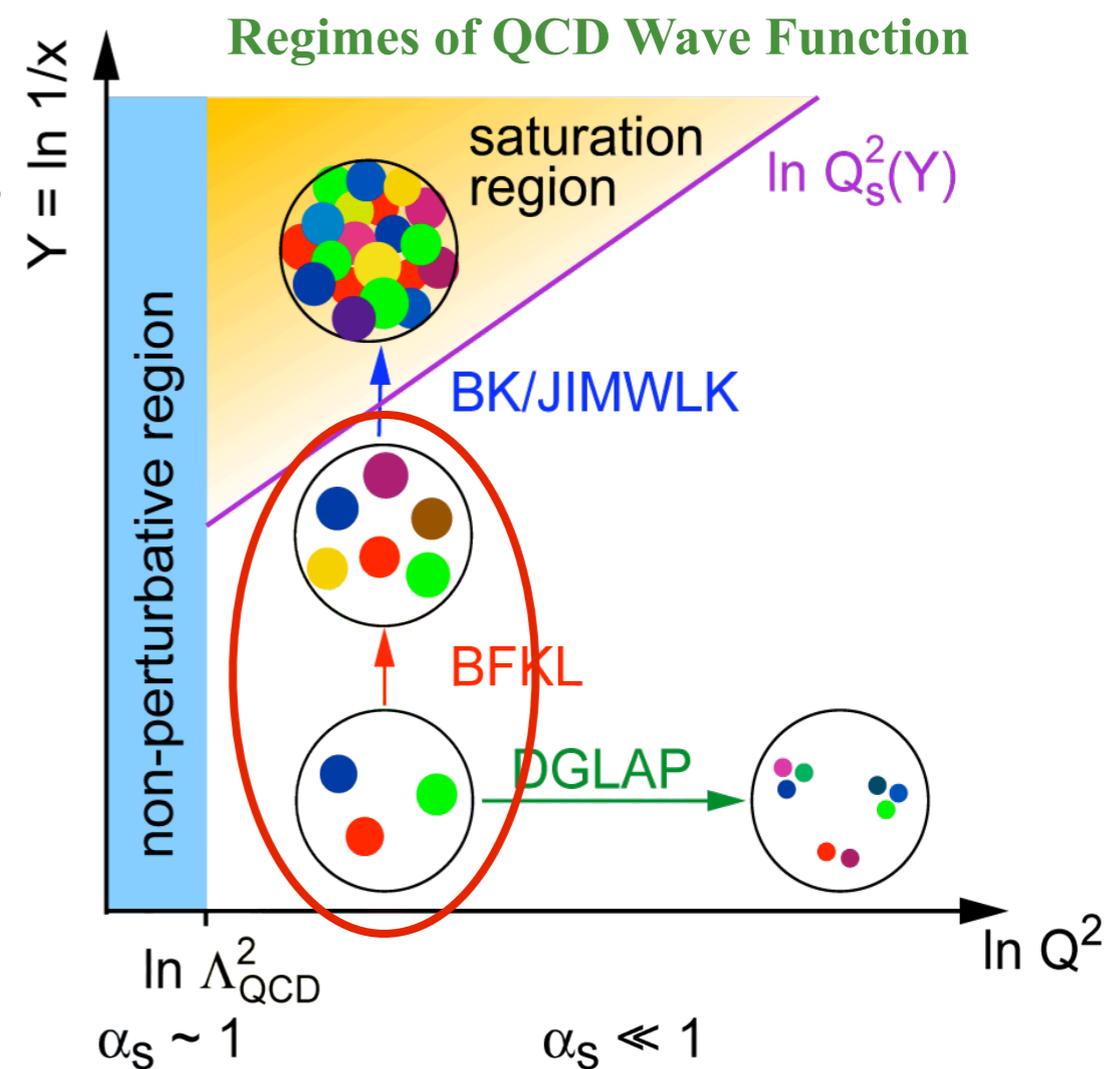
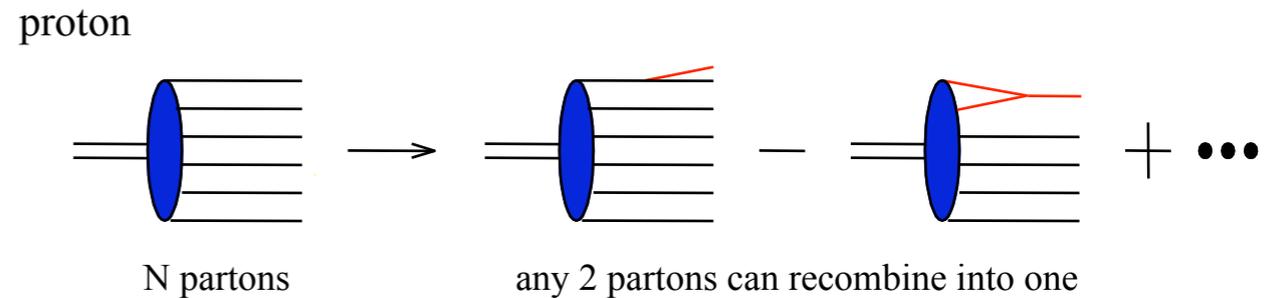
▶ explosion in colour field at low- x

- Non-linear **BK/JIMWLK** equations

➔ non-linearity \Rightarrow saturation

➔ characterised by the saturation scale, $Q_s(x,A)$

➔ arises naturally in the Colour Glass Condensate (CGC) EFT



Non-linear QCD - Saturation

- **BFKL**: evolution in x

➔ linear

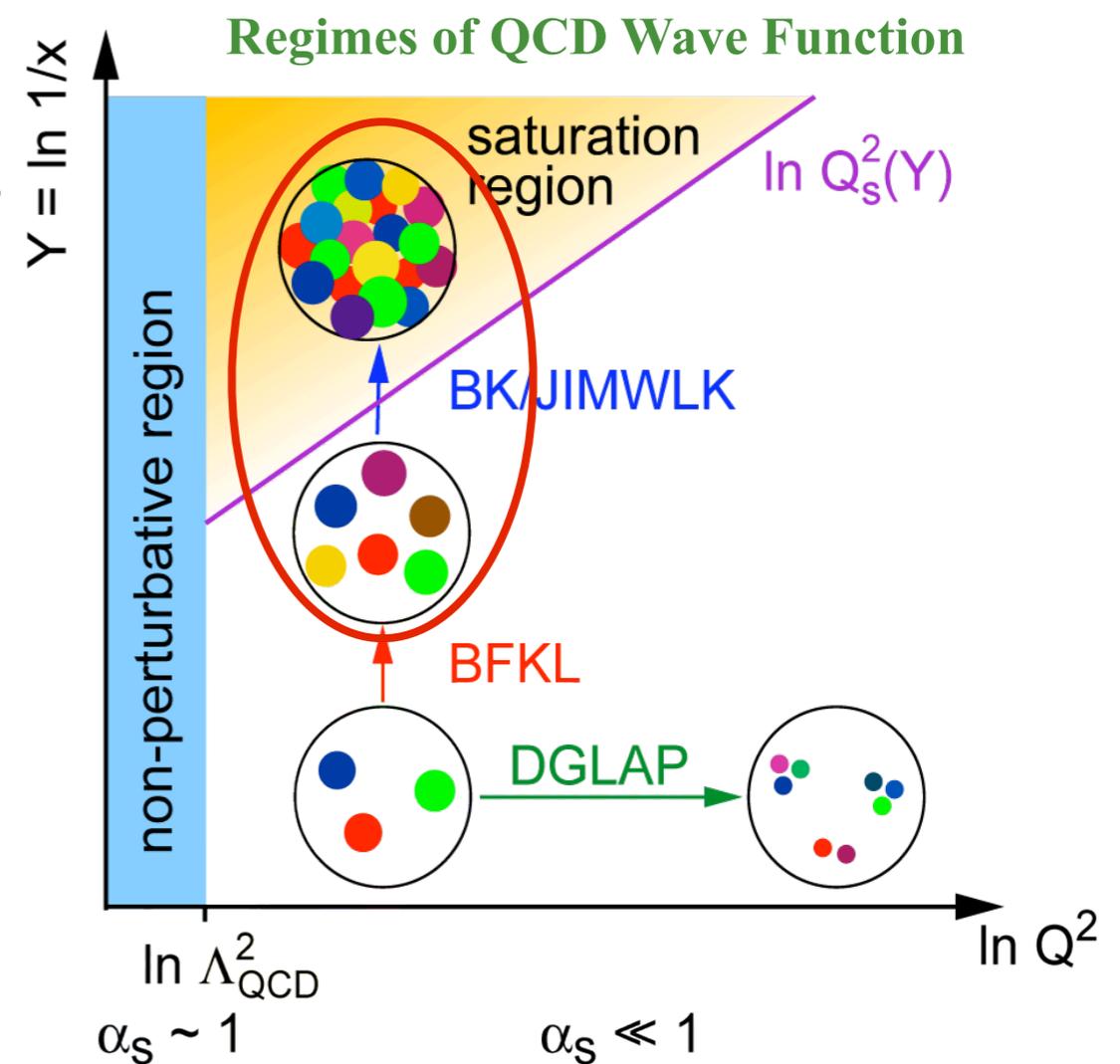
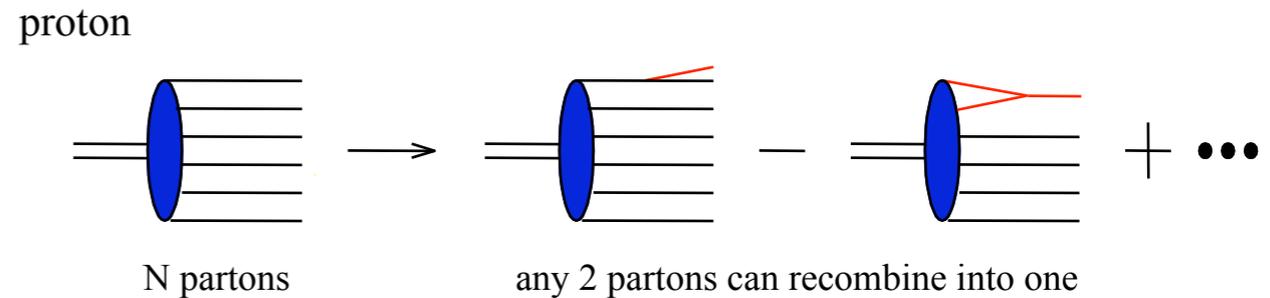
▶ explosion in colour field at low- x

- Non-linear **BK/JIMWLK** equations

➔ non-linearity \Rightarrow saturation

➔ characterised by the saturation scale, $Q_s(x,A)$

➔ arises naturally in the Colour Glass Condensate (CGC) EFT



Why study $e+A$ instead of $e+p$?

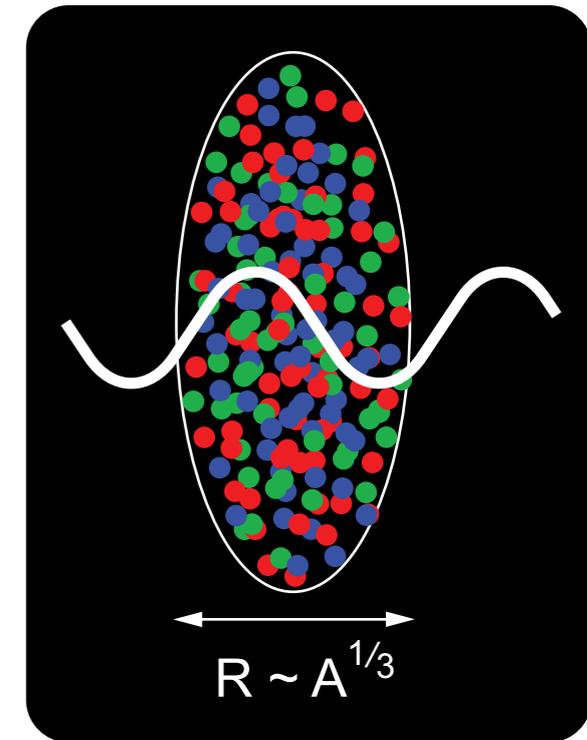
Enhancing Saturation Effects:

Scattering of electrons off nuclei:

Probes interact over distances $L \sim (2m_N x)^{-1}$

For $L > 2 R_A \sim A^{1/3}$ probe cannot distinguish between nucleons in front or back of nuclei

\Rightarrow Probe interacts *coherently* with all nucleons



Why study $e+A$ instead of $e+p$?

Enhancing Saturation Effects:

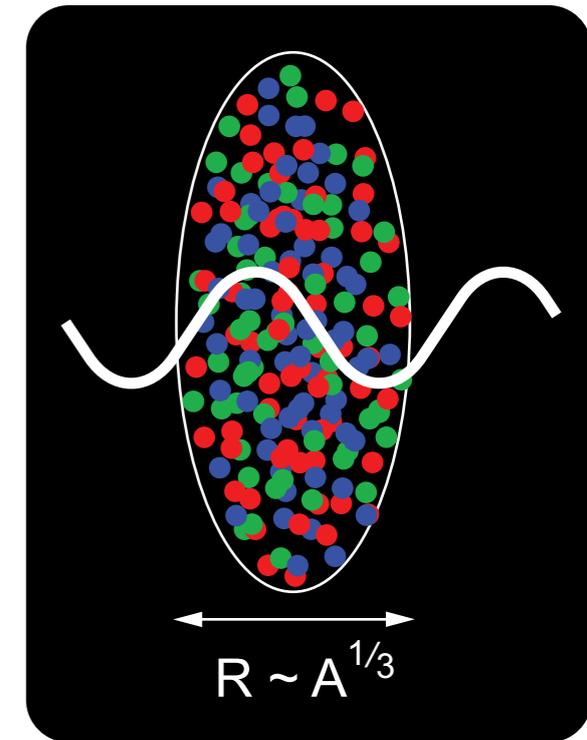
Scattering of electrons off nuclei:

Probes interact over distances $L \sim (2m_N x)^{-1}$

For $L > 2 R_A \sim A^{1/3}$ probe cannot distinguish between nucleons in front or back of nuclei

\Rightarrow Probe interacts *coherently* with all nucleons

A probe of transverse resolution $1/Q^2 (\ll \Lambda^2_{\text{QCD}}) \sim 1 \text{ fm}^2$ will experience large colour charge fluctuations. This kick experienced in a random walk is the saturation scale.



Why study $e+A$ instead of $e+p$?

Enhancing Saturation Effects:

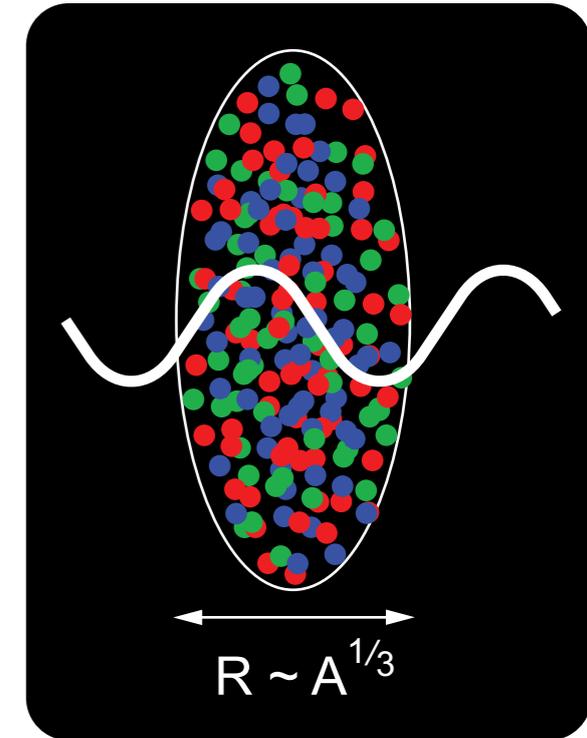
Scattering of electrons off nuclei:

Probes interact over distances $L \sim (2m_N x)^{-1}$

For $L > 2 R_A \sim A^{1/3}$ probe cannot distinguish between nucleons in front or back of nuclei

⇒ Probe interacts *coherently* with all nucleons

A probe of transverse resolution $1/Q^2 (\ll \Lambda^2_{\text{QCD}}) \sim 1 \text{ fm}^2$ will experience large colour charge fluctuations. This kick experienced in a random walk is the saturation scale.



$$Q_s^2 \propto \frac{\alpha_s x G(x, Q_s^2)}{\pi R_A^2} \quad \text{HERA : } xG \propto \frac{1}{x^{1/3}} \quad \text{A dependence : } xG_A \propto A$$

Nuclear “Oomph” Factor: $(Q_s^A)^2 \approx c Q_0^2 \left(\frac{A}{x}\right)^{1/3}$

Why study $e+A$ instead of $e+p$?

Enhancing Saturation Effects:

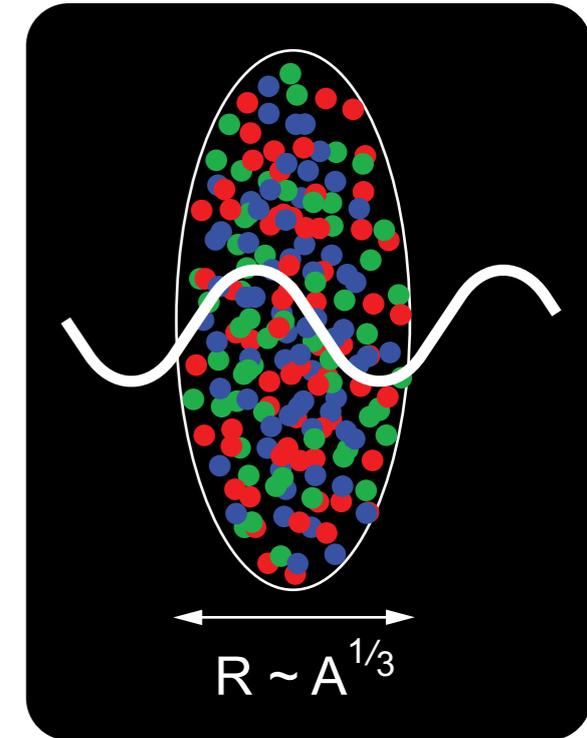
Scattering of electrons off nuclei:

Probes interact over distances $L \sim (2m_N x)^{-1}$

For $L > 2 R_A \sim A^{1/3}$ probe cannot distinguish between nucleons in front or back of nuclei

\Rightarrow Probe interacts *coherently* with all nucleons

A probe of transverse resolution $1/Q^2 (\ll \Lambda^2_{\text{QCD}}) \sim 1 \text{ fm}^2$ will experience large colour charge fluctuations. This kick experienced in a random walk is the saturation scale.



$$Q_s^2 \propto \frac{\alpha_s x G(x, Q_s^2)}{\pi R_A^2} \quad \text{HERA : } xG \propto \frac{1}{x^{1/3}} \quad \text{A dependence : } xG_A \propto A$$

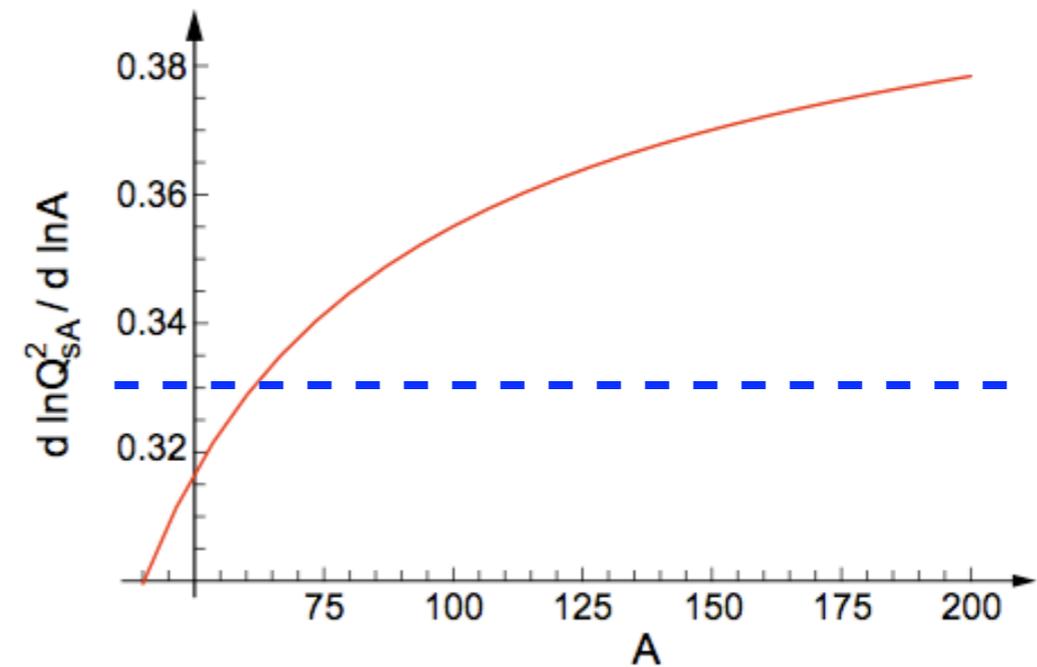
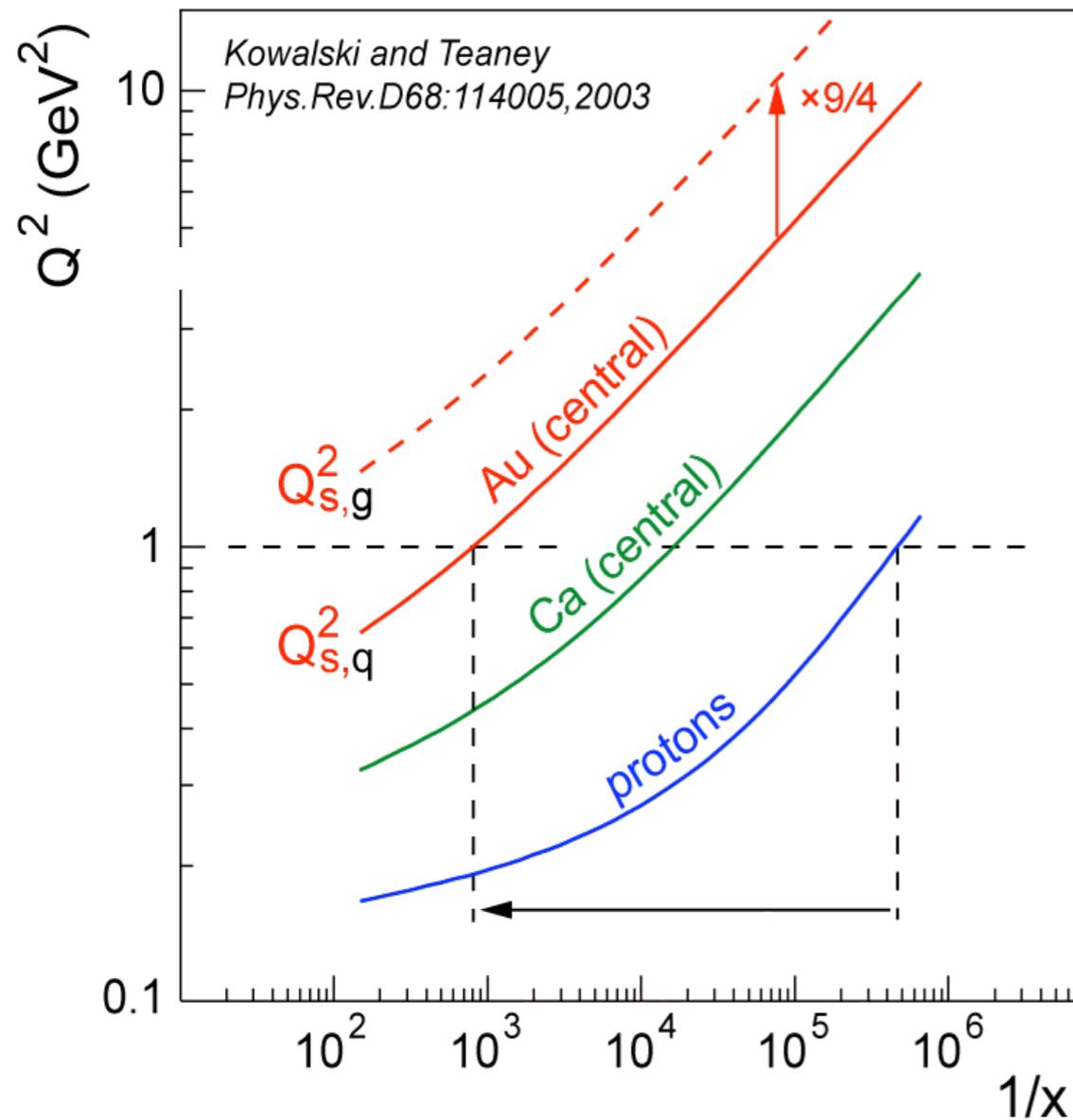
Nuclear “Oomph” Factor: $(Q_s^A)^2 \approx c Q_0^2 \left(\frac{A}{x}\right)^{1/3}$

Enhancement of Q_s with A : \Rightarrow non-linear QCD regime reached at significantly lower energy in $e+A$ than in $e+p$

The Nuclear “Oomph” factor

More sophisticated analyses \Rightarrow confirm (exceed) pocket formula

(e.g. Kowalski, Lappi and Venugopalan, PRL 100, 022303 (2008); Armesto et al., PRL 94:022002; Kowalski, Teaney, PRD 68:114005)

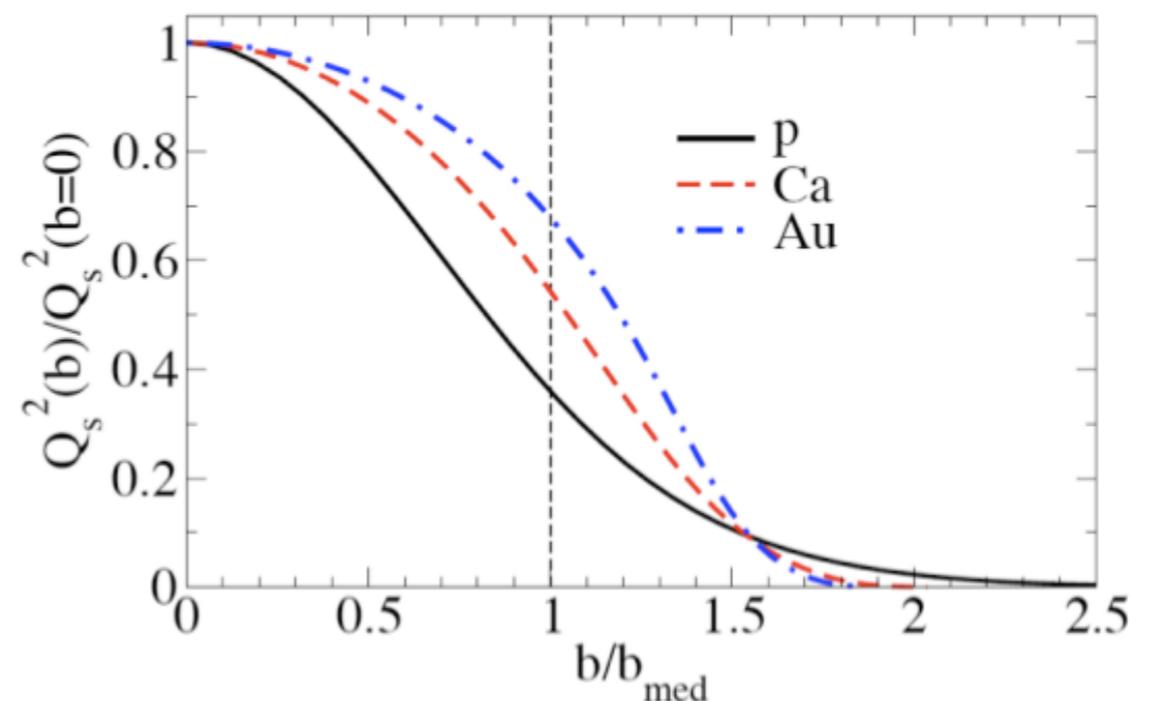
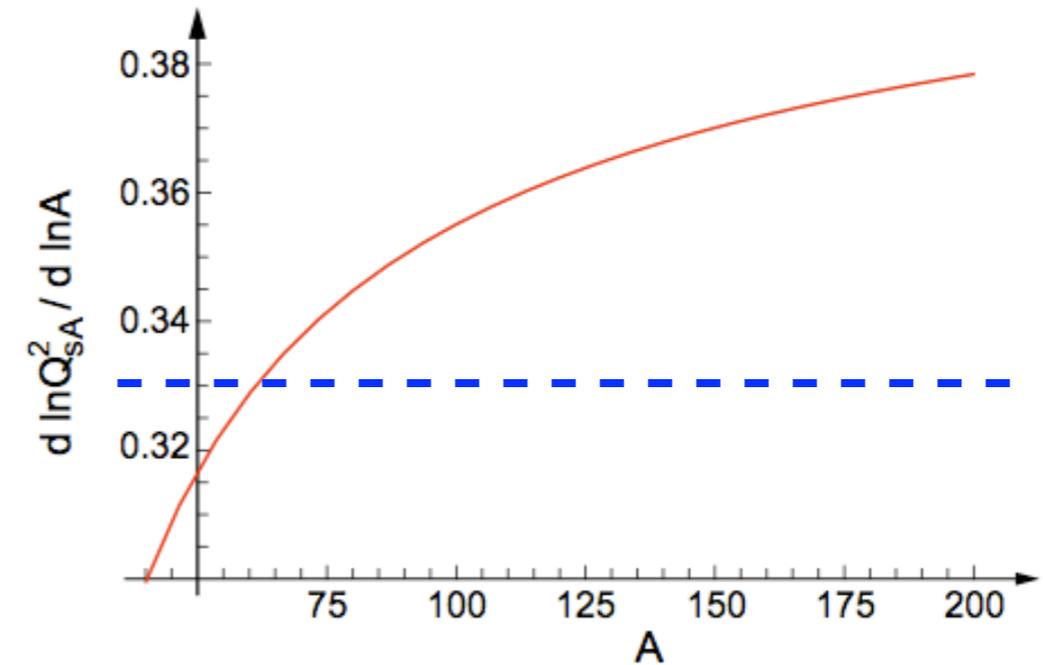
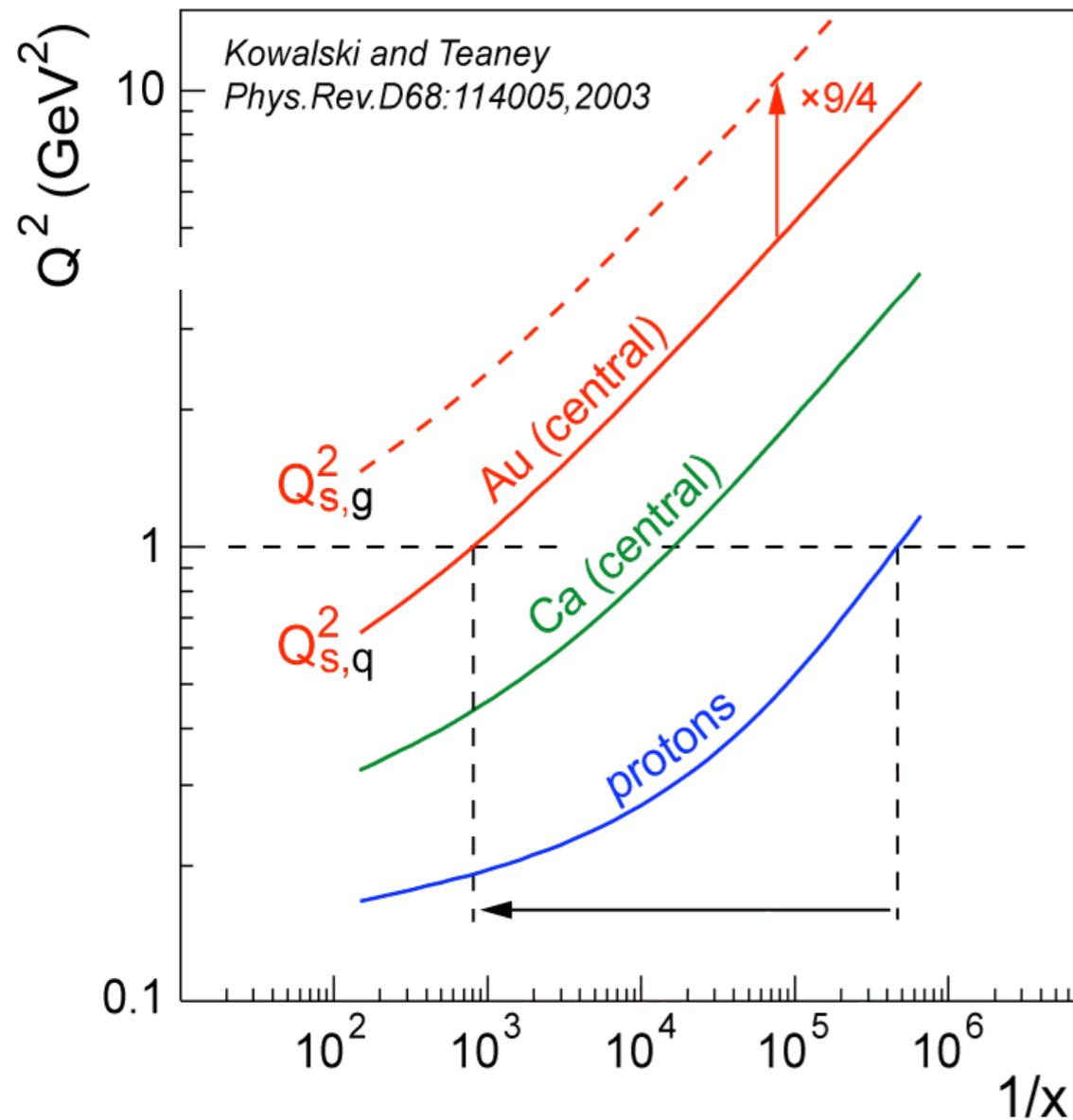


Models need to use realistic b -dependence for nuclei and nucleons
 $\Rightarrow b = 0$ for proton $\neq b_{\text{med}}$

The Nuclear ‘‘Oomph’’ factor

More sophisticated analyses \Rightarrow confirm (exceed) pocket formula

(e.g. Kowalski, Lappi and Venugopalan, PRL 100, 022303 (2008); Armesto et al., PRL 94:022002; Kowalski, Teaney, PRD 68:114005)

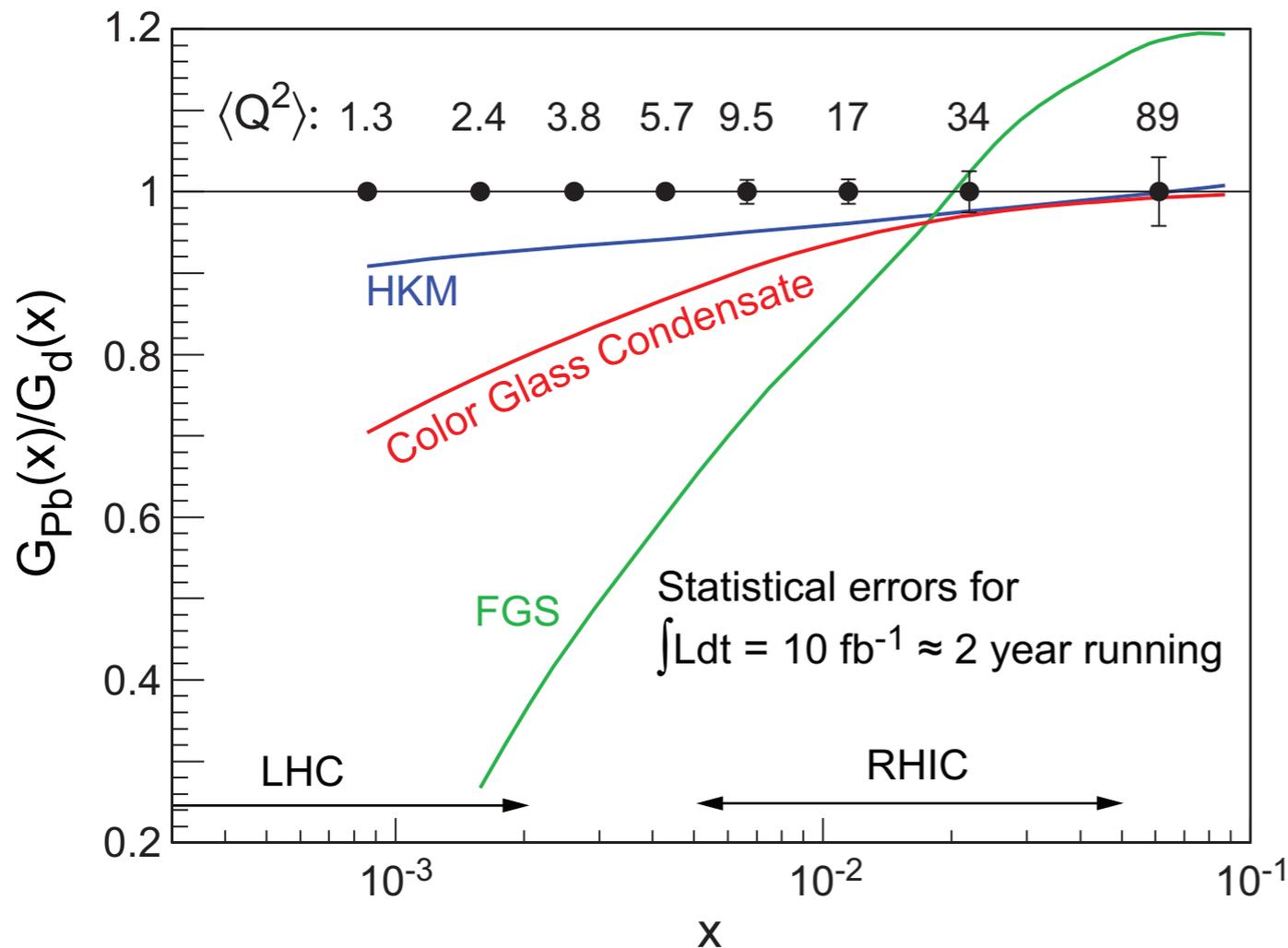


Key Measurements in $e+A$

- *Momentum distribution of gluons $G(x, Q^2)$*
 - ➔ Extract via scaling violation in F_2 : $\delta F_2 / \delta \ln Q^2$
 - ➔ Direct measurement: $F_L \sim xG(x, Q^2)$ (requires \sqrt{s} scan)
 - ➔ 2+1 jet rates
 - ➔ Inelastic vector meson production (e.g. J/ψ)
 - ➔ Diffractive vector meson production $\sim [xG(x, Q^2)]^2$

Example of Key Measurements:

$$\frac{d^2\sigma^{ep\rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$



HKM and FGS are "standard" shadowing parameterizations that are evolved with DGLAP

$$F_L \sim \alpha_s xG(x, Q^2)$$

requires \sqrt{s} scan, $Q^2/xs = y$

Here:

$$\begin{aligned} \int Ldt &= 4/A \text{ fb}^{-1} \text{ (10+100) GeV} \\ &= 4/A \text{ fb}^{-1} \text{ (10+50) GeV} \\ &= 2/A \text{ fb}^{-1} \text{ (5+50) GeV} \end{aligned}$$

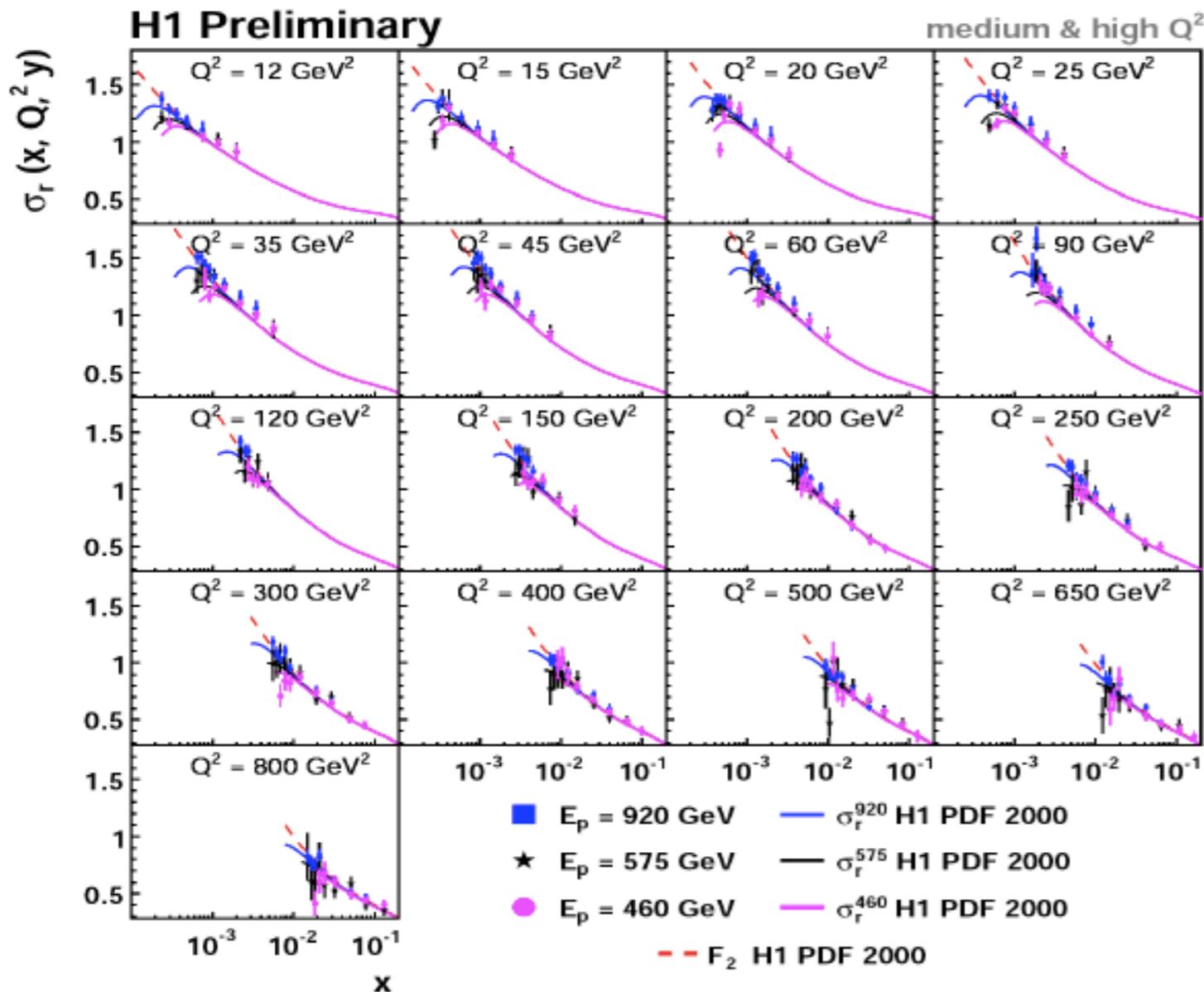
statistical error only

Syst. studies of $F_L(A, x, Q^2)$:

- $xG(x, Q^2)$ with great precision
- Distinguish between models

Example of Key Measurements:

$$\frac{d^2\sigma^{ep\rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[\left(1 - y + \frac{y^2}{2}\right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$



$$F_L \sim \alpha_s xG(x, Q^2)$$

requires \sqrt{s} scan, $Q^2/xs = y$

Here:

$$\begin{aligned} \int L dt &= 4/A \text{ fb}^{-1} (10+100) \text{ GeV} \\ &= 4/A \text{ fb}^{-1} (10+50) \text{ GeV} \\ &= 2/A \text{ fb}^{-1} (5+50) \text{ GeV} \end{aligned}$$

statistical error only

Syst. studies of $F_L(A, x, Q^2)$:

- $xG(x, Q^2)$ with great precision
- Distinguish between models

Preliminary F_L measurements

Key Measurements in $e+A$

- *Momentum distribution of gluons $G(x, Q^2)$*
 - ➔ Extract via scaling violation in F_2 : $\delta F_2 / \delta \ln Q^2$
 - ➔ Direct measurement: $F_L \sim xG(x, Q^2)$ (requires \sqrt{s} scan)
 - ➔ 2+1 jet rates
 - ➔ Inelastic vector meson production (e.g. J/ψ)
 - ➔ Diffractive vector meson production $\sim [xG(x, Q^2)]^2$

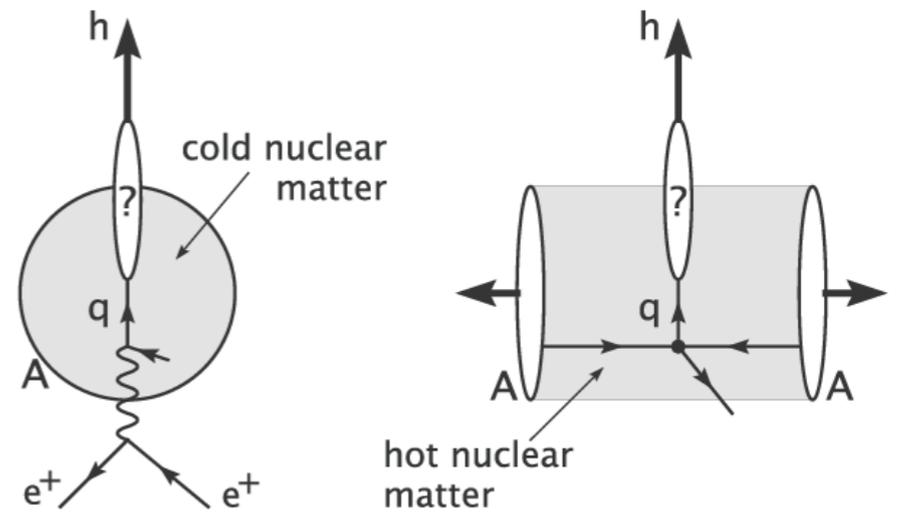
Key Measurements in $e+A$

- **Momentum distribution of gluons $G(x, Q^2)$**
 - ➔ Extract via scaling violation in F_2 : $\delta F_2 / \delta \ln Q^2$
 - ➔ Direct measurement: $F_L \sim xG(x, Q^2)$ (requires \sqrt{s} scan)
 - ➔ 2+1 jet rates
 - ➔ Inelastic vector meson production (e.g. J/ψ)
 - ➔ Diffractive vector meson production $\sim [xG(x, Q^2)]^2$
- **Space-time distributions of gluons in matter**
 - ➔ Exclusive final states (e.g. vector meson production $\rho, J/\psi$)
 - ➔ Deep Virtual Compton Scattering (DVCS) - $\sigma \sim A^{4/3}$
 - ➔ F_2, F_L for various A and impact parameter dependence

Key Measurements in $e+A$

- **Momentum distribution of gluons $G(x, Q^2)$**
 - ➔ Extract via scaling violation in F_2 : $\delta F_2 / \delta \ln Q^2$
 - ➔ Direct measurement: $F_L \sim xG(x, Q^2)$ (requires \sqrt{s} scan)
 - ➔ 2+1 jet rates
 - ➔ Inelastic vector meson production (e.g. J/ψ)
 - ➔ Diffractive vector meson production $\sim [xG(x, Q^2)]^2$
- **Space-time distributions of gluons in matter**
 - ➔ Exclusive final states (e.g. vector meson production ρ , J/ψ)
 - ➔ Deep Virtual Compton Scattering (DVCS) - $\sigma \sim A^{4/3}$
 - ➔ F_2, F_L for various A and impact parameter dependence
- **Interaction of fast probes with *gluonic* medium?**
 - ➔ Hadronization, Fragmentation
 - ➔ Energy loss (charm, bottom!)

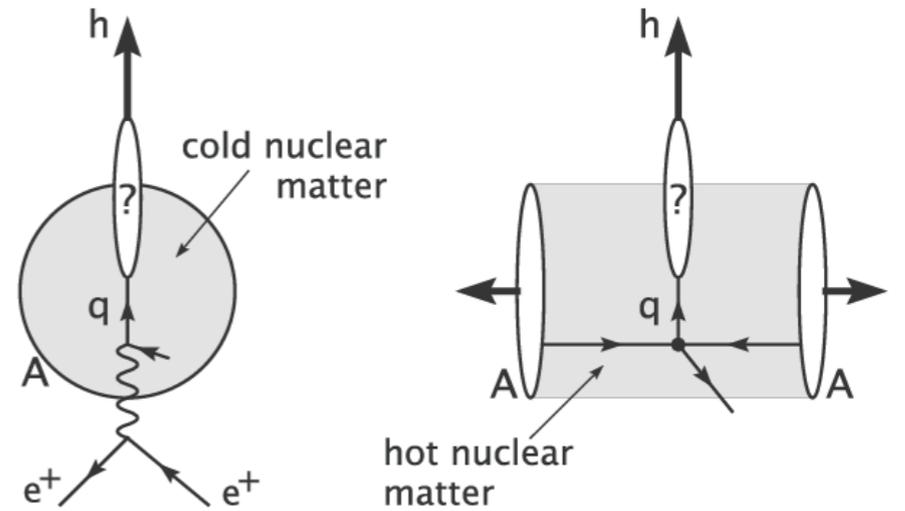
Hadronization and Energy Loss



Hadronization and Energy Loss

nDIS:

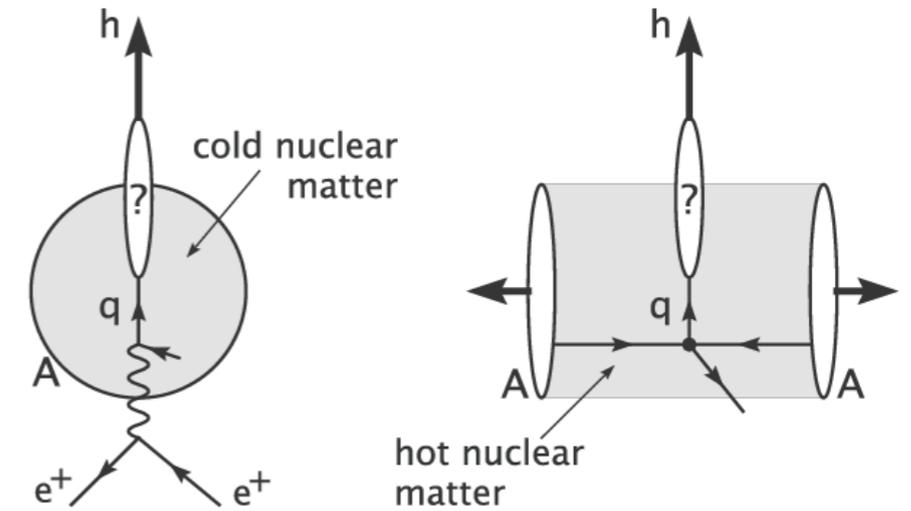
- Clean measurement in 'cold' nuclear matter



Hadronization and Energy Loss

nDIS:

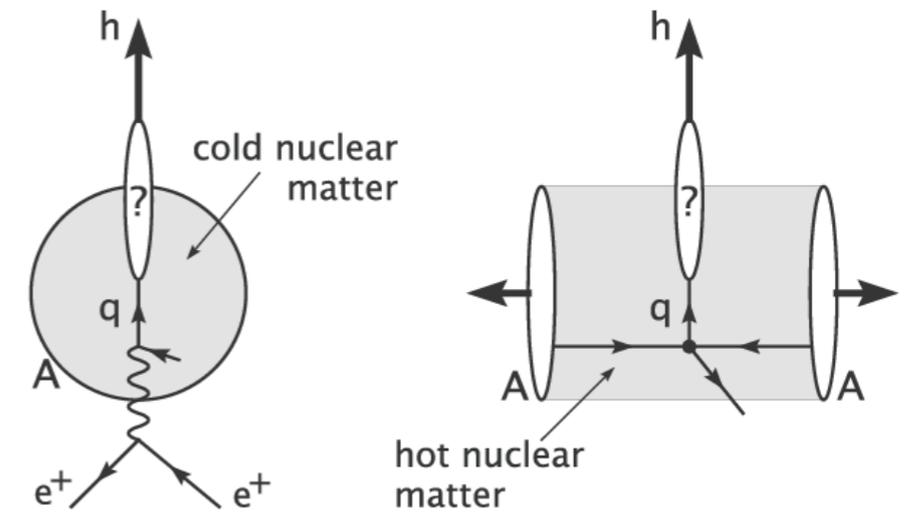
- Clean measurement in ‘cold’ nuclear matter
- Suppression of high- p_T hadrons analogous but *weaker* than at RHIC



Hadronization and Energy Loss

nDIS:

- Clean measurement in 'cold' nuclear matter
- Suppression of high- p_T hadrons analogous but *weaker* than at RHIC



Fundamental question:

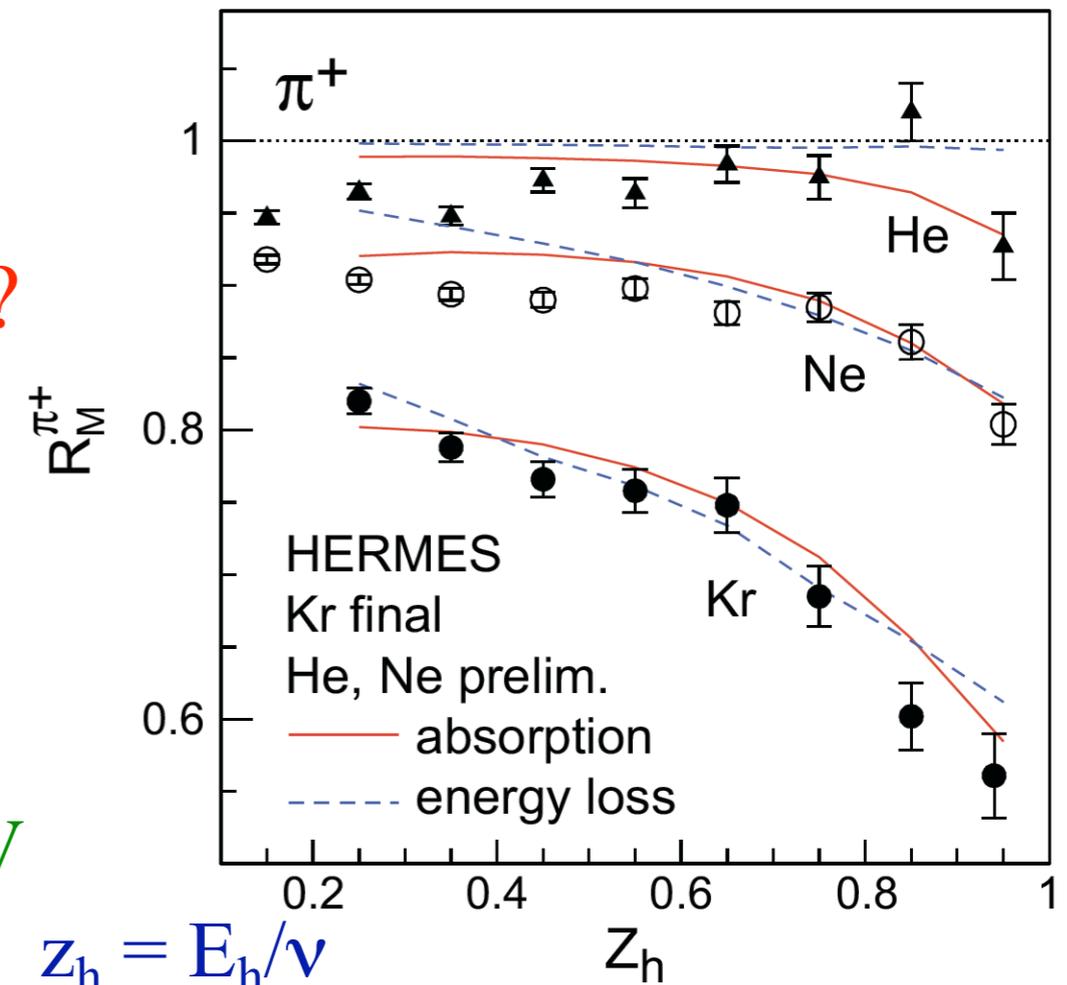
When do coloured partons get neutralized?

Parton energy loss vs.
(pre)hadron absorption

Energy transfer in lab rest frame

EIC: $10 < \nu < 1600$ GeV HERMES: 2-25 GeV

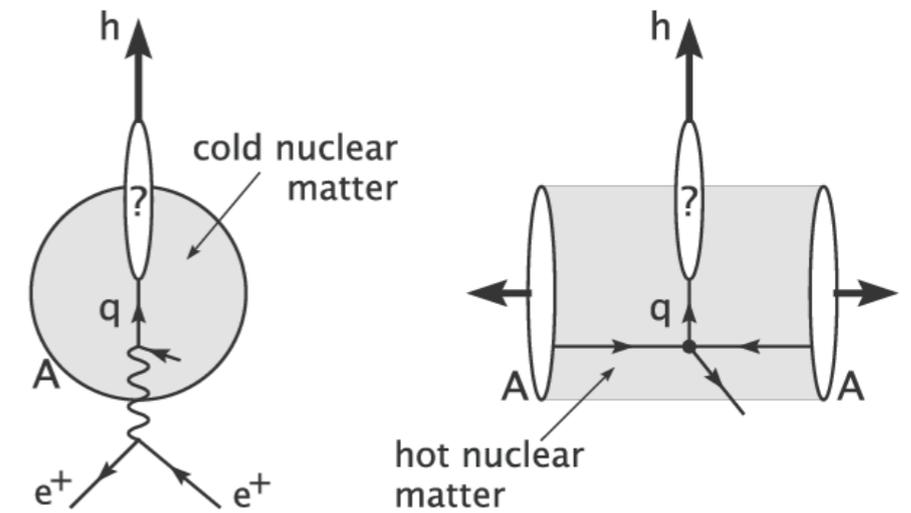
EIC: can measure *heavy flavour* energy loss $z_h = E_h/\nu$



Hadronization and Energy Loss

nDIS:

- Clean measurement in 'cold' nuclear matter
- Suppression of high- p_T hadrons analogous but *weaker* than at RHIC



Fundamental question:

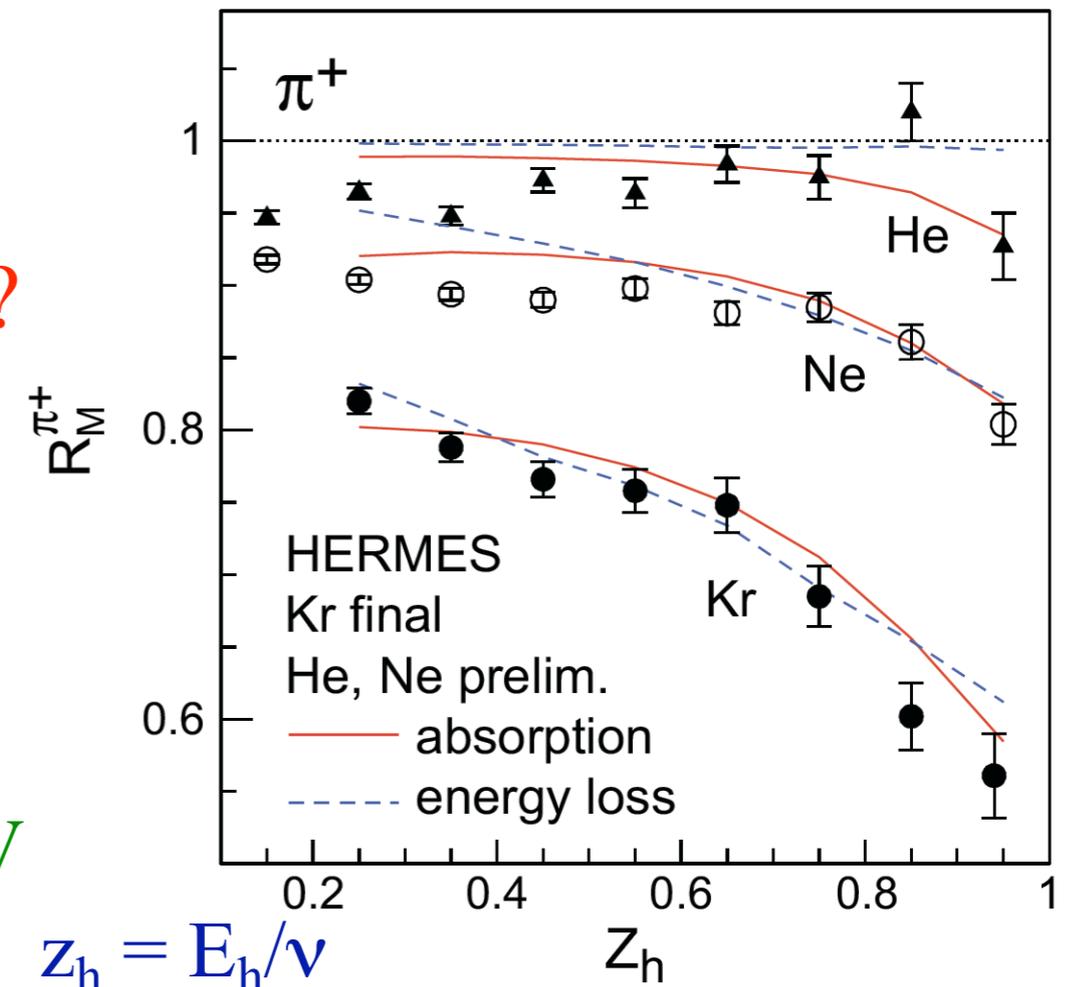
When do coloured partons get neutralized?

Parton energy loss vs. (pre)hadron absorption

Energy transfer in lab rest frame

EIC: $10 < \nu < 1600$ GeV HERMES: 2-25 GeV

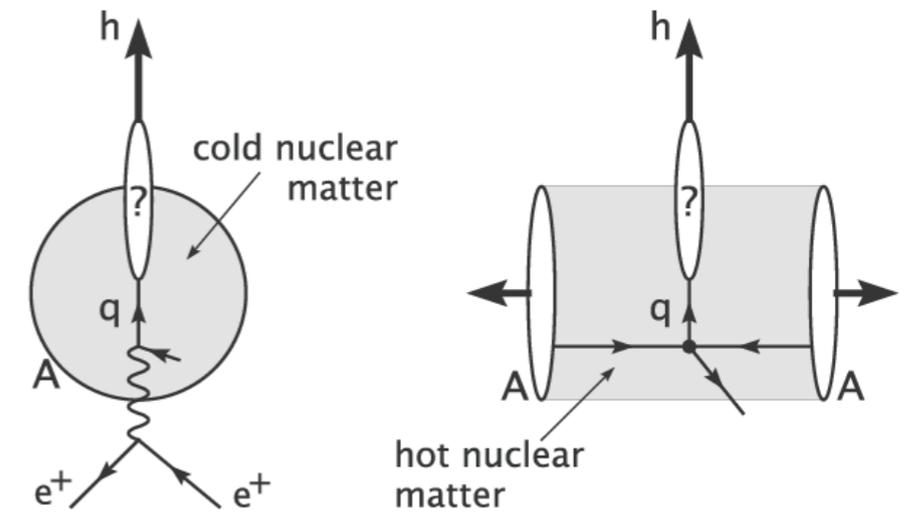
EIC: can measure *heavy flavour* energy loss $z_h = E_h/\nu$



Hadronization and Energy Loss

nDIS:

- Clean measurement in ‘cold’ nuclear matter
- Suppression of high- p_T hadrons analogous but *weaker* than at RHIC



Fundamental question:

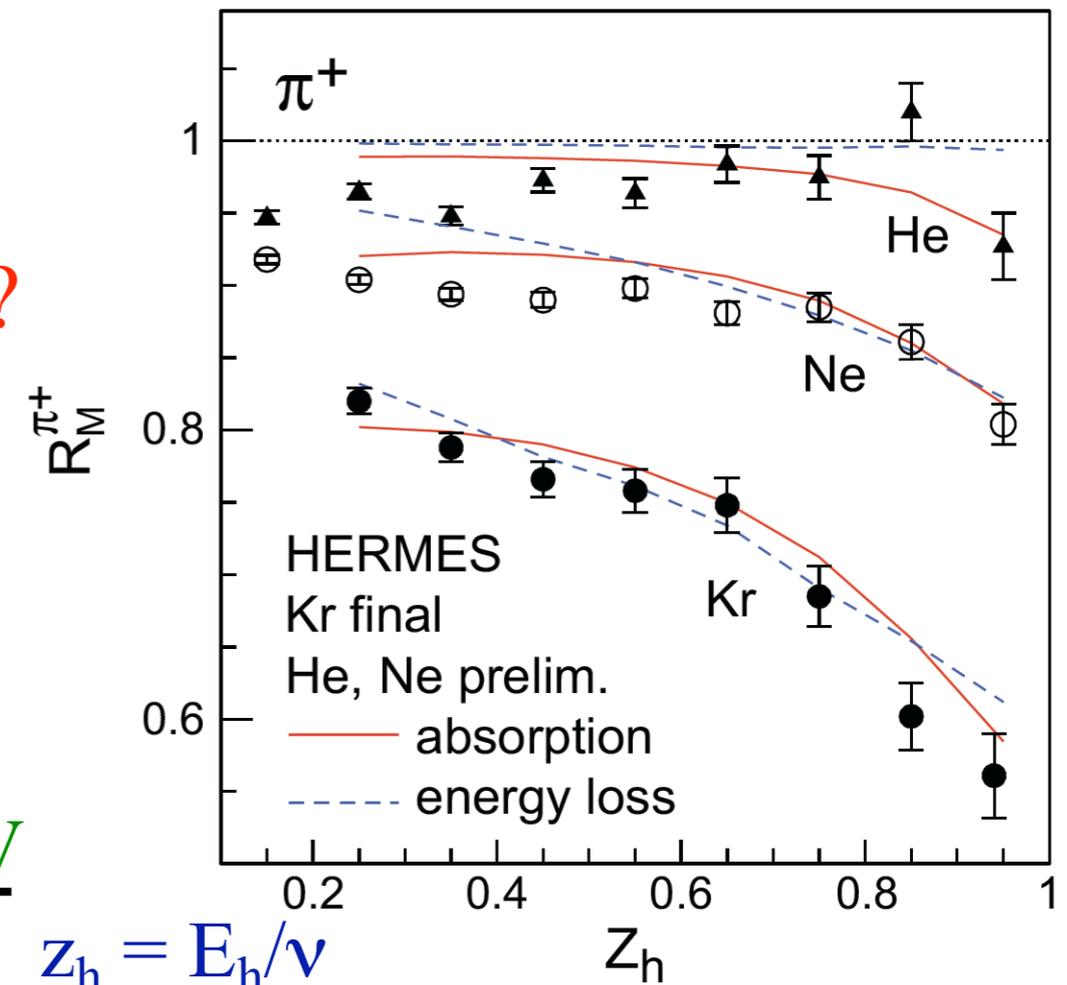
When do coloured partons get neutralized?

Parton energy loss vs. (pre)hadron absorption

Energy transfer in lab rest frame

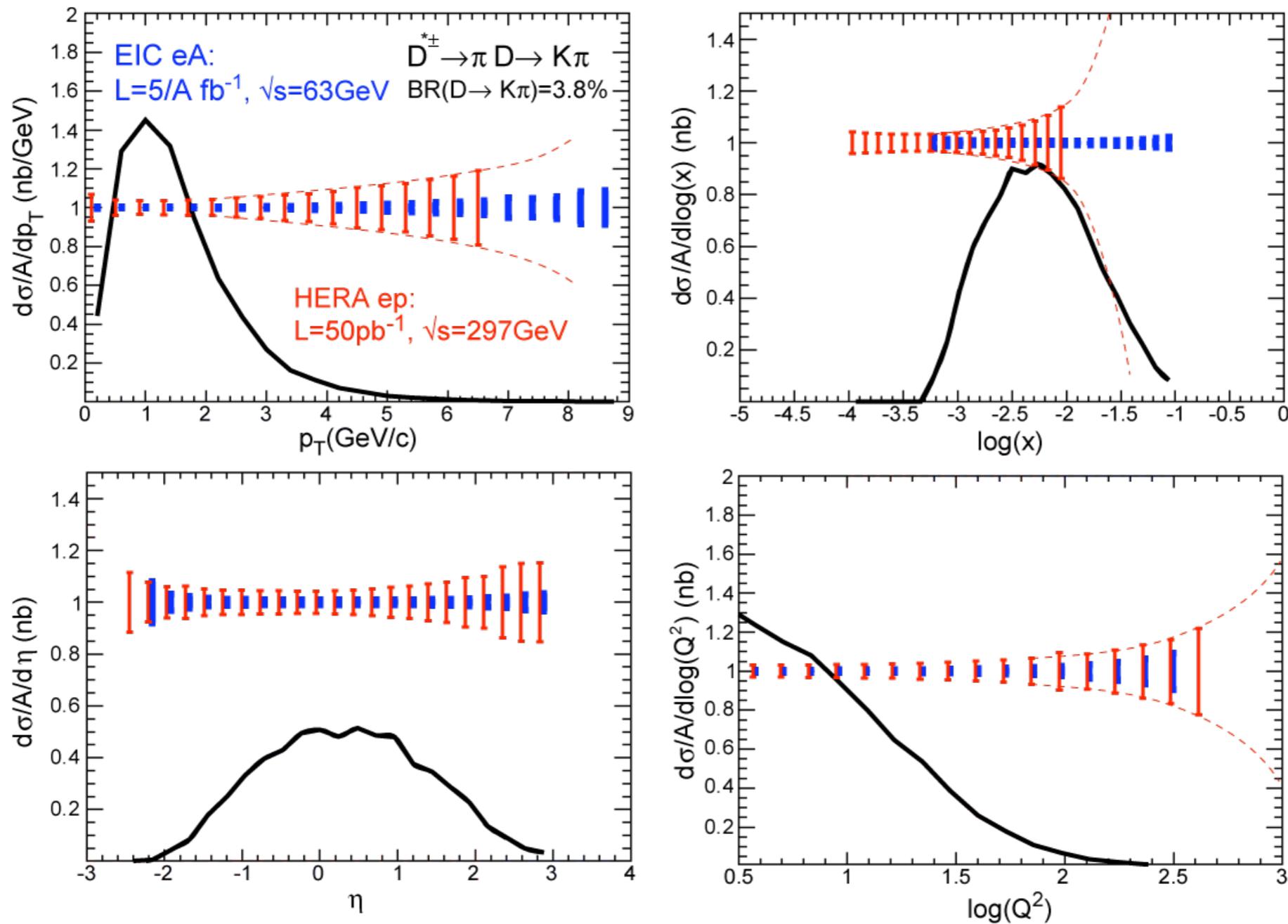
EIC: $10 < \nu < 1600$ GeV HERMES: 2-25 GeV

EIC: can measure *heavy flavour* energy loss $z_h = E_h/\nu$



Charm at an EIC

Based on HVQDIS model, J. Smith



- EIC: allows multi-differential measurements of heavy flavour
- covers and extends energy range of SLAC, EMC, HERA, and JLAB allowing for the study of wide range of formation lengths

Key Measurements in $e+A$

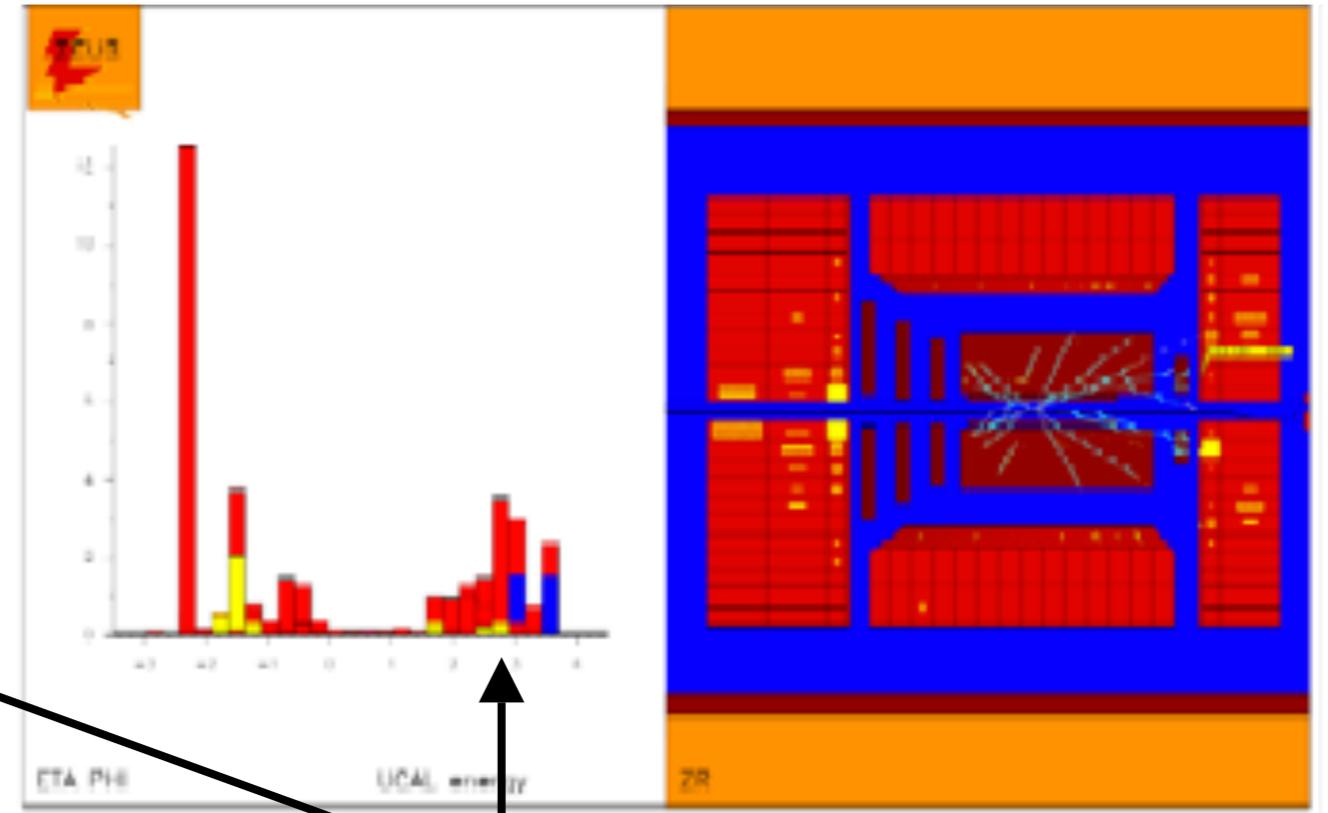
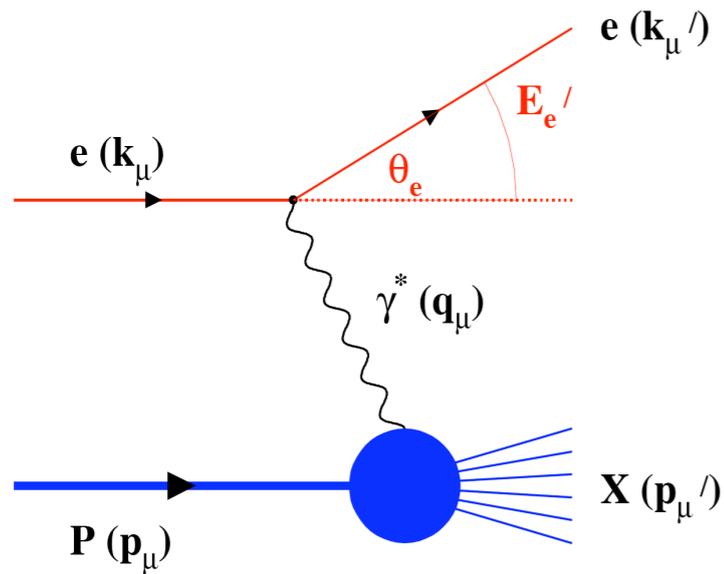
- **Momentum distribution of gluons $G(x, Q^2)$**
 - ➔ Extract via scaling violation in F_2 : $\delta F_2 / \delta \ln Q^2$
 - ➔ Direct measurement: $F_L \sim xG(x, Q^2)$ (requires \sqrt{s} scan)
 - ➔ 2+1 jet rates
 - ➔ Inelastic vector meson production (e.g. J/ψ)
 - ➔ Diffractive vector meson production $\sim [xG(x, Q^2)]^2$
- **Space-time distributions of gluons in matter**
 - ➔ Exclusive final states (e.g. vector meson production $\rho, J/\psi$)
 - ➔ Deep Virtual Compton Scattering (DVCS) - $\sigma \sim A^{4/3}$
 - ➔ F_2, F_L for various A and impact parameter dependence
- **Interaction of fast probes with *gluonic* medium?**
 - ➔ Hadronization, Fragmentation
 - ➔ Energy loss (charm!)
- **Role of colour neutral excitations (Pomerons)**
 - ➔ Diffractive cross-section $\sigma_{diff}/\sigma_{tot}$ (HERA/ ep : 10% , EIC/ eA : 30%?)
 - ➔ Diffractive structure functions and vector meson production
 - ➔ Abundance and distribution of rapidity gaps

Key Measurements in $e+A$

- **Momentum distribution of gluons $G(x, Q^2)$**
 - ➔ Extract via scaling violation in F_2 : $\delta F_2 / \delta \ln Q^2$
 - ➔ Direct measurement: $F_L \sim xG(x, Q^2)$ (requires \sqrt{s} scan)
 - ➔ 2+1 jet rates
 - ➔ Inelastic vector meson production (e.g. J/ψ)
 - ➔ Diffractive vector meson production $\sim [xG(x, Q^2)]^2$
- **Space-time distributions of gluons in matter**
 - ➔ Exclusive final states (e.g. vector meson production $\rho, J/\psi$)
 - ➔ Deep Virtual Compton Scattering (DVCS) - $\sigma \sim A^{4/3}$
 - ➔ F_2, F_L for various A and impact parameter dependence
- **Interaction of fast probes with *gluonic* medium?**
 - ➔ Hadronization, Fragmentation
 - ➔ Energy loss (charm!)
- **Role of colour neutral excitations (Pomerons)**
 - ➔ Diffractive cross-section $\sigma_{diff}/\sigma_{tot}$ (HERA/ ep : 10% , EIC/ eA : 30%?)
 - ➔ Diffractive structure functions and vector meson production
 - ➔ Abundance and distribution of rapidity gaps

Diffractive Physics in $e+A$

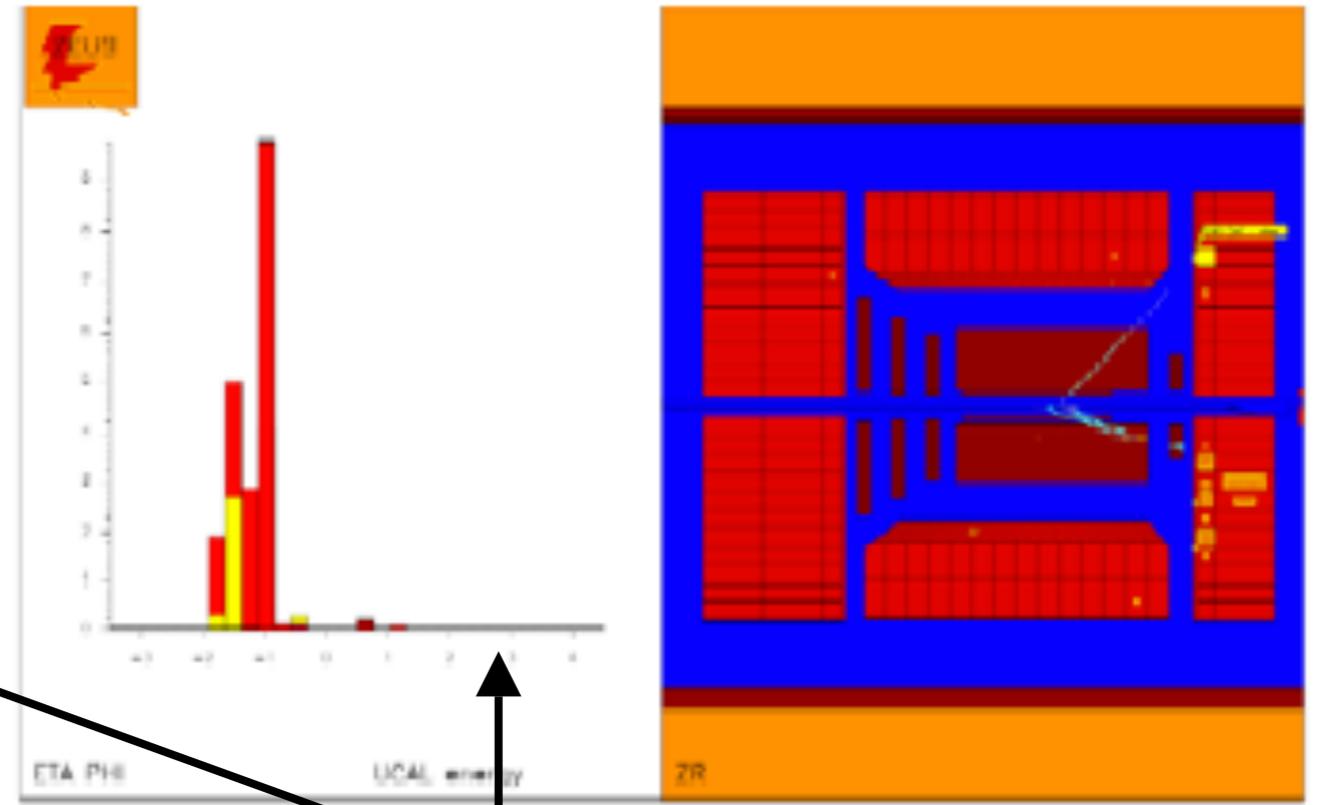
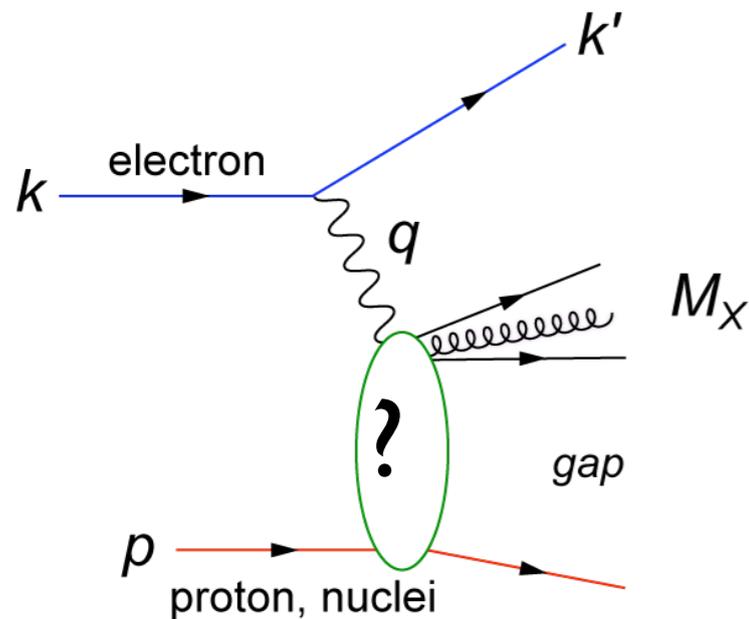
‘Standard DIS event’



Activity in proton direction

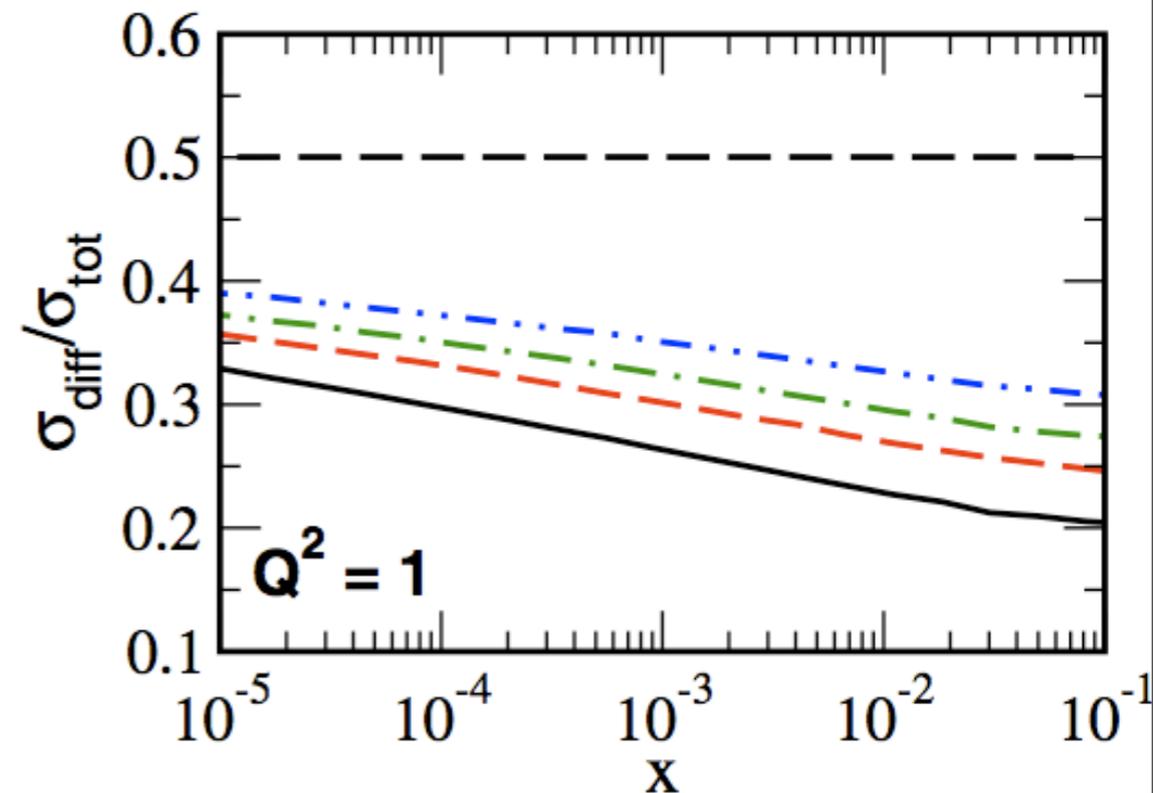
Diffractive Physics in $e+A$

Diffractive event



- HERA/ep: 15% of all events are hard diffractive
- Diffractive cross-section $\sigma_{\text{diff}}/\sigma_{\text{tot}}$ in $e+A$?
- ➔ Predictions: ~25-40%?

Activity in proton direction

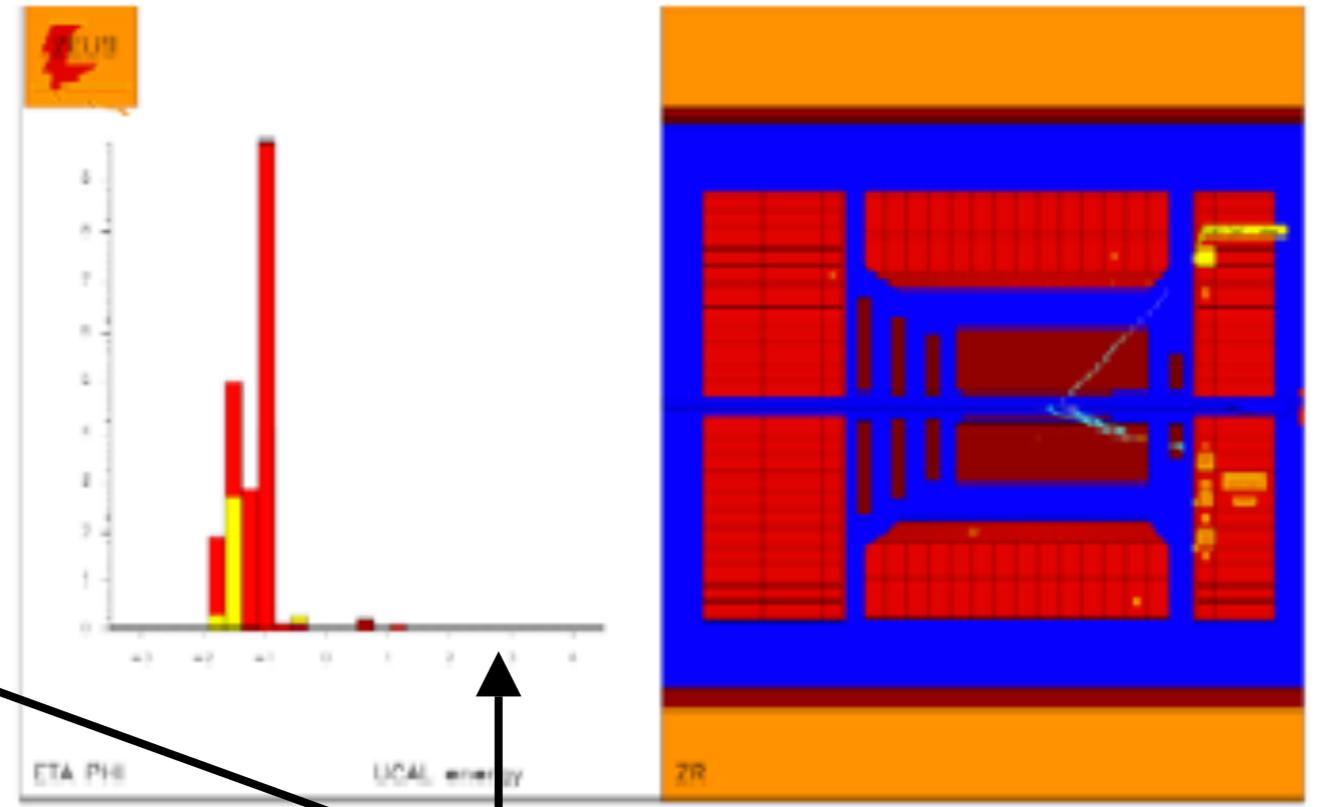
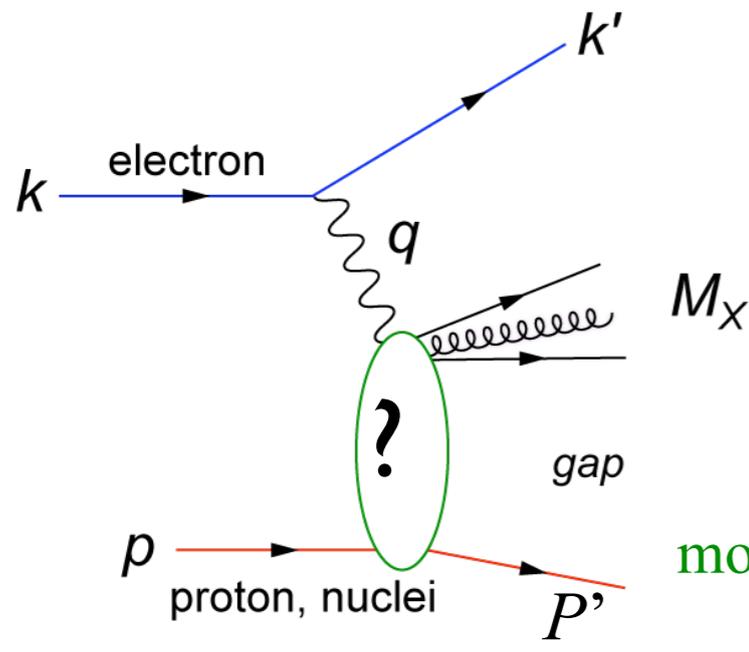


Curves: Kugeratski, Goncalves,
Navarra, EPJ C46, 413

HQ2008: macl@bnl.gc

Diffraction Physics in $e+A$

Diffractive event



- **HERA/ep**: 15% of all events are hard diffractive
- Diffractive cross-section $\sigma_{\text{diff}}/\sigma_{\text{tot}}$ in $e+A$?

➔ Predictions: $\sim 25-40\%$?

- Look inside the “Pomeron”

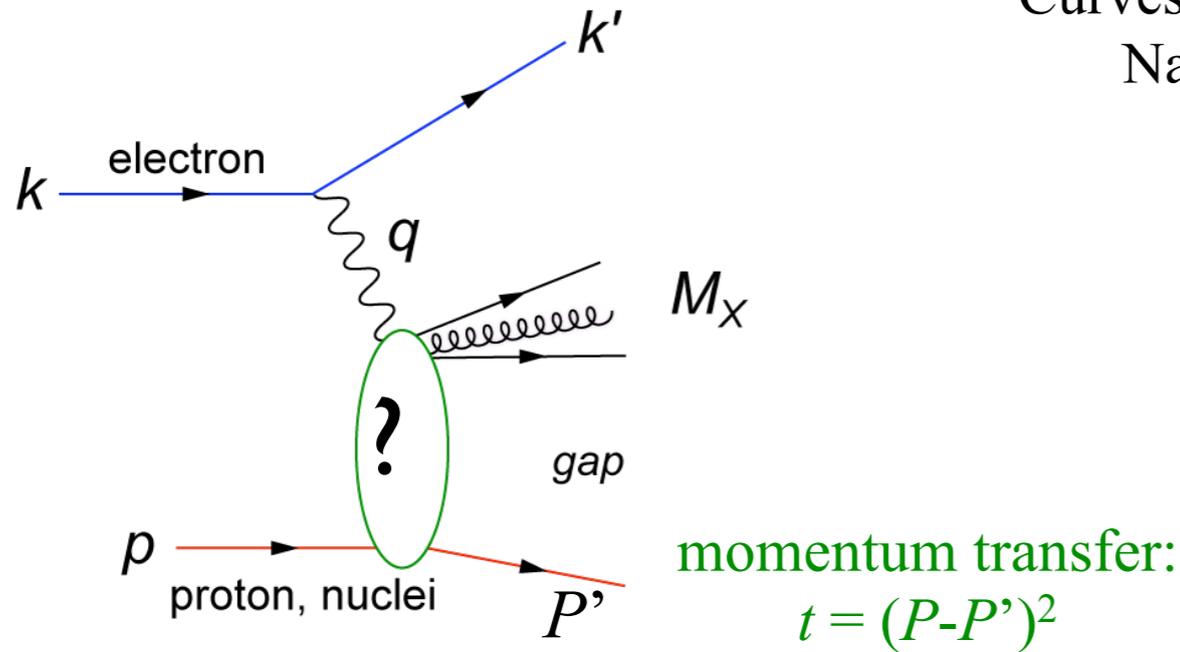
➔ Diffractive structure functions

➔ Diffractive vector meson production: $d\sigma/dt \sim [xG(x, Q^2)]^2$!!

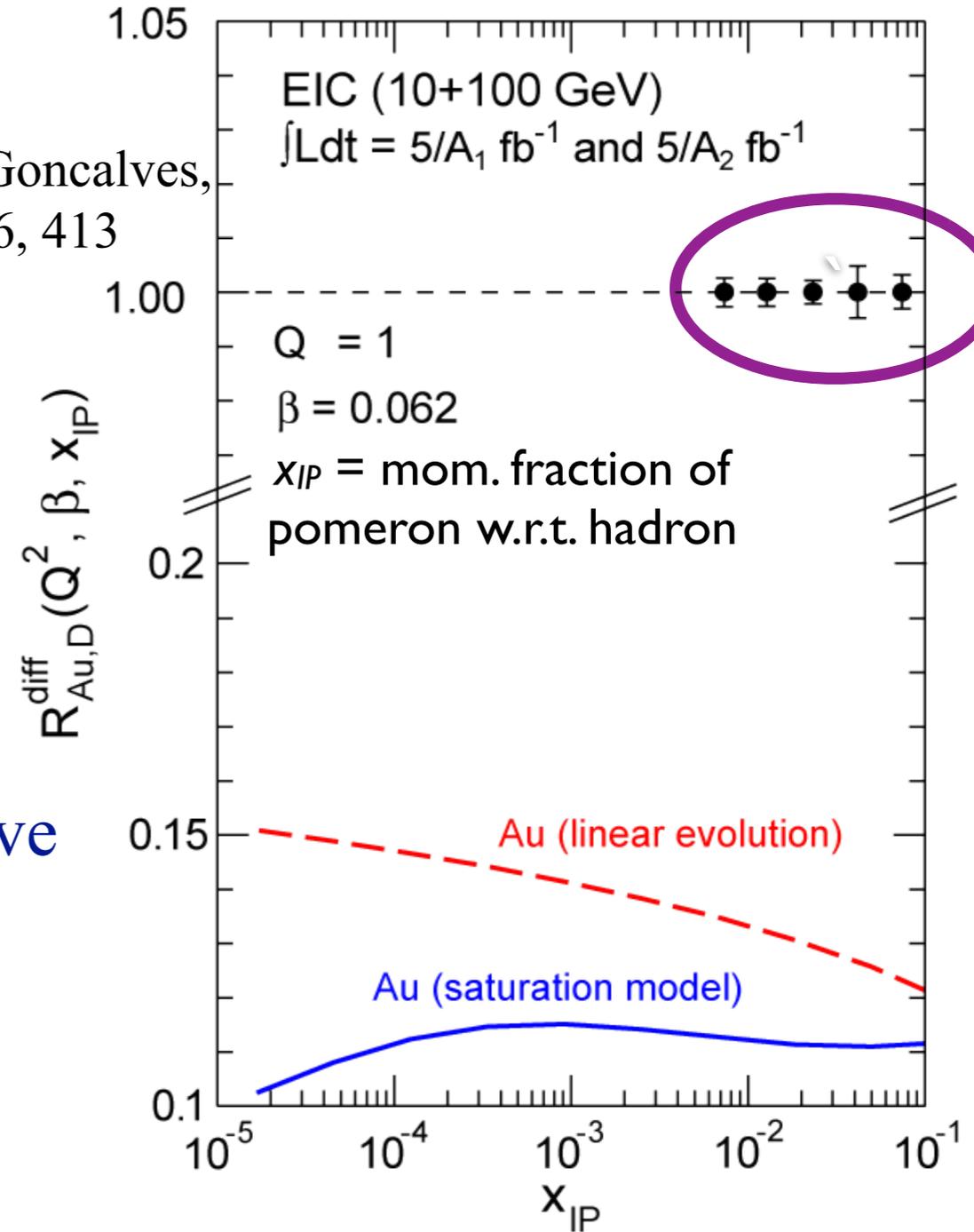
Activity in proton direction

Diffraction Physics in e+A

Diffractive event



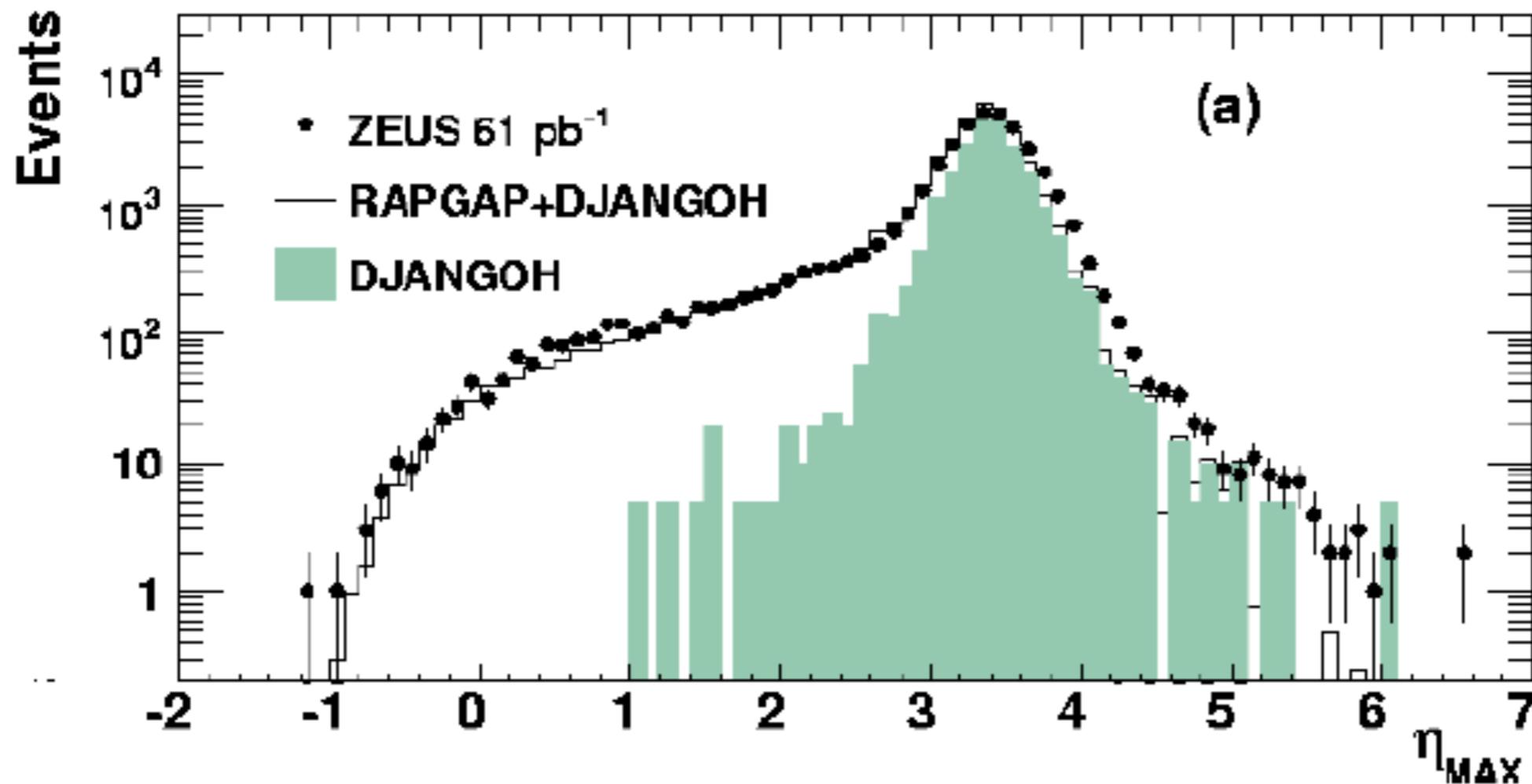
- HERA/ep: 15% of all events are hard diffractive
- Diffractive cross-section $\sigma_{\text{diff}}/\sigma_{\text{tot}}$ in e+A ?
 - ➔ Predictions: ~25-40%?
- Look inside the “Pomeron”
 - ➔ Diffractive structure functions
 - ➔ Diffractive vector meson production: $d\sigma/dt \sim [xG(x, Q^2)]^2 !!$
- Distinguish between linear evolutions and saturation models



Diffraction Physics in e+A

- How to measure diffraction in e+A?
 - ➔ Use HERA method of **Large Rapidity Gaps**
 - ➔ Ideal gap of ~ 7.7 at HERA units reduced to 3-4 due to **spread from hadronisation**

ZEUS



Diffraction Physics in e+A

- How to measure diffraction in e+A?

- ➔ Use HERA method of **Large Rapidity Gaps**

- ➔ Ideal gap of ~ 7.7 at HERA units reduced to 3-4 due to **spread from hadronisation**

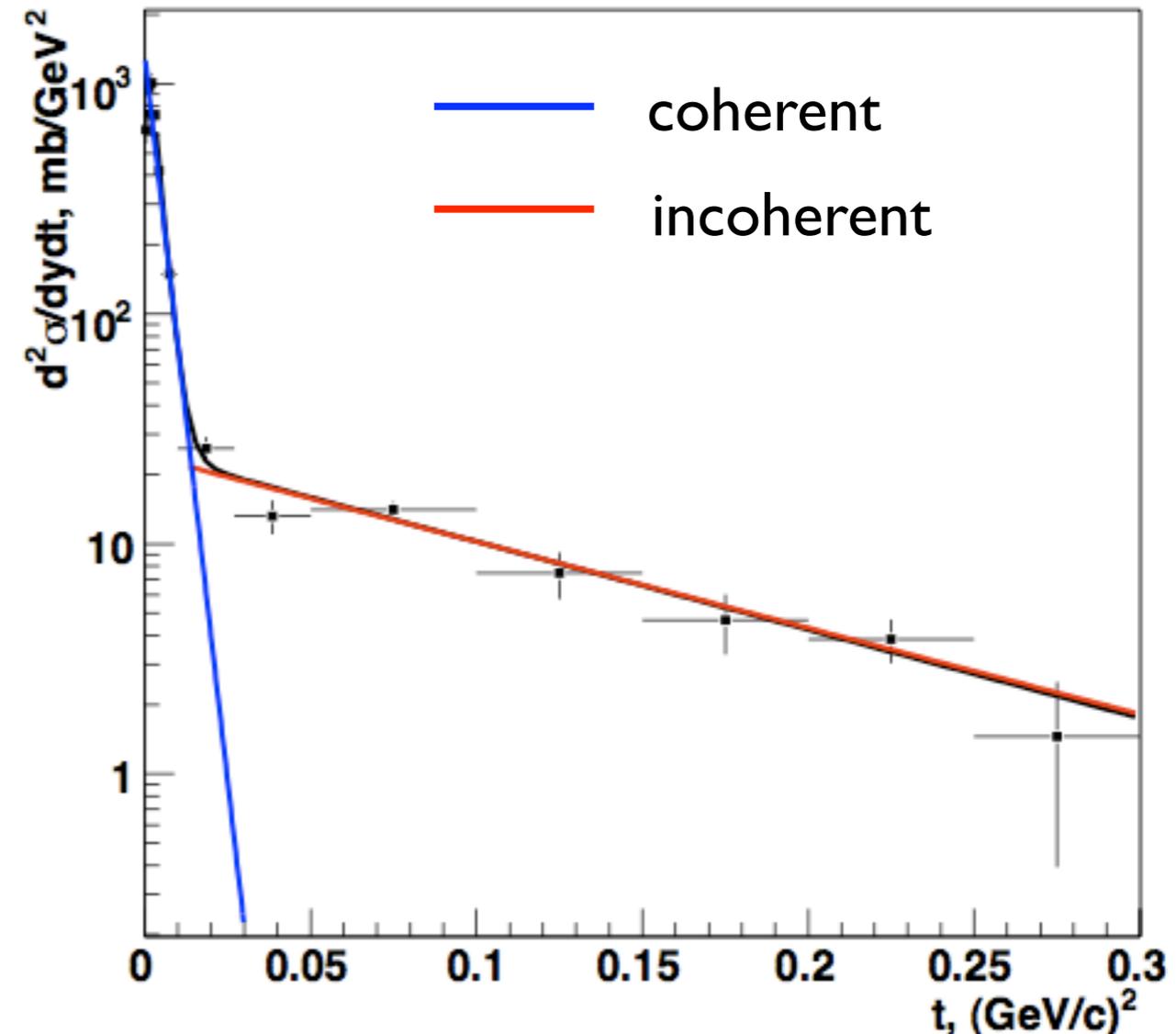
- Issues with measuring diffractive physics in e+A:

- ➔ t required for nucleus to break-up is small ($\sim 30 \text{ MeV}/c^2$)

- ➔ t required for nucleus to be measured in detector $\gg 30 \text{ MeV}/c^2$

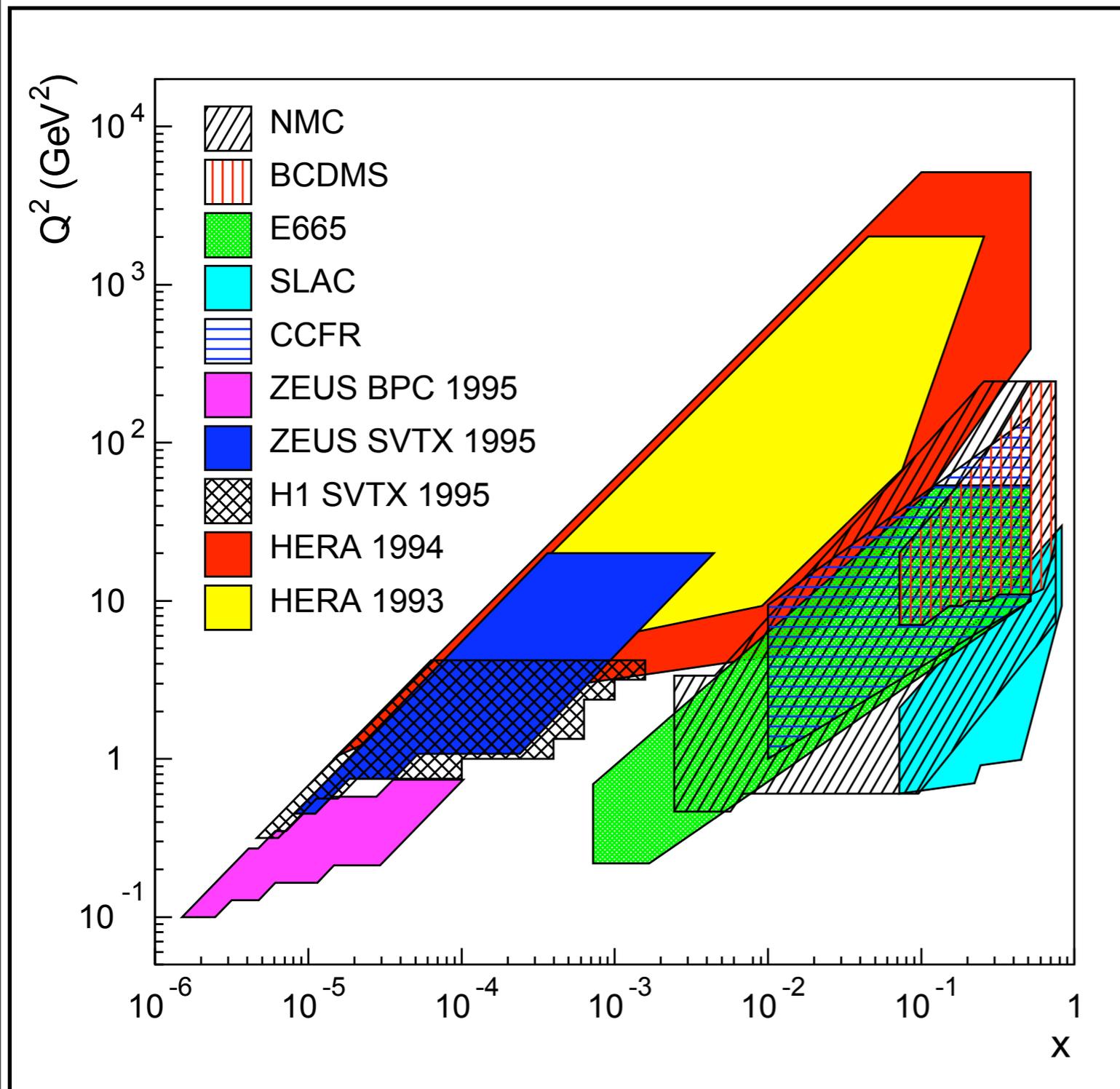
- ➔ To measure t dependence, must measure exclusive diffraction (e.g. **vector mesons** - $t \sim p_T^2$)

STAR - UPC Collisions

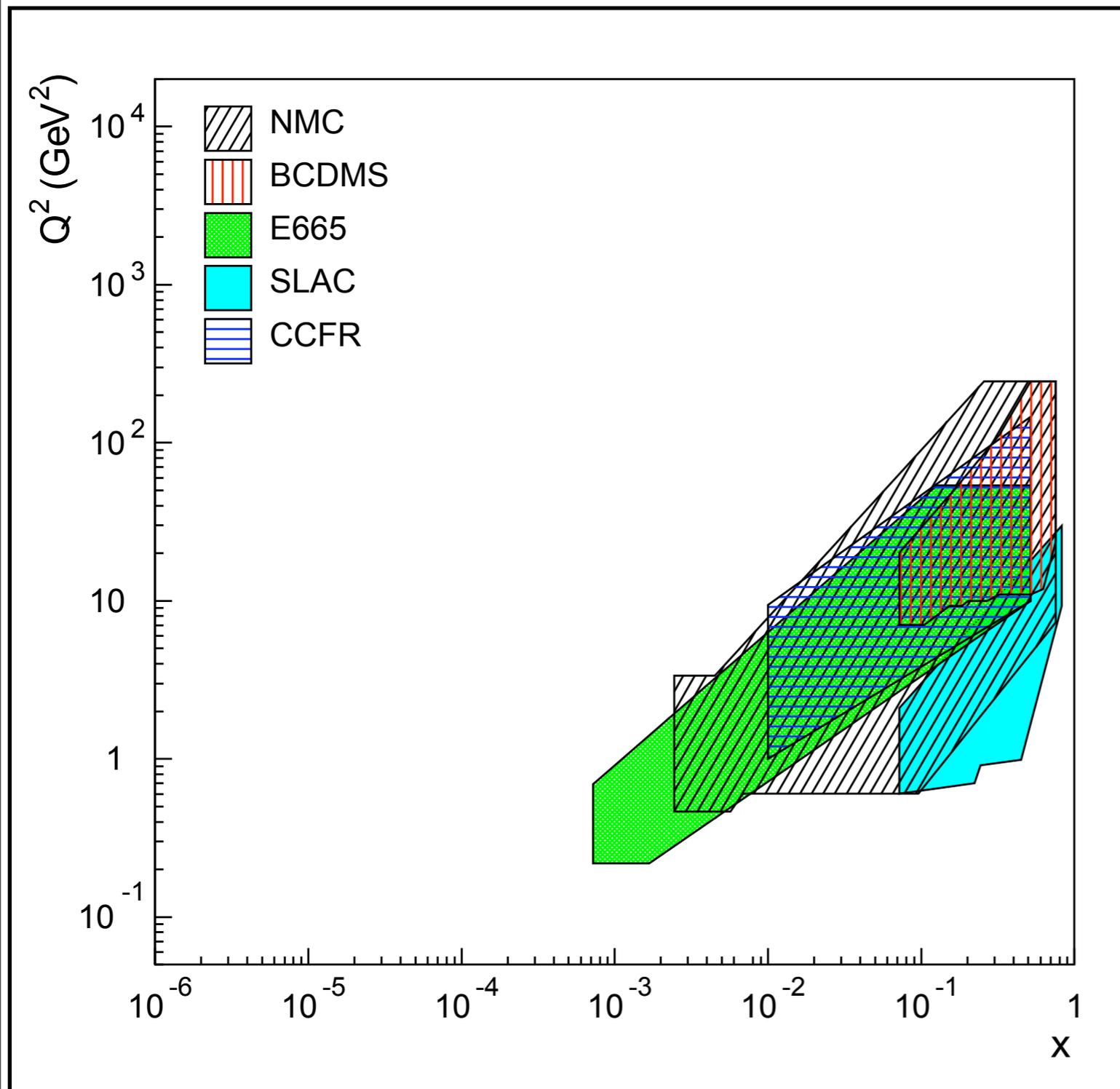


Requirements for an Electron Ion Collider

Well mapped in $e+p$



Requirements for an Electron Ion Collider

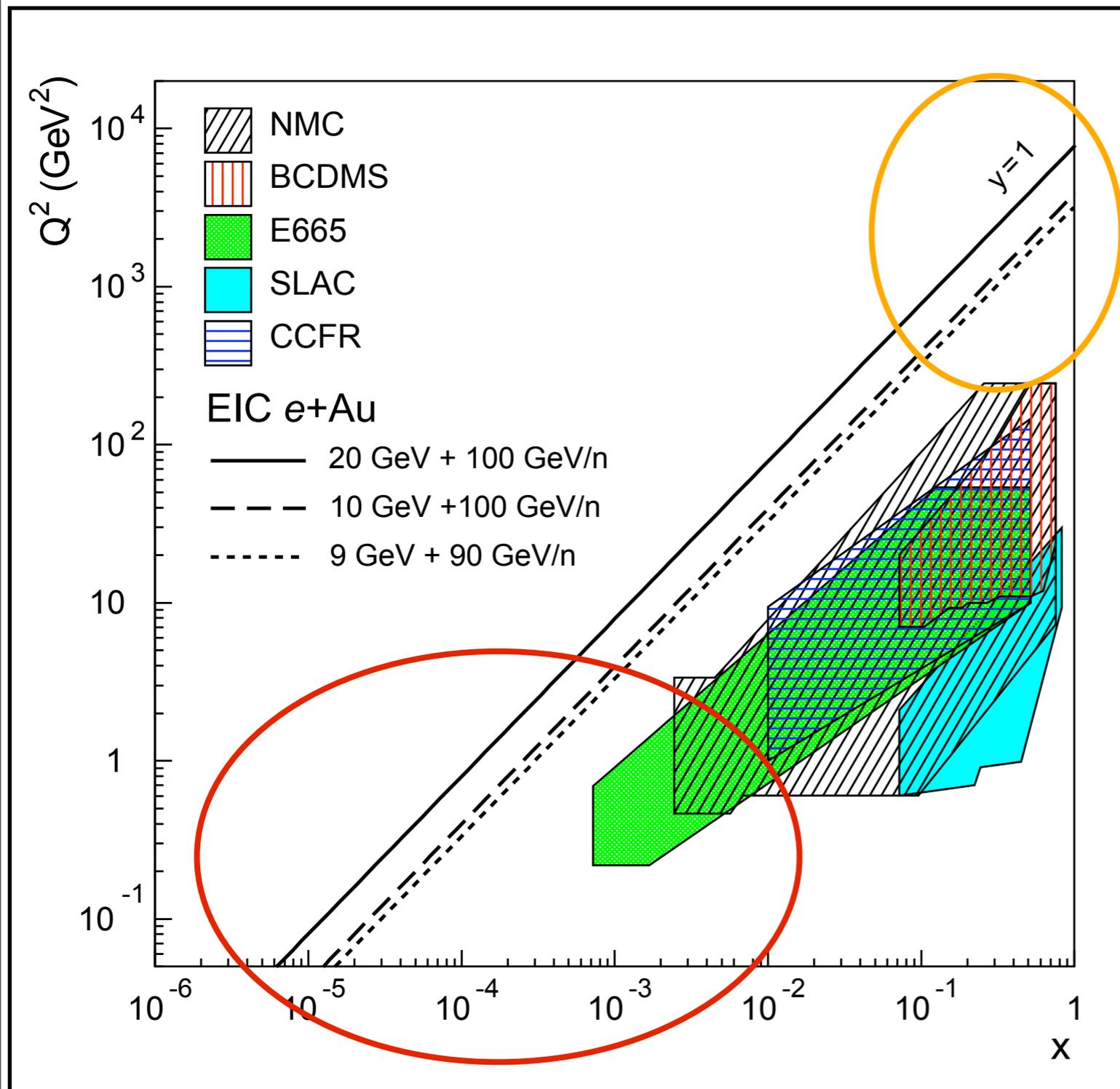


Well mapped in $e+p$

Not so for $\ell+A$ ($\nu+A$)

- many with small A
- low statistics

Requirements for an Electron Ion Collider



Well mapped in $e+p$

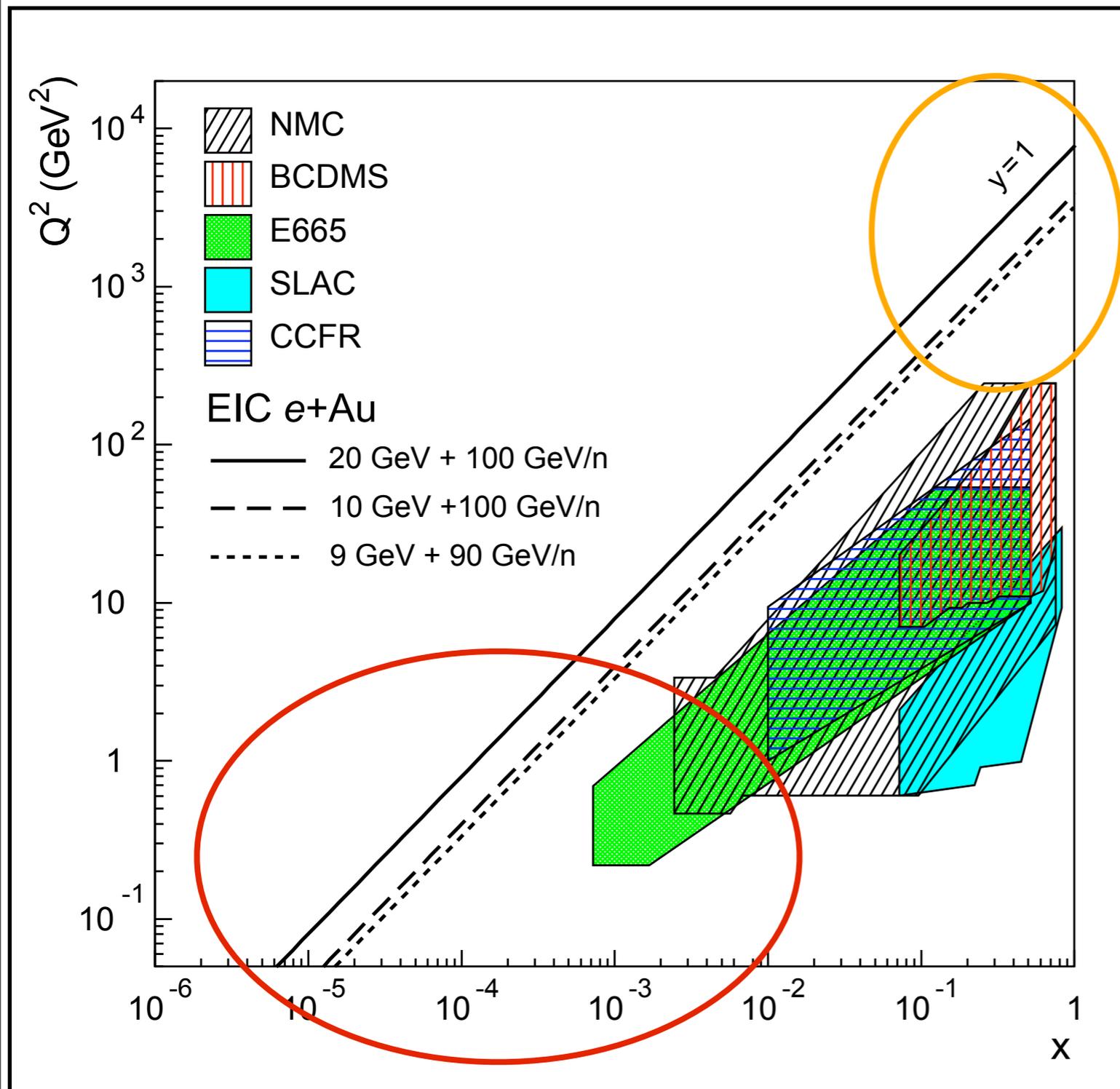
Not so for $\ell+A$ ($\nu+A$)

- many with small A
- low statistics

Electron Ion Collider:

- $\mathcal{L}(\text{EIC}) > 100 \times \mathcal{L}(\text{HERA})$
- Electrons
 - $E_e = 3 - 20$ GeV
 - polarized
- Hadron Beams
 - $E_A = 100$ GeV
 - $A = p \rightarrow U$
 - polarized p & light ions

Requirements for an Electron Ion Collider



Well mapped in $e+p$

Not so for $\ell+A$ ($\nu+A$)

- many with small A
- low statistics

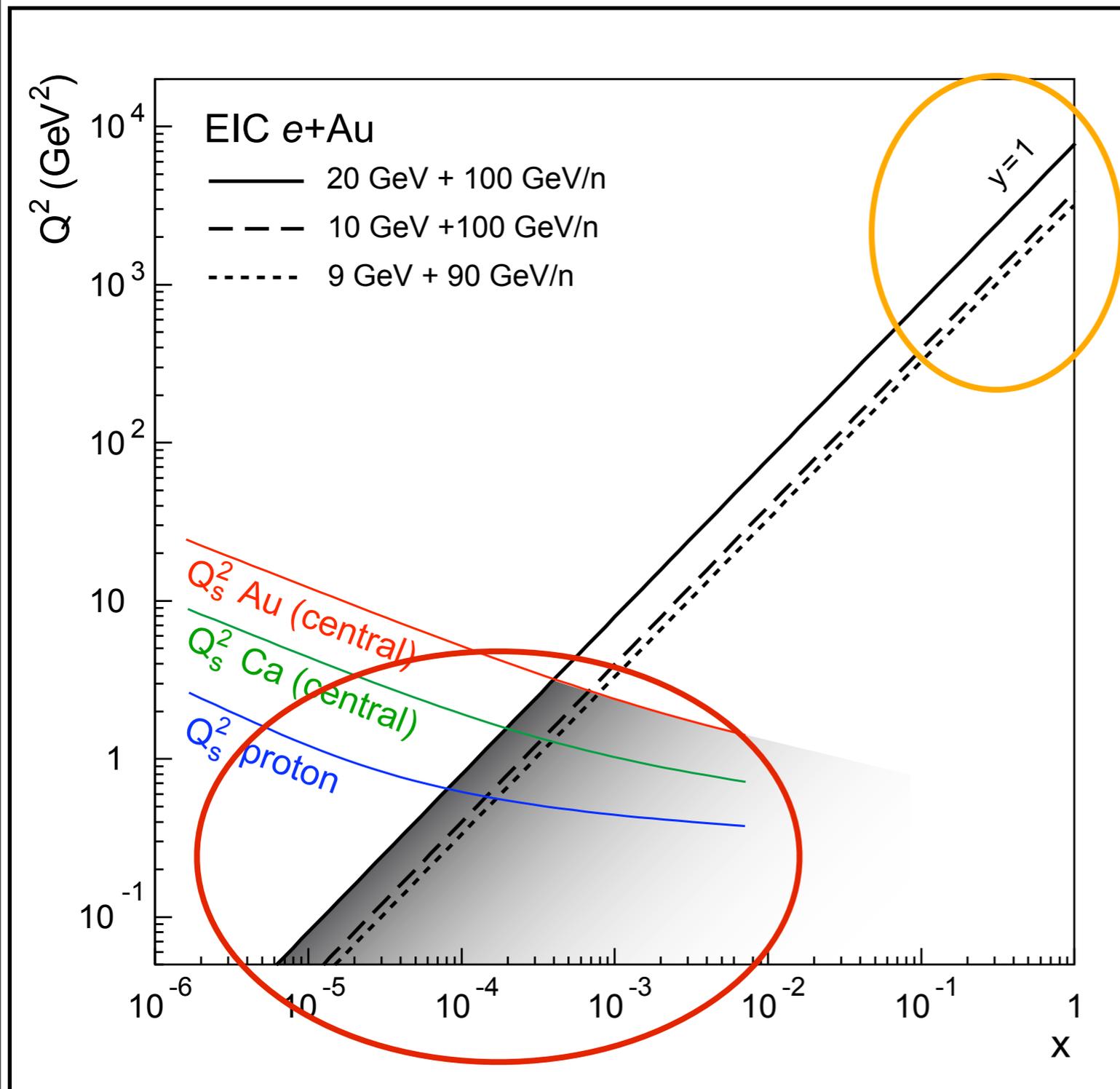
Electron Ion Collider:

- $\mathcal{L}(\text{EIC}) > 100 \times \mathcal{L}(\text{HERA})$
- Electrons
 - $E_e = 3 - 20$ GeV
 - polarized
- Hadron Beams
 - $E_A = 100$ GeV
 - $A = p \rightarrow U$
 - polarized p & light ions

Terra incognita:

small- x , $Q \leq Q_s$
 high- x , large Q^2

Requirements for an Electron Ion Collider



Well mapped in $e+p$

Not so for $\ell+A$ ($\nu+A$)

- many with small A
- low statistics

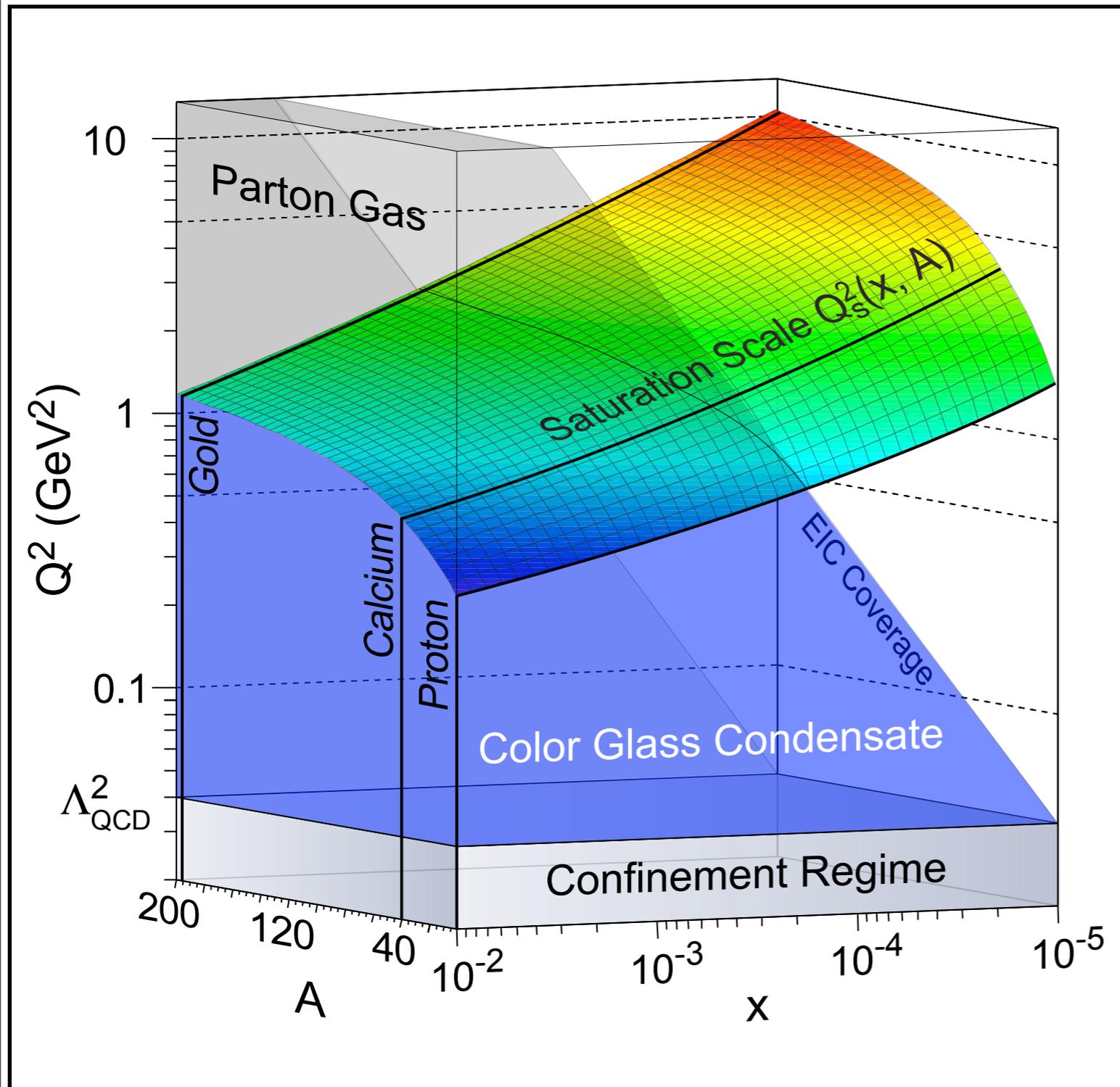
Electron Ion Collider:

- $\mathcal{L}(\text{EIC}) > 100 \times \mathcal{L}(\text{HERA})$
- Electrons
 - $E_e = 3 - 20$ GeV
 - polarized
- Hadron Beams
 - $E_A = 100$ GeV
 - $A = p \rightarrow U$
 - polarized p & light ions

Terra incognita:

small- x , $Q \leq Q_s$
high- x , large Q^2

Requirements for an Electron Ion Collider



Well mapped in $e+p$

Not so for $\ell+A$ ($\nu+A$)

- many with small A
- low statistics

Electron Ion Collider:

- $\mathcal{L}(\text{EIC}) > 100 \times \mathcal{L}(\text{HERA})$
- Electrons
 - $E_e = 3 - 20 \text{ GeV}$
 - polarized
- Hadron Beams
 - $E_A = 100 \text{ GeV}$
 - $A = p \rightarrow U$
 - polarized p & light ions

Terra incognita:

small- x , $Q \leq Q_s$
high- x , large Q^2

EIC Collider concepts

*e*RHIC (RHIC/BNL):

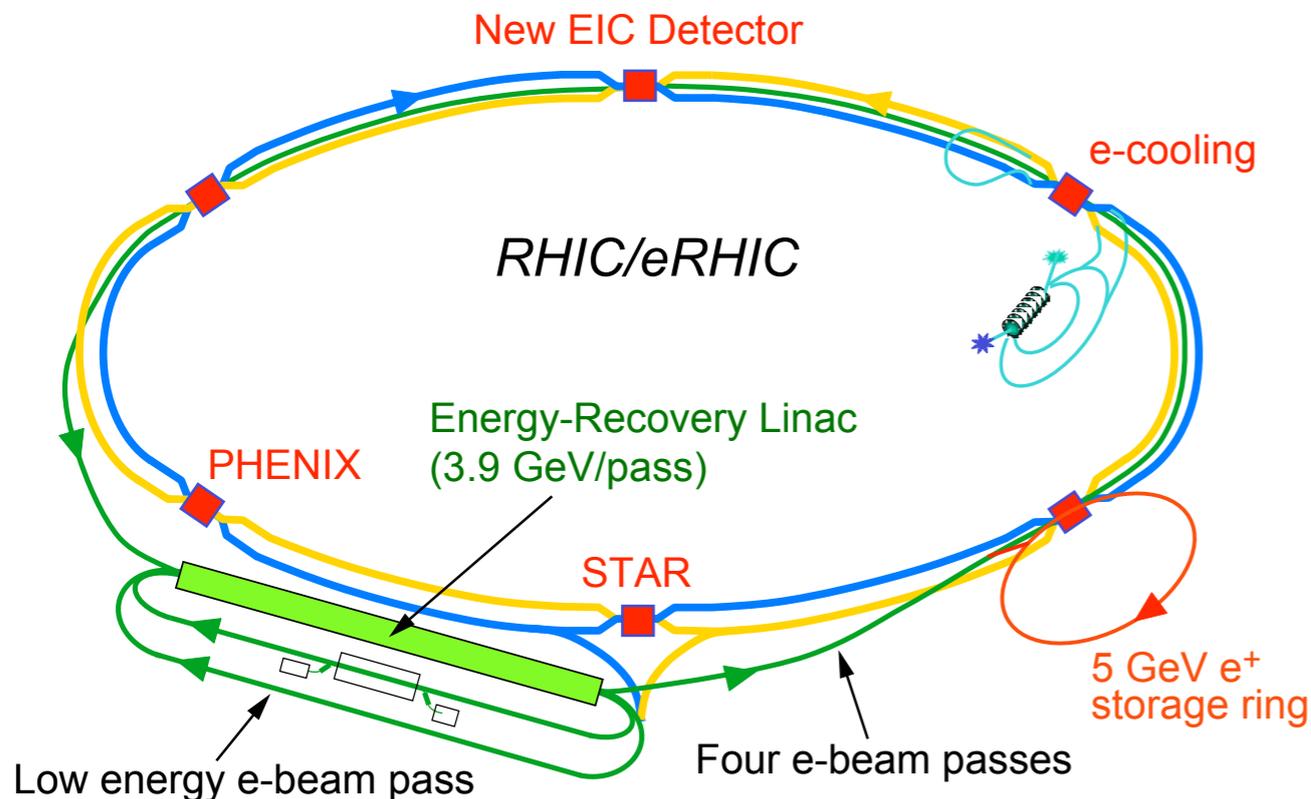
Add Energy Recovery Linac

$$E_e = 10 \text{ (20) GeV}$$

$$E_A = 100 \text{ GeV (up to U)}$$

$$\sqrt{s_{eN}} = 63 \text{ (90) GeV}$$

$$L_{eAu} \text{ (peak)}/n \sim 2.9 \cdot 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$$



EIC Collider concepts

eRHIC (RHIC/BNL):

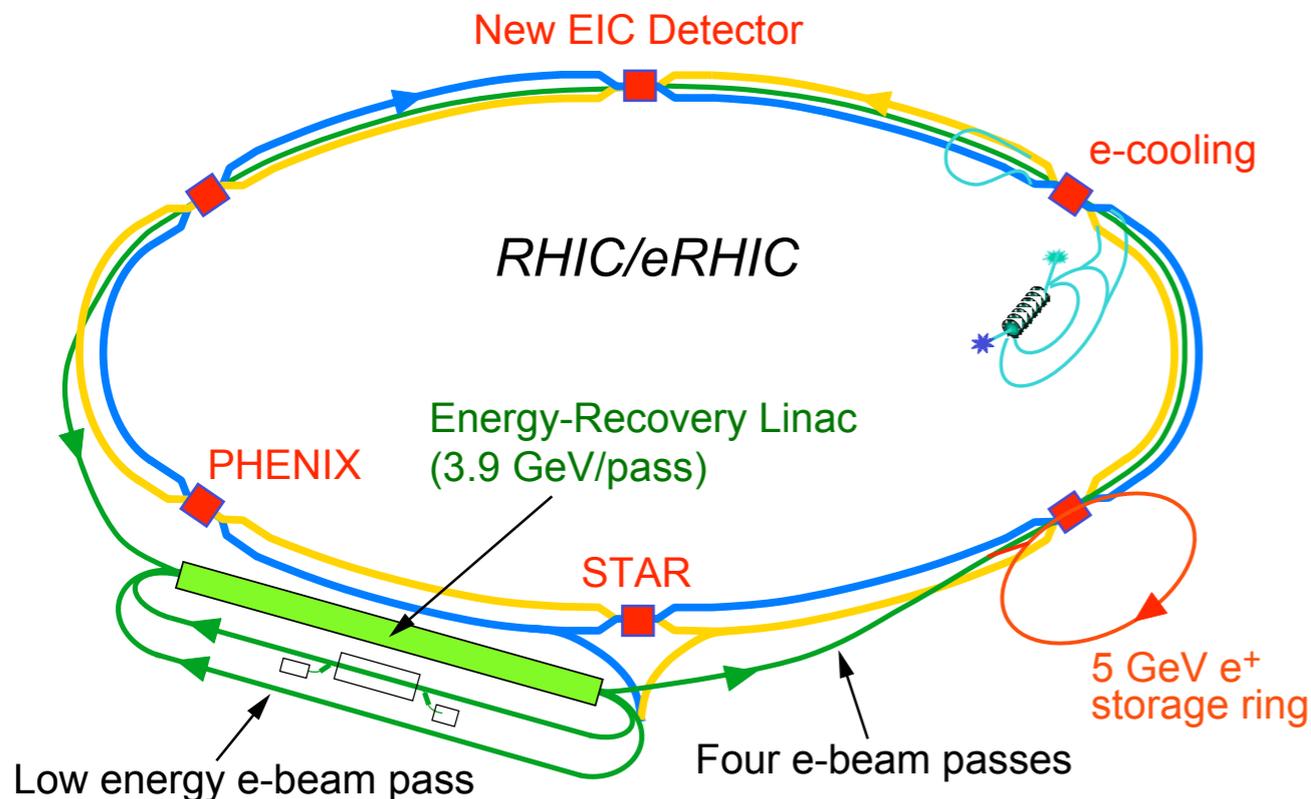
Add Energy Recovery Linac

$$E_e = 10 \text{ (20) GeV}$$

$$E_A = 100 \text{ GeV (up to U)}$$

$$\sqrt{s_{eN}} = 63 \text{ (90) GeV}$$

$$L_{eAu} \text{ (peak)}/n \sim 2.9 \cdot 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$$



ELIC (CEBAF/JLAB):

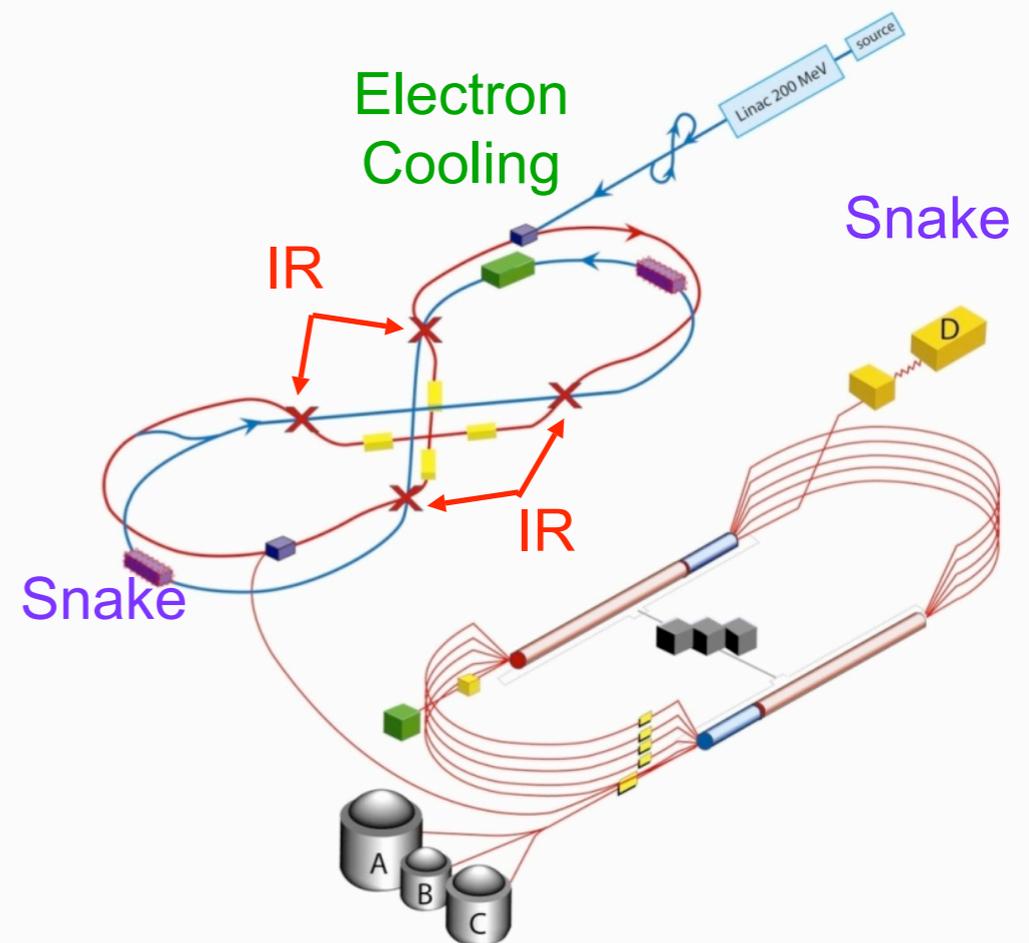
Add hadron machine

$$E_e = 9 \text{ GeV}$$

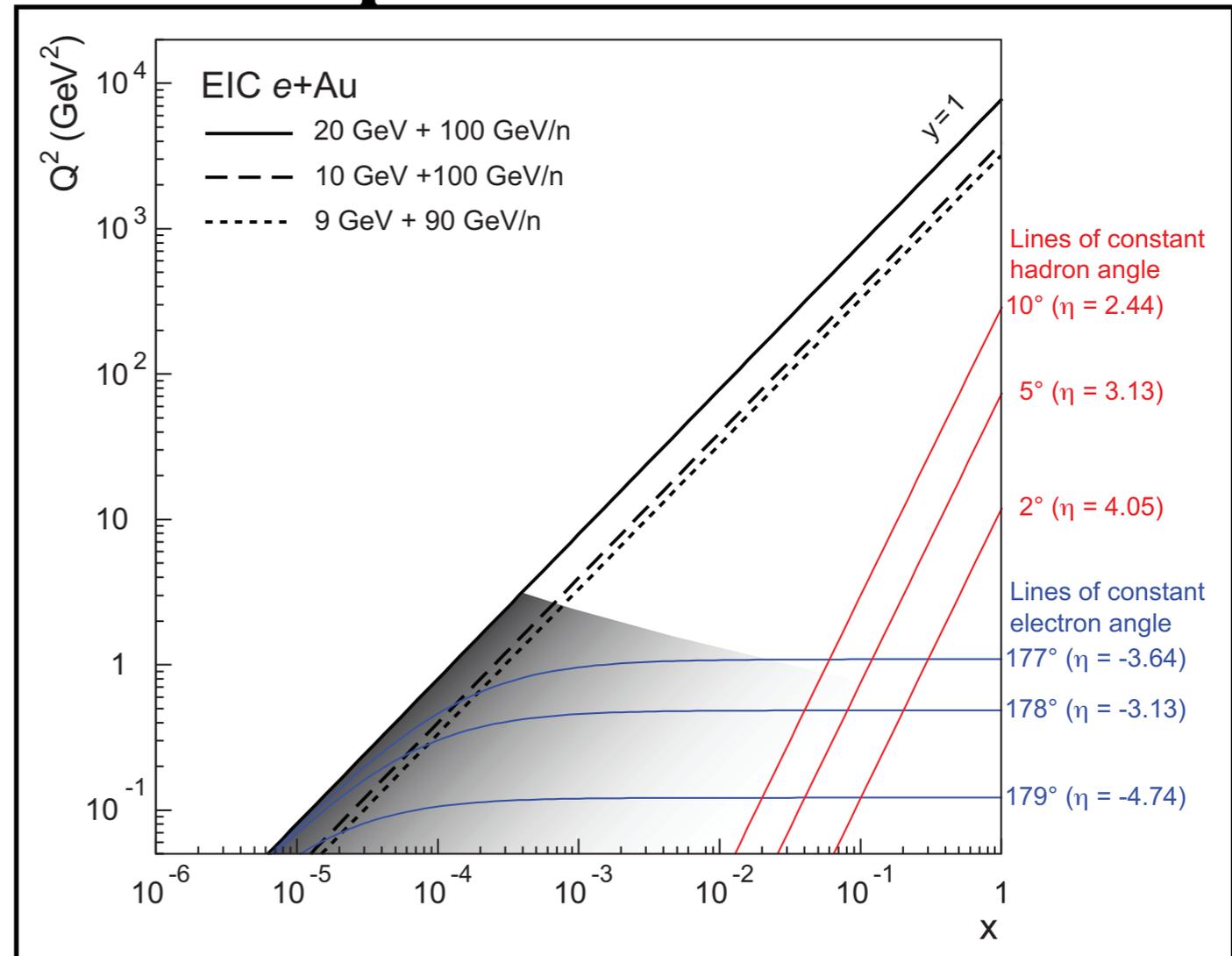
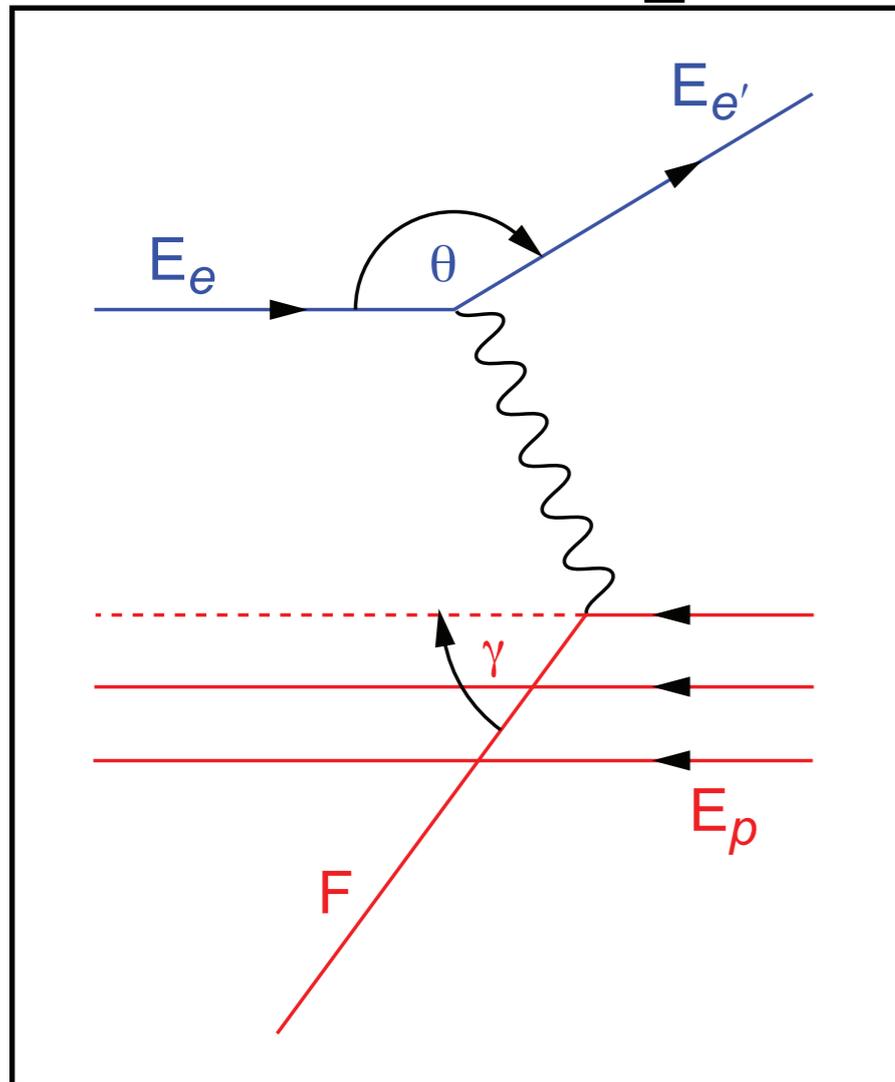
$$E_A = 90 \text{ GeV (up to Au)}$$

$$\sqrt{s_{eN}} = 57 \text{ GeV}$$

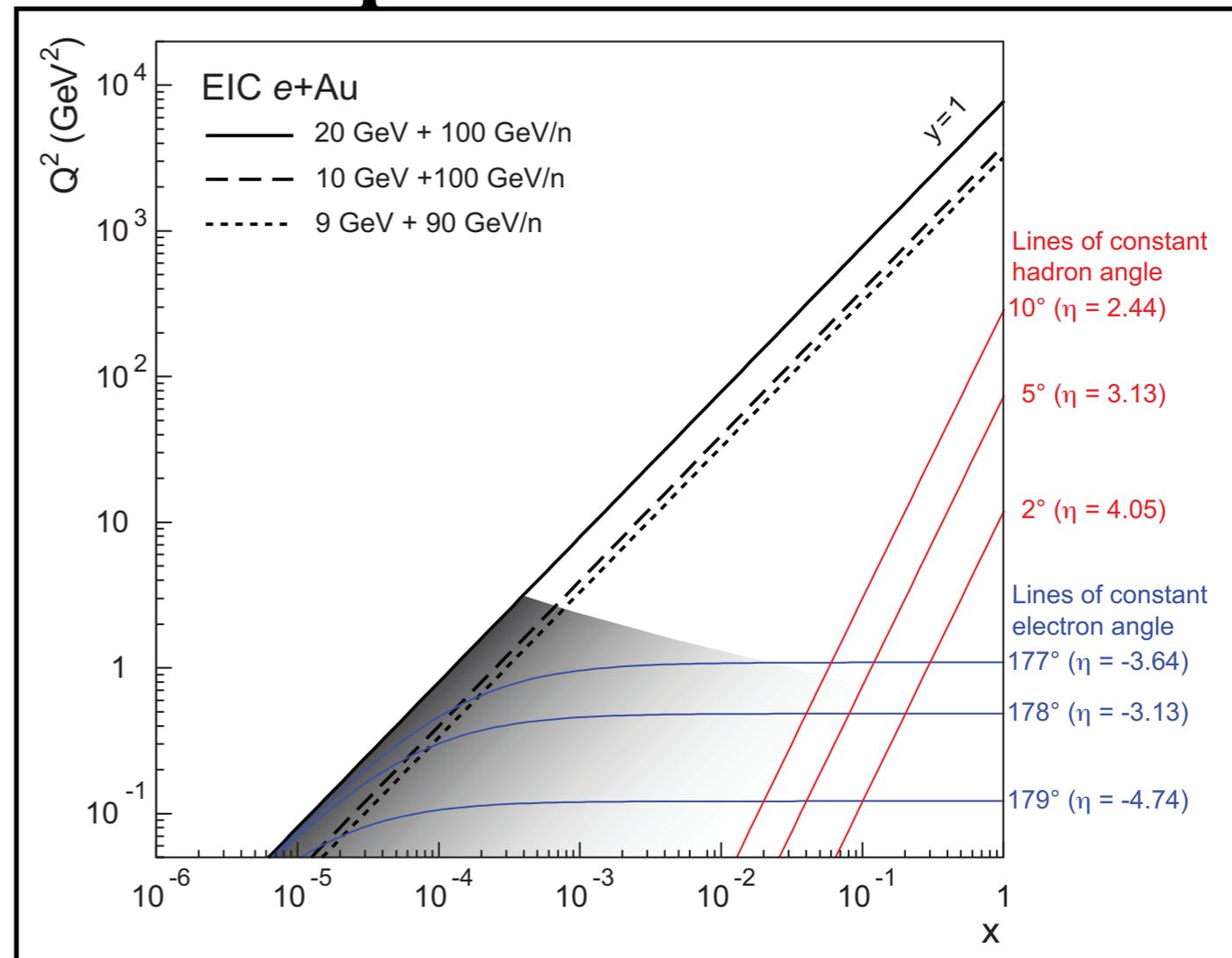
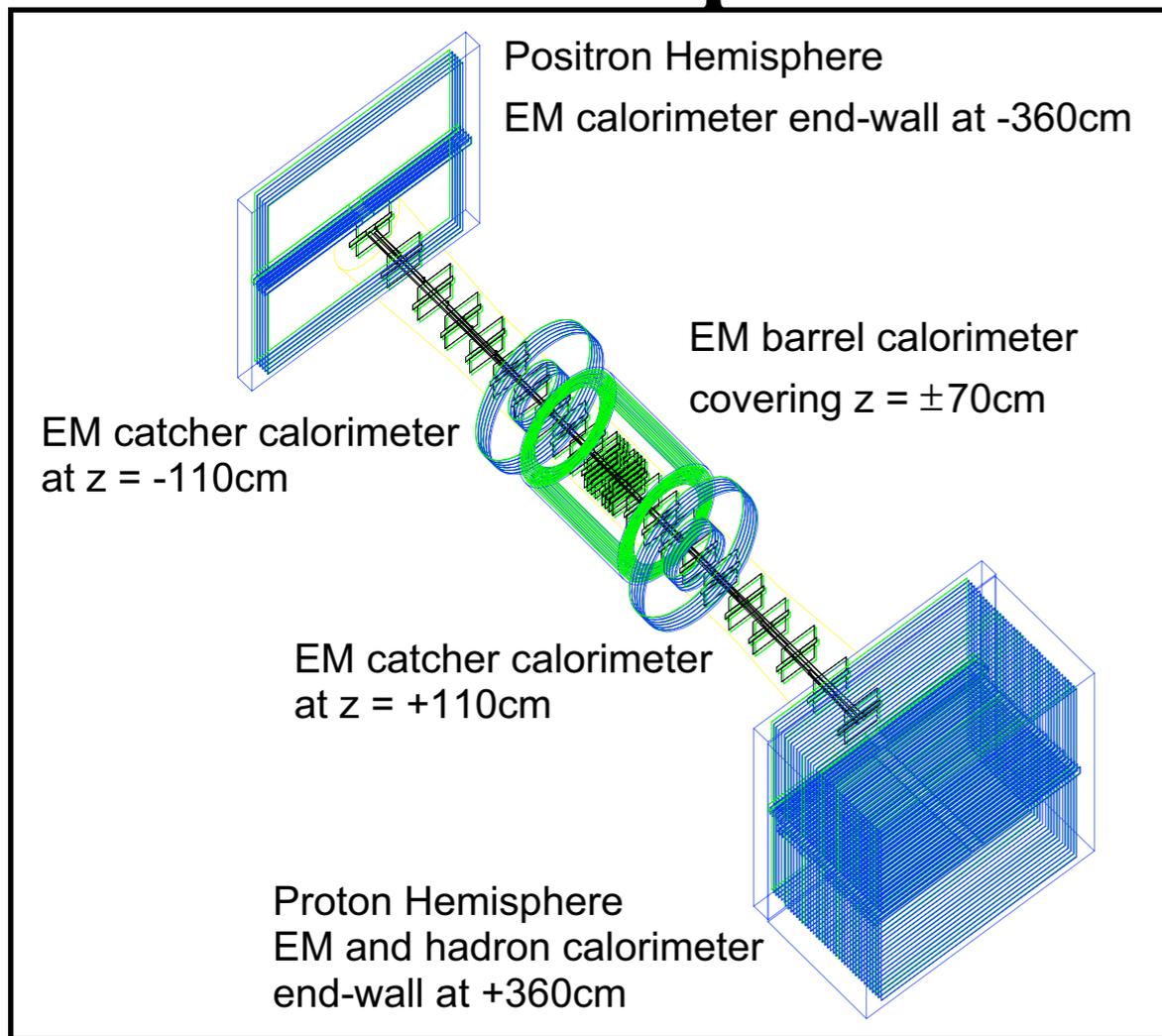
$$L_{eAu} \text{ (peak)}/n \sim 1.6 \cdot 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$$



Experimental Aspects



Experimental Aspects



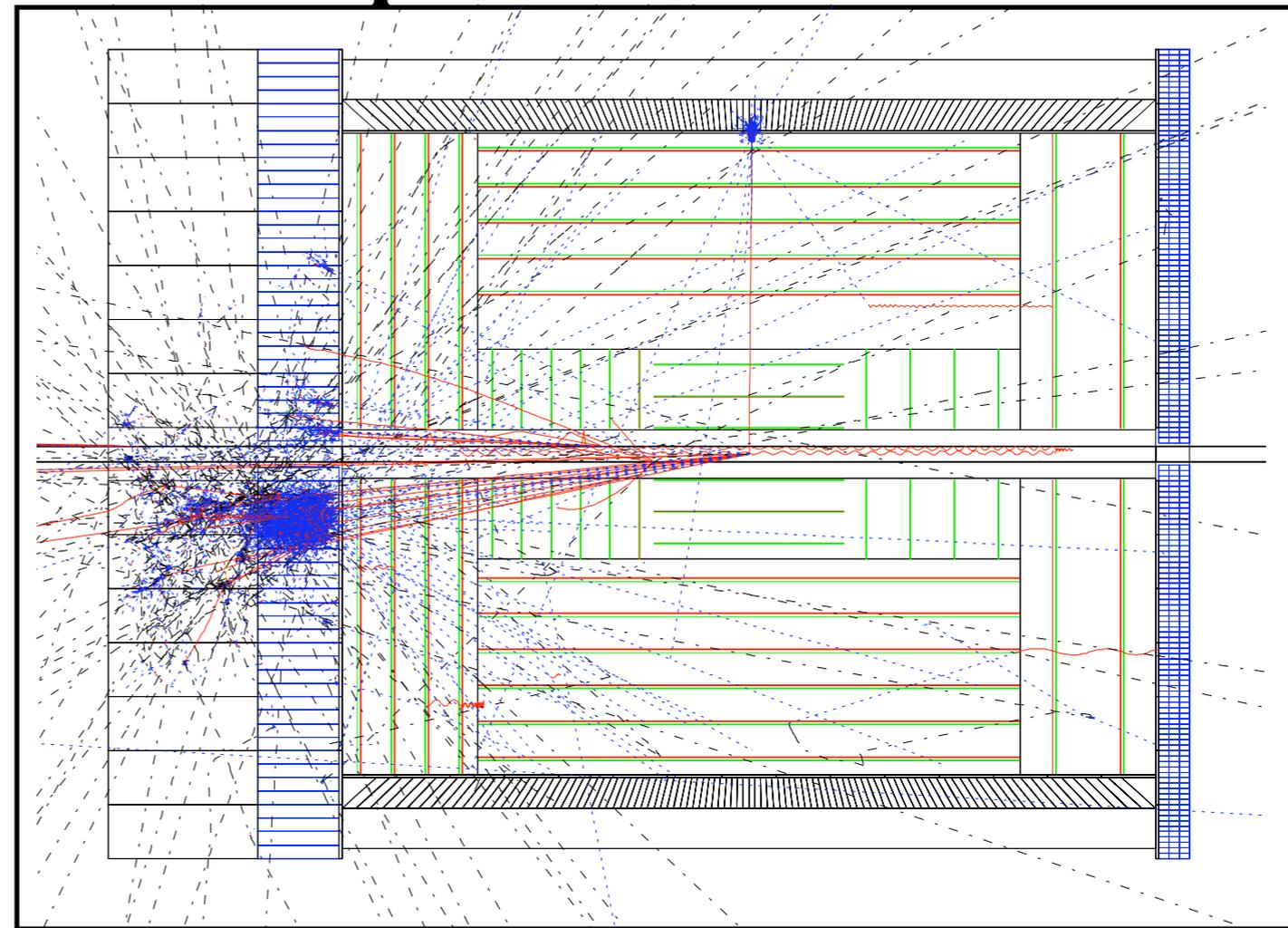
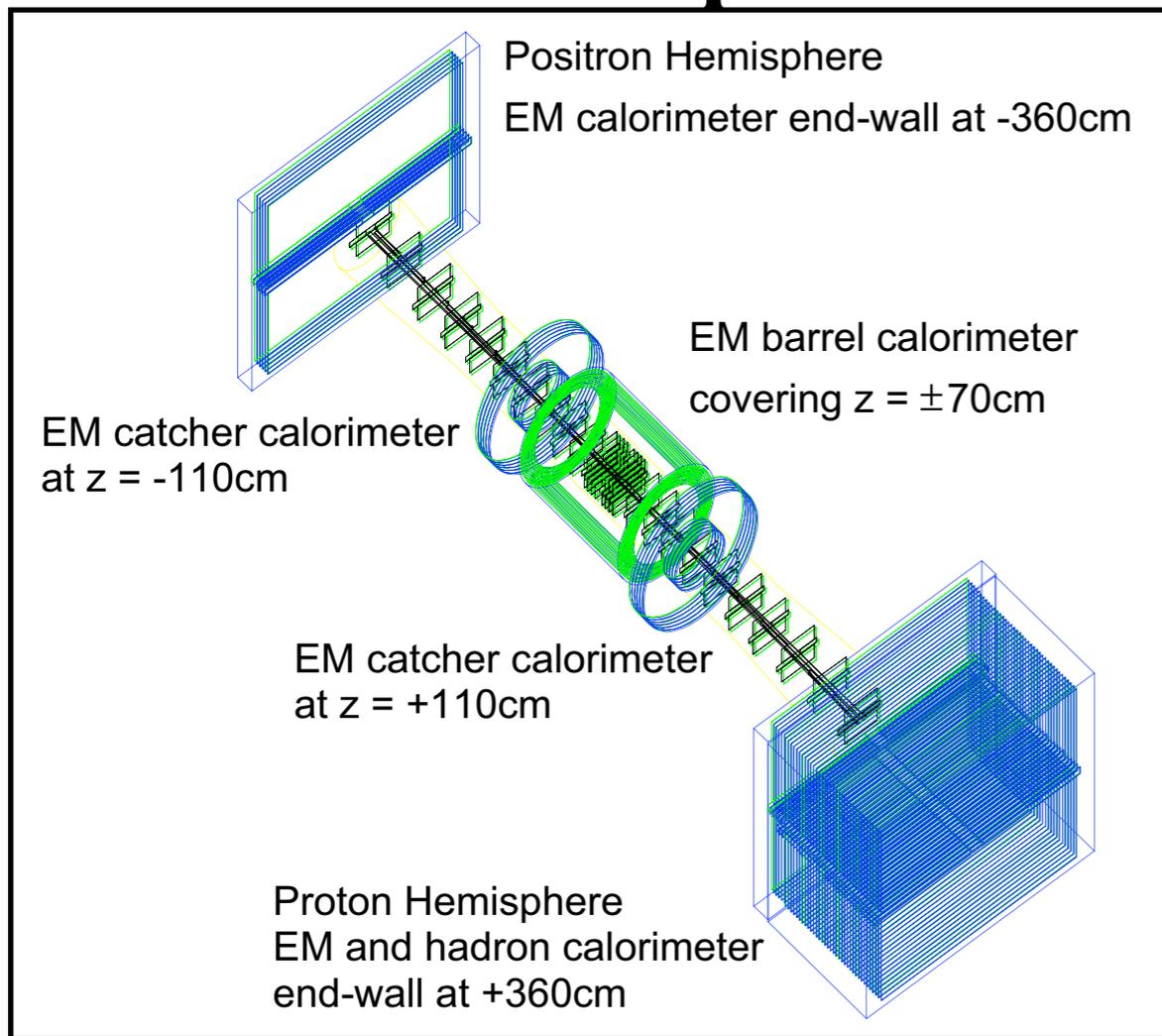
I. Abt, A. Caldwell, X. Liu, J. Sutiak, hep-ex 0407053

Concepts:

(a) Focus on the rear/forward acceptance and thus on low- x / high- x physics

- compact system of tracking and central electromagnetic calorimetry inside a magnetic dipole field and calorimetric end-walls outside

Experimental Aspects



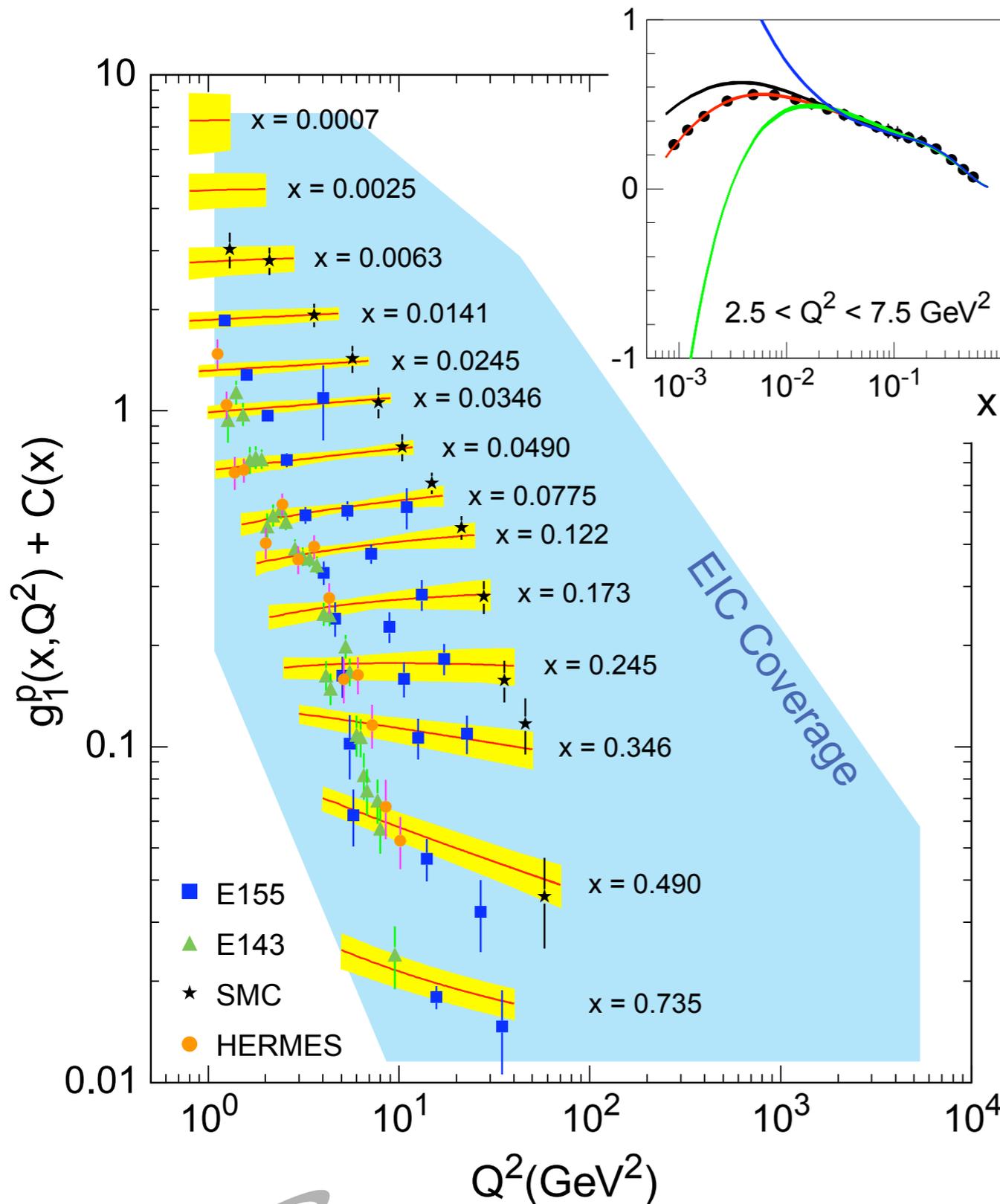
I. Abt, A. Caldwell, X. Liu, J. Sutiak, hep-ex 0407053

J. Pasukonis, B.Surrow, physics/0608290

Concepts:

- (a) Focus on the rear/forward acceptance and thus on low- x / high- x physics
 - compact system of tracking and central electromagnetic calorimetry inside a magnetic dipole field and calorimetric end-walls outside
- (b) Focus on a wide acceptance detector system similar to HERA experiments
 - allow for the maximum possible Q^2 range.

EIC as an $e+\vec{p}$ machine - The Quest for ΔG



Spin Structure of the Proton

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + L_q + L_g$$

quark contribution $\Delta\Sigma \approx 0.3$

gluon contribution $\Delta G \approx 1 \pm 1 ?$

ΔG : a “quotable” property of the proton
(like mass, charge)

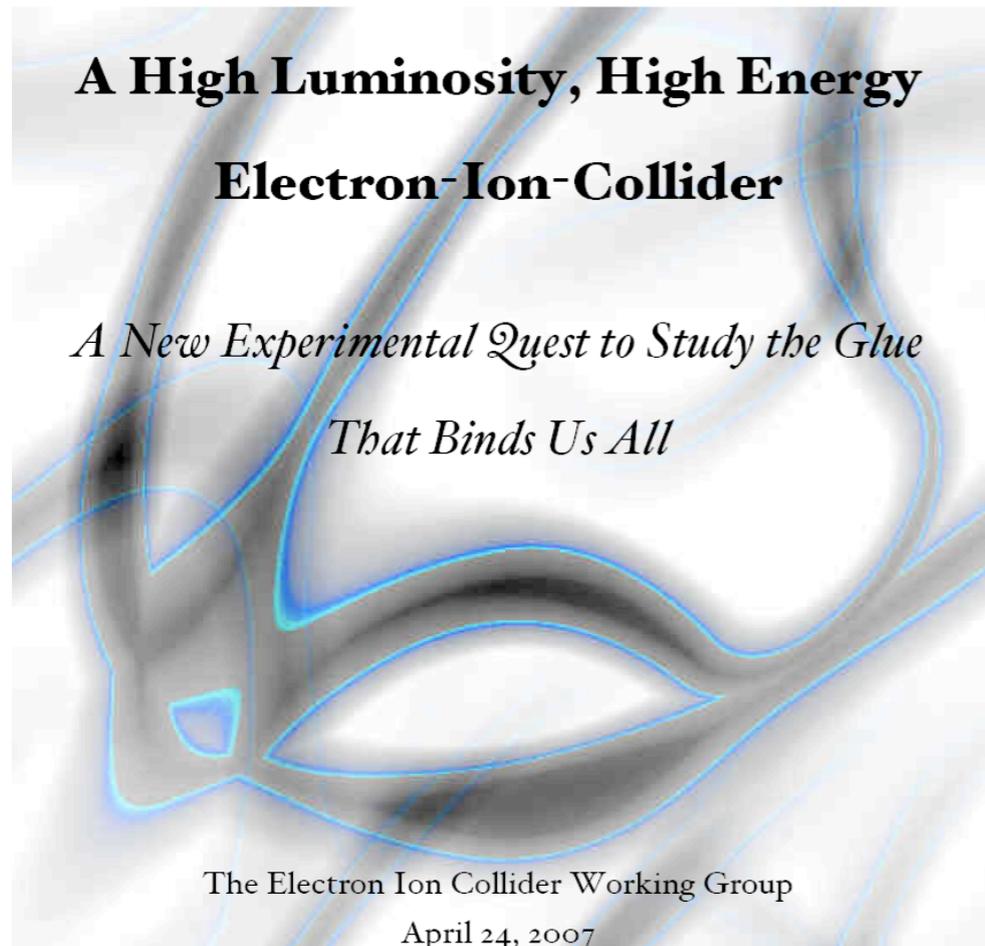
Measure through scaling violation:

$$\frac{dg_1}{d\log(Q^2)} \propto -\Delta g(x, Q^2)$$

$$\Delta G = \int_{x=0}^{x=1} \Delta g(x, Q^2) dx$$

Superb sensitivity to ΔG at
small x !

Status of the EIC Project:



Status of the EIC Project:

DRAFT V1 10-JAN-07

Exploring the 3D quark and gluon structure of the proton:
Electron scattering with present and future facilities*

H. Abramowicz,¹ A. Afanasev,² H. Avakian,³ M. Burkardt,⁴ V. Burkert,³ C. Munoz Camacho,⁵ A. Camsonne,³
A. Deshpande,⁶ F. Ellinghaus,⁷ L. Elouadrhiri,³ R. Ent,³ M. Garcon,⁸ G. Gavalian,⁹ M. Guidal,¹⁰
V. Guzey,¹¹ C. E. Hyde-Wright,⁹ X.-D. Ji,¹² A. Levy,¹ S. Liuti,¹³ W. Melnitchouk,³ R. Milner,¹⁴
Ch. Montag,¹⁵ D. Müller,¹⁶ R. Niyazov,³ B. Pasquini,¹⁷ S. Procureur,⁸ A. Radyushkin,^{9,3} J. Roche,¹⁸
F. Sabatie,⁸ A. Sandacz,¹⁹ A. Schäfer,¹⁶ M. Strikman,²⁰ M. Vanderhaeghen,^{21,3} E. Voutier,²² and Ch. Weiss³

¹Tel Aviv University, Tel Aviv, Israel

²Hampton University, Hampton, VA 23668, USA

³Jefferson Lab, Newport News, VA 23606, USA

⁴New Mexico State U.

⁵LANL

⁶Stony Brook U. and RIKEN-BNL Res. C.

⁷Colorado U.

⁸DAPNIA Saclay

⁹Old Dominion University, Norfolk, VA 23529, USA

¹⁰IPN Orsay

¹¹Bochum U.

¹²U. of Maryland

¹³U. of Virginia

¹⁴MIT

¹⁵BNL

¹⁶Regensburg U.

¹⁷Pavia U. and INFN

¹⁸Ohio U.

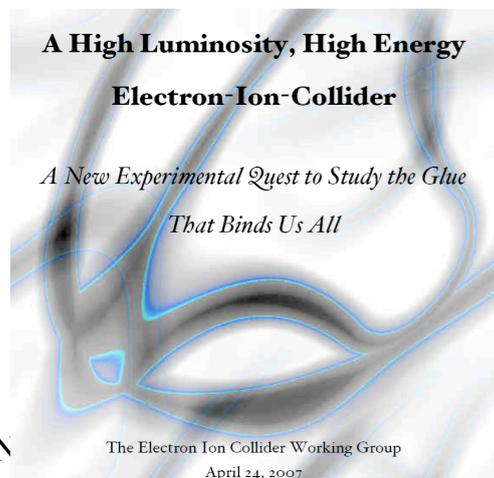
¹⁹Soltan I. Warsaw

²⁰Pennsylvania State University, University Park, PA 16802, USA

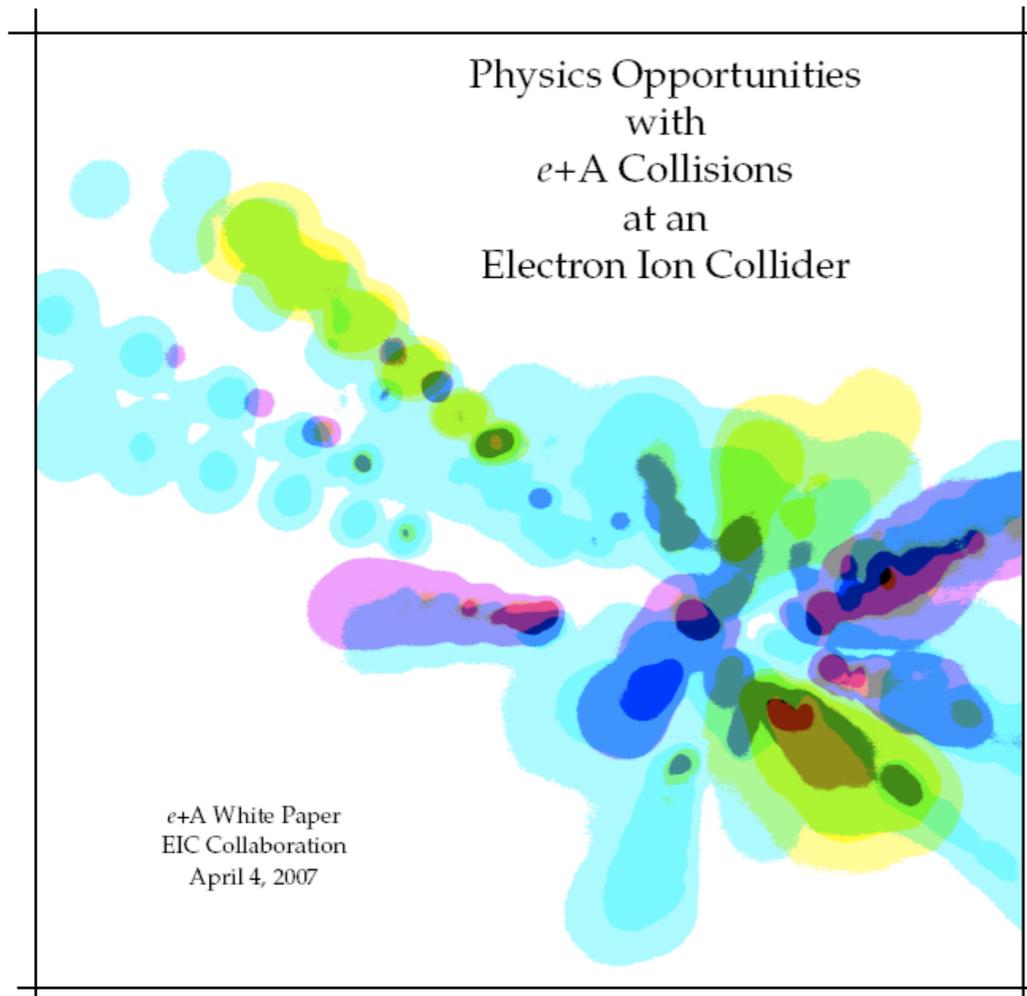
²¹William and Mary U.

²²LPSC Grenoble

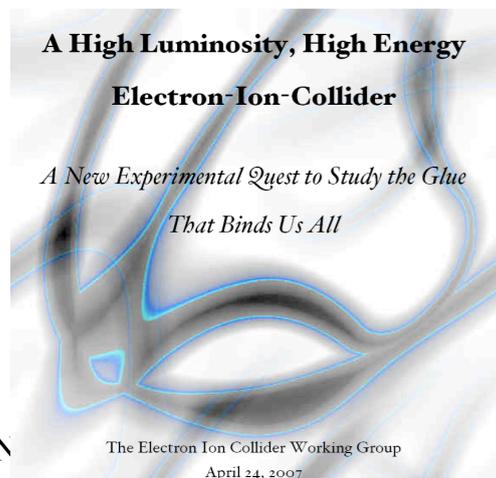
- The Electron Ion Collider (EIC) White Paper
- The GPD/DVCS White Paper
- Position Paper: $e+A$ Physics at an Electron Ion Collider
- The eRHIC machine: Accelerator Position Paper
- ELIC ZDR Draft



Status of the EIC Project:



- The Electron Ion Collider (EIC) White Paper
- The GPD/DVCS White Paper
- Position Paper: $e+A$ Physics at an Electron Ion Collider
- The eRHIC machine: Accelerator Position Paper
- ELIC ZDR Draft



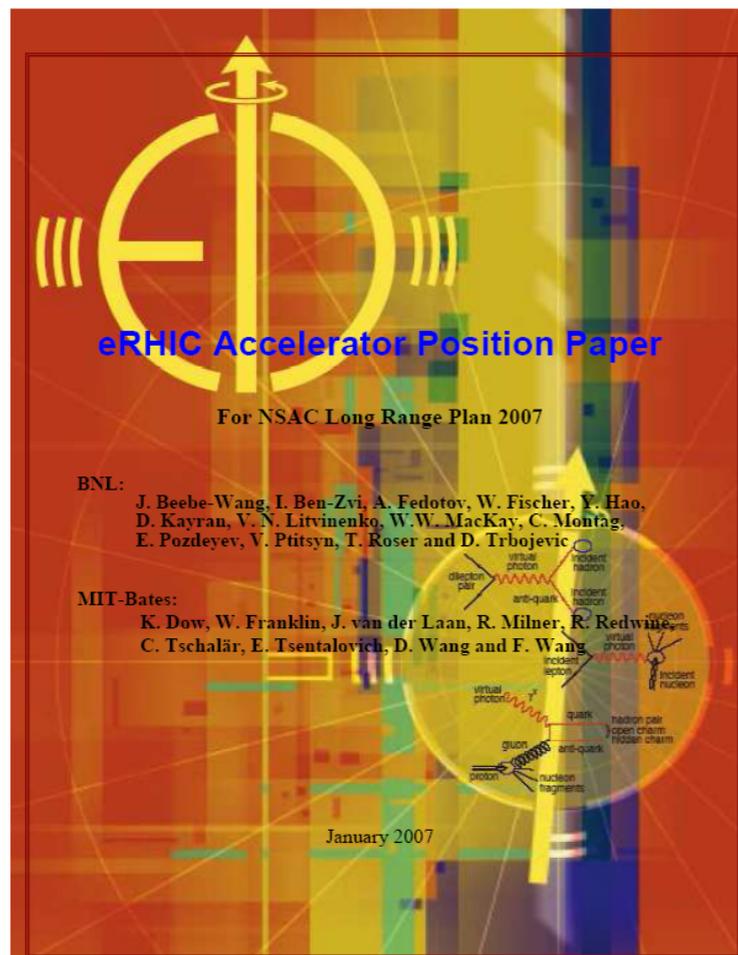
Exploring the 3D quark and gluon structure of the proton:
Electron scattering with present and future facilities*

H. Abramowicz,¹ A. Afanasev,² H. Avakian,³ M. Burkardt,⁴ V. Burkert,³ C. Munoz Camacho,⁵ A. Camsonne,³
A. Deshpande,⁶ F. Ellinghaus,⁷ L. Elouadrhiri,³ R. Ent,³ M. Garcon,⁸ G. Gavalian,⁹ M. Guidal,¹⁰
V. Guzey,¹¹ C. E. Hyde-Wright,⁹ X.-D. Ji,¹² A. Levy,¹ S. Liuti,¹³ W. Melnitchouk,³ R. Milner,¹⁴
Ch. Montag,¹⁵ D. Müller,¹⁶ R. Niyazov,³ B. Pasquini,¹⁷ S. Procureur,⁸ A. Radyushkin,^{9,3} J. Roche,¹⁸
F. Sabatie,⁸ A. Sandacz,¹⁹ A. Schäfer,¹⁶ M. Strikman,²⁰ M. Vanderhaeghen,^{21,3} E. Voutier,²² and Ch. Weiss³

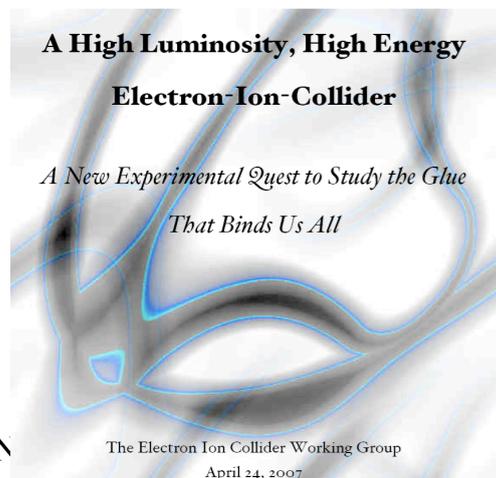
¹Tel Aviv University, Tel Aviv, Israel
²Hampton University, Hampton, VA 23668, USA
³Jefferson Lab, Newport News, VA 23606, USA
⁴New Mexico State U.
⁵LANL
⁶Stony Brook U. and RIKEN-BNL Res. C.
⁷Colorado U.
⁸DAPNIA Saclay
⁹Old Dominion University, Norfolk, VA 23529, USA
¹⁰IPN Orsay
¹¹Bochum U.
¹²U. of Maryland
¹³U. of Virginia
¹⁴MIT
¹⁵BNL
¹⁶Regensburg U.
¹⁷Pavia U. and INFN
¹⁸Ohio U.
¹⁹Sultan I. Warsaw
²⁰Pennsylvania State University, University Park, PA 16802, USA
²¹William and Mary U.
²²LPSC Grenoble

macl@bnl.gov

Status of the EIC Project:



- The Electron Ion Collider (EIC) White Paper
- The GPD/DVCS White Paper
- Position Paper: $e+A$ Physics at an Electron Ion Collider
- The eRHIC machine: Accelerator Position Paper
- ELIC ZDR Draft

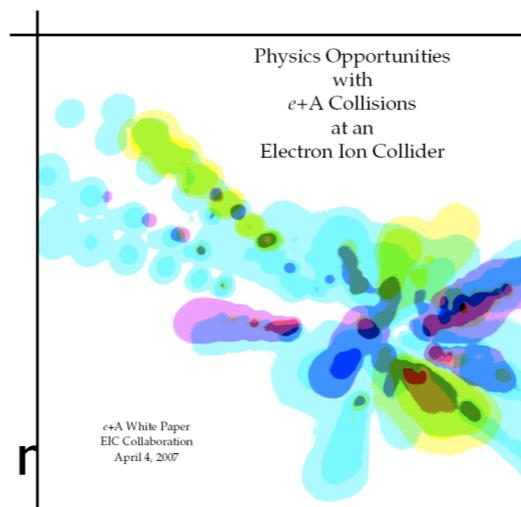


DRAFT V1 10-JAN-07

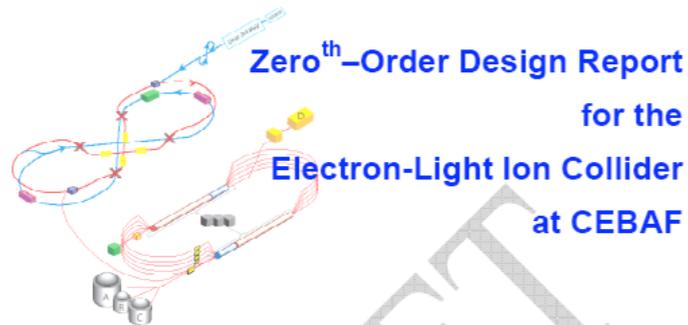
**Exploring the 3D quark and gluon structure of the proton:
Electron scattering with present and future facilities***

H. Abramowicz,¹ A. Afanasev,² H. Avakian,³ M. Burkardt,⁴ V. Burkert,³ C. Munoz Camacho,⁵ A. Camsonne,³ A. Deshpande,⁶ F. Ellinghaus,⁷ L. Elouadrhiri,³ R. Ent,³ M. Garcon,⁸ G. Gavalian,⁹ M. Guidal,¹⁰ V. Guzey,¹¹ C. E. Hyde-Wright,⁹ X.-D. Ji,¹² A. Levy,¹ S. Liuti,¹³ W. Melnitchouk,³ R. Milner,¹⁴ Ch. Montag,¹⁵ D. Müller,¹⁶ R. Niyazov,³ B. Pasquini,¹⁷ S. Procureur,⁸ A. Radyushkin,^{9,3} J. Roche,¹⁸ F. Sabatie,⁸ A. Sandacz,¹⁹ A. Schäfer,¹⁶ M. Strikman,²⁰ M. Vanderhaeghen,^{21,3} E. Voutier,²² and Ch. Weiss³

¹Tel Aviv University, Tel Aviv, Israel
²Hampton University, Hampton, VA 23668, USA
³Jefferson Lab, Newport News, VA 23606, USA
⁴New Mexico State U.
⁵LANL
⁶Stony Brook U. and RIKEN-BNL Res. C.
⁷Colorado U.
⁸DAPNIA Saclay
⁹Old Dominion University, Norfolk, VA 23529, USA
¹⁰IPN Orsay
¹¹Bochum U.
¹²U. of Maryland
¹³U. of Virginia
¹⁴MIT
¹⁵BNL
¹⁶Regensburg U.
¹⁷Pavia U. and INFN
¹⁸Ohio U.
¹⁹Sultan I. Warsaw
²⁰Pennsylvania State University, University Park, PA 16802, USA
²¹William and Mary U.
²²LPSC Grenoble



Status of the EIC Project:



**Zeroth-Order Design Report
for the
Electron-Light Ion Collider
at CEBAF**

A. Afanasev, A. Bogacz, A. Bruell, L. Cardman, Y. Chao, S. Chattopadhyay, E. Chudakov, P. Degtiarenko, J. Delaysen, Ya. Derbenev, R. Ent, P. Evtushenko, A. Freyberger, J. Grames, A. Hutton, R. Kazimi, G. Krafft, R. Li, L. Merminga, M. Poelker, A. Thomas, C. Weiss, B. Wojtsekhowski, B. Yunn, Y. Zhang
Thomas Jefferson National Accelerator Facility
Newport News, Virginia, USA

W. Fischer, C. Montag
Brookhaven National Laboratory
Upton, New York, USA

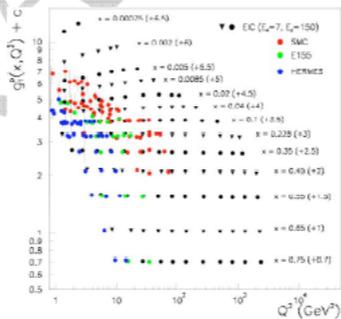
V. Danilov
Oak Ridge National Laboratory
Oak Ridge, Tennessee, USA

V. Dudnikov
Brookhaven Technology Group
New York, New York, USA

P. Ostroumov
Argonne National Laboratory
Argonne, Illinois, USA

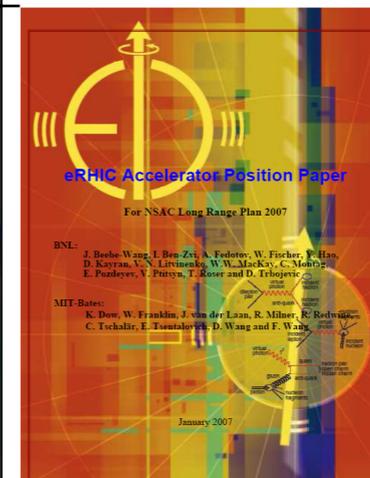
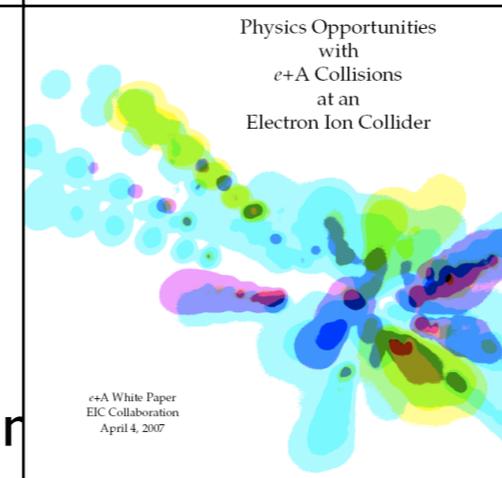
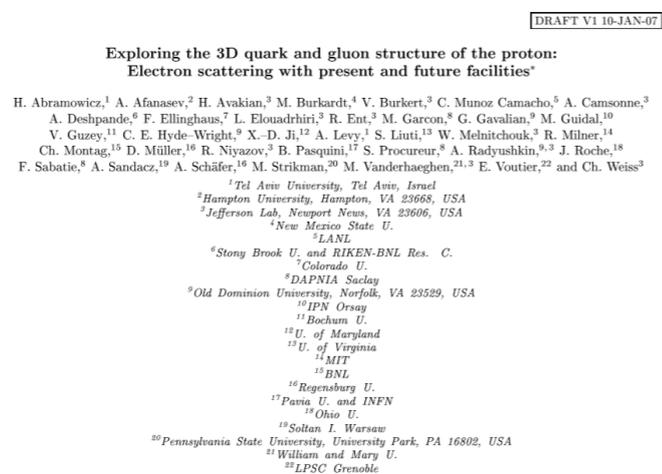
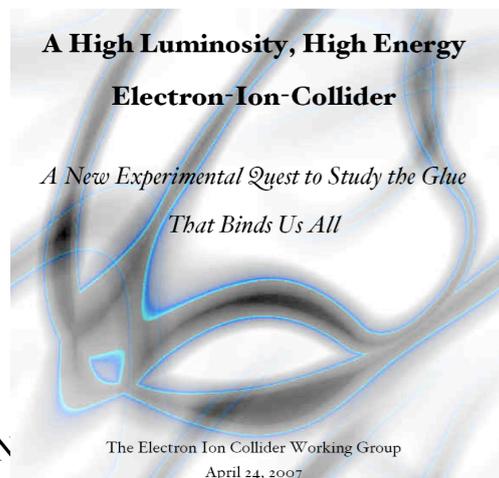
V. Derenchuk
Indiana University Cyclotron Facility
Bloomington, Indiana, USA

A. Belov
Institute of Nuclear Research
Moscow-Troitsk, Russia



Editors: Ya. Derbenev, L. Merminga, Y. Zhang

- The Electron Ion Collider (EIC) White Paper
- The GPD/DVCS White Paper
- Position Paper: $e+A$ Physics at an Electron Ion Collider
- The eRHIC machine: Accelerator Position Paper
- ELIC ZDR Draft

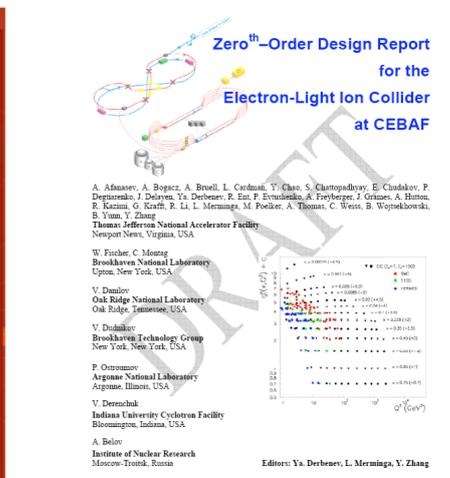
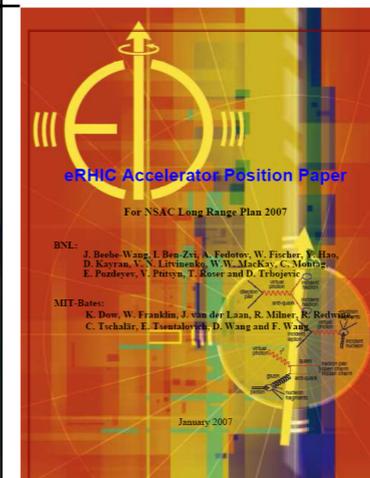
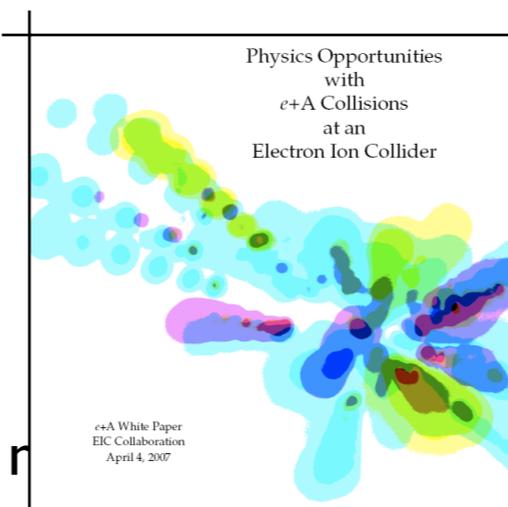
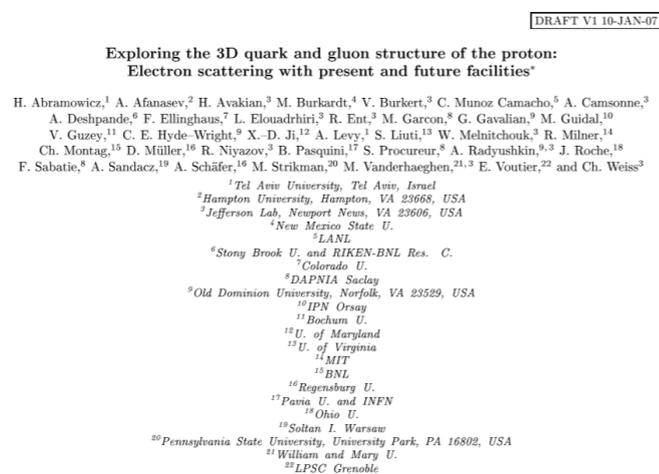
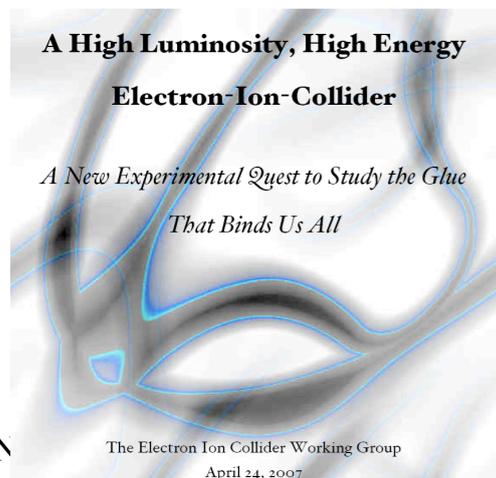


Status of the EIC Project:

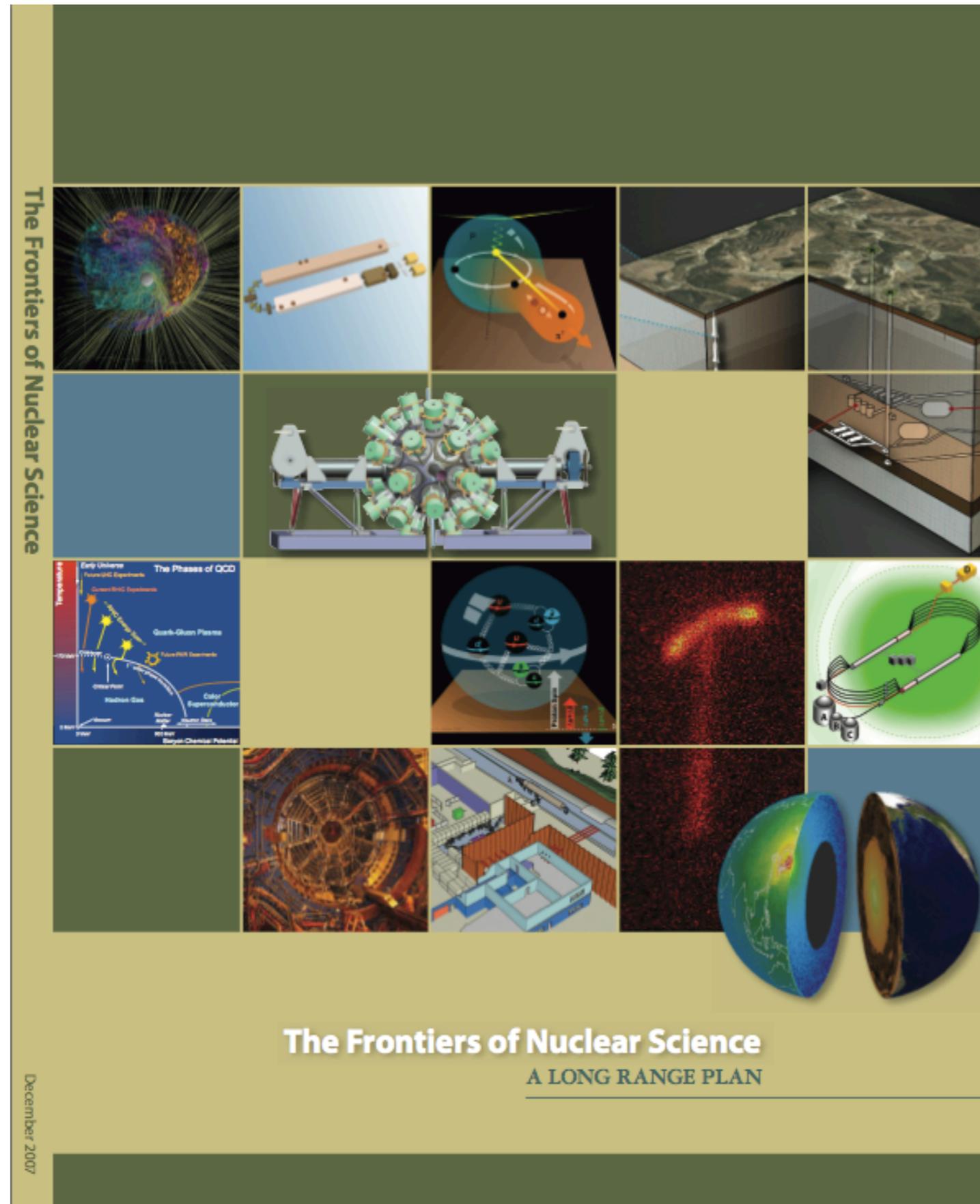
Available at:

- NSAC LRP2007 home page
- Rutgers Town Meeting page
- <http://web.mit.edu/eicc>

- The Electron Ion Collider (EIC) White Paper
- The GPD/DVCS White Paper
- Position Paper: $e+A$ Physics at an Electron Ion Collider
- The eRHIC machine: Accelerator Position Paper
- ELIC ZDR Draft



2007 NSAC Long Range Plan



2007 NSAC Long Range Plan

FURTHER INTO THE FUTURE

Gluons and their interactions are critical to QCD. But their properties and dynamics in matter remain largely unexplored. Recent theoretical breakthroughs and experimental results suggest that both nucleons and nuclei, when viewed at high energies, appear as dense systems of gluons, creating fields whose intensity may be the strongest allowed in nature. The emerging science of this universal gluonic matter drives the development of a next-generation facility, the high-luminosity Electron-Ion Collider (EIC). The EIC's ability to collide high-energy electron beams with high-energy ion beams will provide access to those regions in the nucleon and nuclei where their structure is dominated

by gluons. Moreover, polarized beams in the EIC will give unprecedented access to the spatial and spin structure of gluons in the proton.

An EIC with polarized beams has been embraced by the U.S. nuclear science community as embodying the vision for reaching the next QCD frontier. EIC would provide unique capabilities for the study of QCD well beyond those available at existing facilities worldwide and complementary to those planned for the next generation of accelerators in Europe and Asia. While significant progress has been made in developing concepts for an EIC, many open questions remain. Realization of an EIC will require advancements in accelerator science and technology, and detector research and development. The nuclear science community has recognized the importance of this future facility and makes the following recommendation.

2007 NSAC Long Range Plan

FURTHER INTO THE FUTURE

Gluons and their interactions are critical to QCD. But their properties and dynamics in matter remain largely unexplored. Recent theoretical breakthroughs and experimental results suggest that both nucleons and nuclei, when viewed at high energies, appear as dense systems of gluons, creating fields whose intensity may be the strongest allowed

in nature. The emerging science of this universal gluonic matter drives the development of a next-generation facility, the high-luminosity Electron-Ion Collider (EIC). The EIC's ability to collide high-energy electron beams with high-energy ion beams will provide access to those regions in the nucleon and nuclei where their structure is dominated

by gluons. Moreover, polarized beams in the EIC will give unprecedented access to the spatial and spin structure of gluons in the proton.

An EIC with polarized beams has been embraced by the U.S. nuclear science community as embodying the vision for reaching the next QCD frontier. EIC would provide unique capabilities for the study of QCD well beyond those available at existing facilities worldwide and complementary to those planned for the next generation of accelerators in Europe and Asia. While significant progress has been made in developing concepts for an EIC, many open questions remain. Realization of an EIC will require advancements in accelerator science and technology, and detector research and development. The nuclear science community has recognized the importance of this future facility and makes the following recommendation.

2007 NSAC Long Range Plan

FURTHER INTO THE FUTURE

Gluons and their interactions are critical to QCD. But their properties and dynamics in matter remain largely unexplored. Recent theoretical breakthroughs and experimental results suggest that both nucleons and nuclei, when viewed at high energies, appear as dense systems of gluons, creating fields whose intensity may be the strongest allowed

in nature. The emerging science of this universal gluonic matter drives the development of a next-generation facility, the high-luminosity Electron-Ion Collider (EIC). The EIC's ability to collide high-energy electron beams with high-energy ion beams will provide access to those regions in the nucleon and nuclei where their structure is dominated

by gluons. Moreover, polarized beams in the EIC will give unprecedented access to the spatial and spin structure of gluons in the proton.

An EIC with polarized beams has been embraced by the U.S. nuclear science community as embodying the vision for reaching the next QCD frontier. EIC would provide unique capabilities for the study of QCD well beyond those available at existing facilities worldwide and complementary to those planned for the next generation of accelerators in

Europe and Asia. While significant progress has been made in developing concepts for an EIC, many open questions remain. Realization of an EIC will require advancements in accelerator science and technology, and detector research and development. The nuclear science community has recognized the importance of this future facility and makes the following recommendation.

2007 NSAC Long Range Plan

FURTHER INTO THE FUTURE

Gluons and their interactions are critical to QCD. But their properties and dynamics in matter remain largely unexplored. Recent theoretical breakthroughs and experimental results suggest that both nucleons and nuclei, when viewed at high energies, appear as dense systems of gluons, creating fields whose intensity may be the strongest allowed

in nature. The emerging science of this universal gluonic matter drives the development of a next-generation facility, the high-luminosity Electron-Ion Collider (EIC). The EIC's ability to collide high-energy electron beams with high-energy ion beams will provide access to those regions in the nucleon and nuclei where their structure is dominated

by gluons. Moreover, polarized beams in the EIC will give unprecedented access to the spatial and spin structure of gluons in the proton.

An EIC with polarized beams has been embraced by the U.S. nuclear science community as embodying the vision for reaching the next QCD frontier. EIC would provide unique capabilities for the study of QCD well beyond those available at existing facilities worldwide and complementary to those planned for the next generation of accelerators in

Europe and Asia. While significant progress has been made in developing concepts for an EIC, many open questions remain. Realization of an EIC will require advancements in accelerator science and technology, and detector research and development. The nuclear science community has recognized the importance of this future facility and makes the following recommendation.

We recommend the allocation of resources to develop accelerator and detector technology necessary to lay the foundation for a polarized Electron-Ion Collider. The EIC would explore the new QCD frontier of strong color fields in nuclei and precisely image the gluons in the proton.

What is happening now

- EIC “Collaboration” formed in 2007
 - ➔ Bi-Annual collaboration meetings
 - ▶ Next meeting, 11th - 13th December, 2008, LBNL
- INT
 - ➔ Week long workshop - Autumn 2009
 - ➔ 3-month programme just approved - Autumn 2010
- e+A working group
 - ➔ Convenors: T. Ullrich, D. Morrison, R. Venugopalan, V. Guzey
 - ➔ weekly(ish) meetings at BNL + phone bridge
 - ▶ <http://www.eic.bnl.gov/> for details (and previous seminars)
 - ➔ Current focus of work - understanding diffraction in e+A physics
 - ▶ How do we measure diffractive events? ⇒ detector design
 - ➔ Last week of September, mini-workshop at BNL to start work on putting together an e+A MC code for both DIS and diffractive events

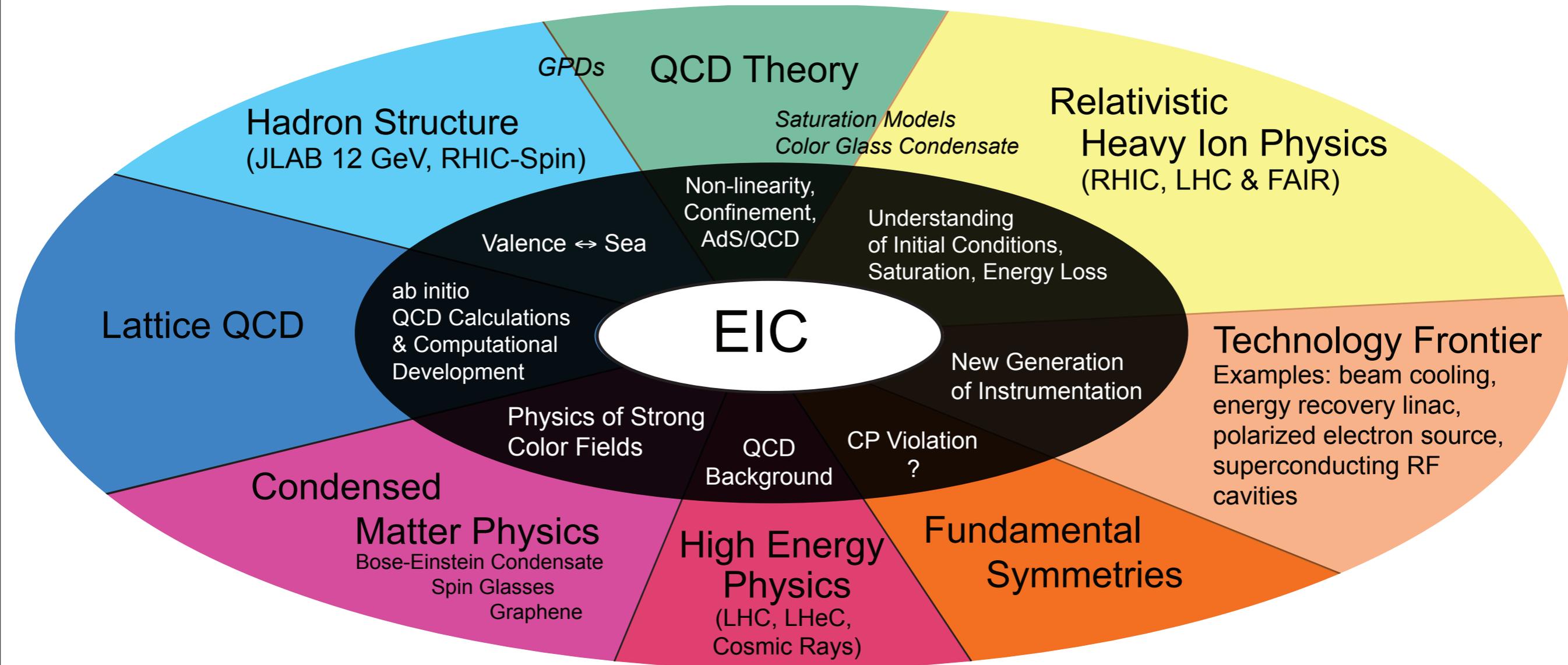
Summary

An EIC presents a unique opportunity in high energy nuclear physics and precision QCD physics

e+A	Polarized e+p
<ul style="list-style-type: none">◆ Study the Physics of Strong Colour Fields<ul style="list-style-type: none">• Establish (or not) the existence of the saturation regime• Explore non-linear QCD• Measure momentum & space-time of glue◆ Study the nature of colour singlet excitations (Pomerons)◆ Test and study the limits of universality (eA vs. pA)	<ul style="list-style-type: none">◆ Precisely image the sea-quarks and gluons to determine the spin, flavour and spatial structure of the nucleon

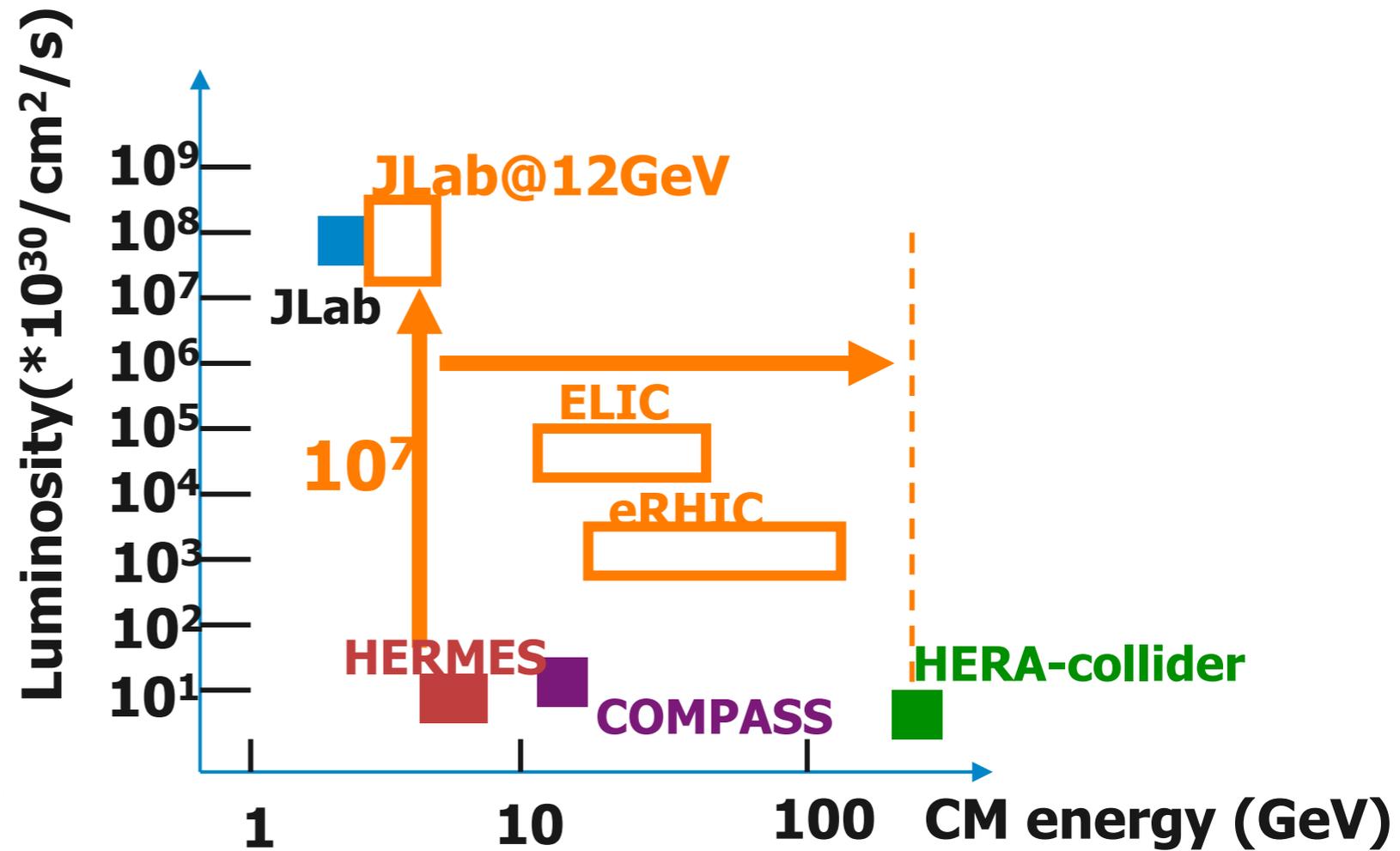
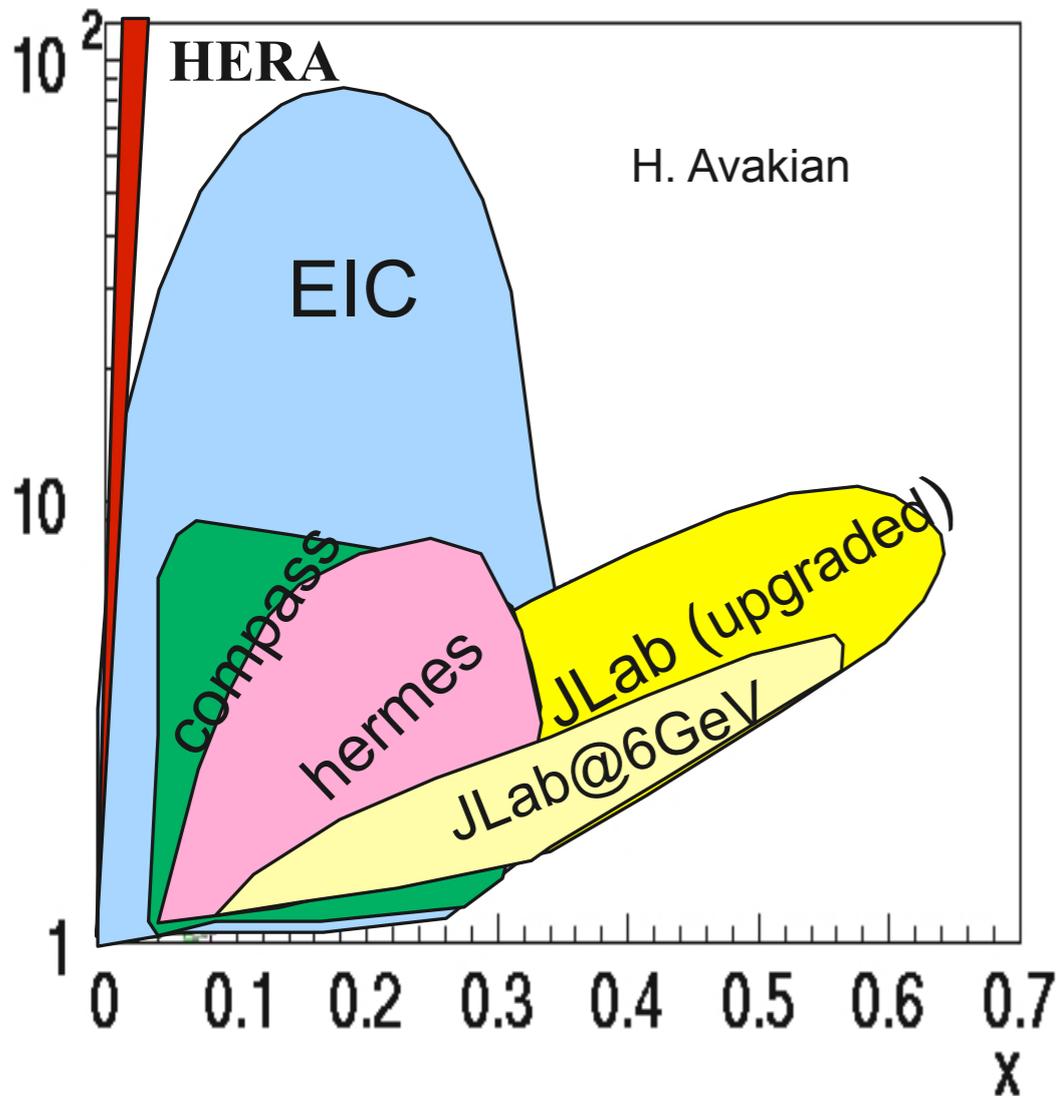
- Embraced by NSAC in Long Range PPlan
 - Recommendation of \$30M for R&D over next 5 years
- EIC Long Term Goal - start construction in next decade
- Possibility of Staged Approach
 - Cheap (no civil construction costs)
 - Early time-scale for realisation (operation by ~2016)
 - Cons - lower energy and luminosity than full design

Connection to other fields



BACKUP SLIDES

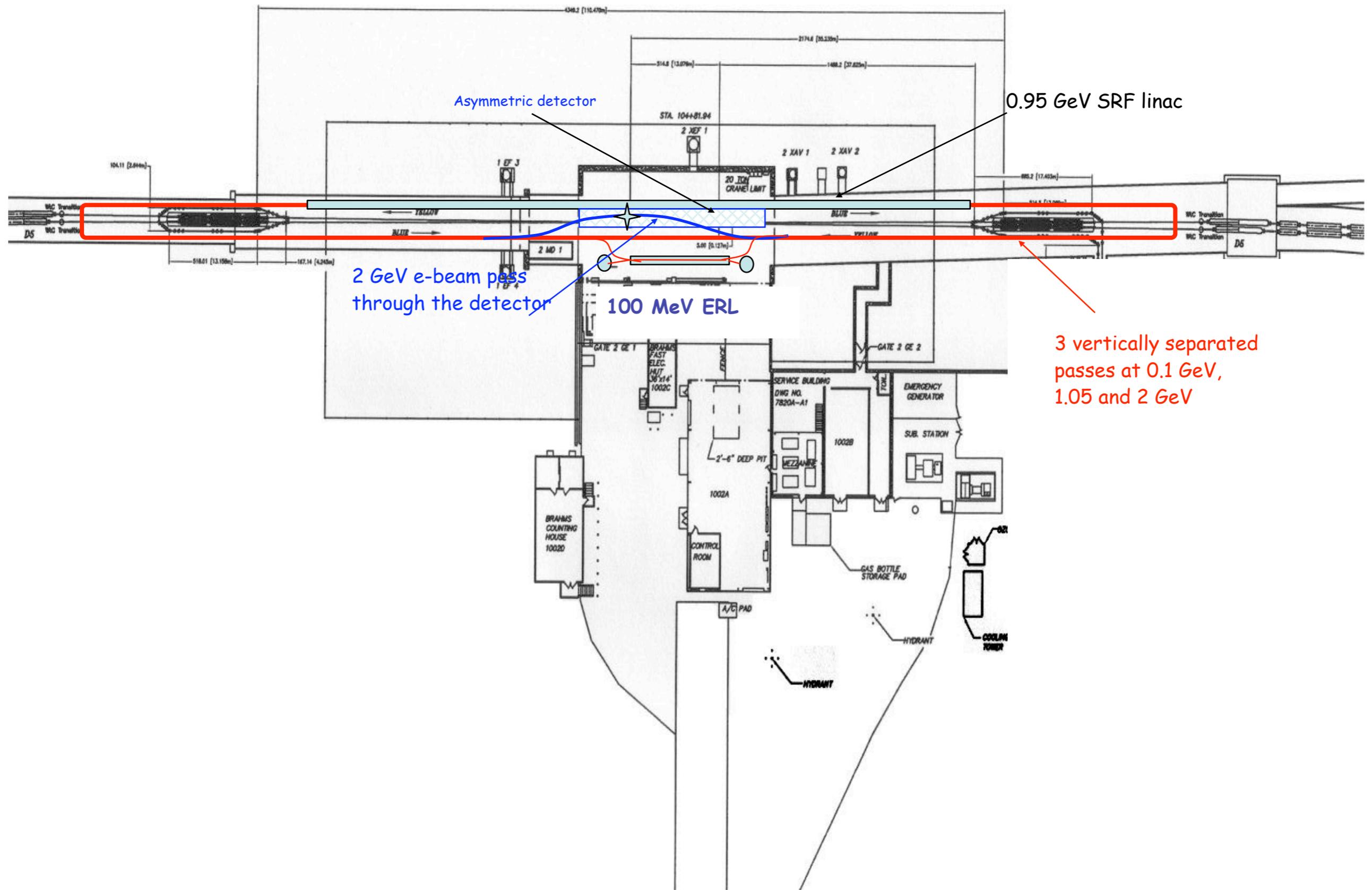
JLab Upgrade



Staged approach to an EIC

- **MEIC: Medium Energy Electron-Ion Collider**
 - Located at IP2 (with a modest detector)
 - 2 GeV e^- x 250 GeV p (45 GeV c.m.), $L \sim 10^{32} \text{ cm}^{-2} \text{ sec}^{-1}$
- **eRHIC - Full energy, nominal luminosity, inside RHIC tunnel**
 - Polarized 20 GeV e^- x 325 GeV p (160 GeV c.m.), $L \sim 4 \cdot 10^{33} \text{ cm}^{-2} \text{ sec}^{-1}$
 - 30 GeV e^- x 120 GeV/n Au (120 GeV c.m.), $L \sim 10^{31} \text{ cm}^{-2} \text{ sec}^{-1}$
 - 20 GeV e^- x 120 GeV/n Au (120 GeV c.m.), $L \sim 5 \cdot 10^{31} \text{ cm}^{-2} \text{ sec}^{-1}$
- **eRHIC - High luminosity at reduced energy, inside RHIC tunnel**
 - Polarized 10 GeV e^- x 325 GeV p, $L \sim 10^{35} \text{ cm}^{-2} \text{ sec}^{-1}$
 - Smaller improvements (3-4 fold) in e-Ion collisions

MEIC with 2 GeV ERL



Staging of eRHIC - Cost

- **MEIC: Medium Energy Electron-Ion Collider**
 - **Cost estimate - \$150M** (in 2007 \$)
 - 90% of ERL hardware will be use in the phase I (and will reduce cost of eRHIC)
 - Possible use of the detector components for eRHIC detectors
- **eRHIC - phase I**
 - **Based on present RHIC beam intensities**
 - With coherent electron cooling requirements on the electron beam current is 25 mA
 - 20 GeV, 25 mA electron beam losses 1.92 MW total for synchrotron radiation.
 - 30 GeV, 5 mA electron beam loses 1.98 MW for synchrotron radiation
 - Power density is 1 kW/meter and is well within B-factory limits (8 kW/m)
- **eRHIC - phase II**
 - **Requires crab cavities, new injections, Cu-coating of RHIC vacuum chambers, new level of intensities in RHIC**
 - Polarized electron source current of 400 mA
 - 10 GeV, 400 mA electron beam losses 1.96 MW total for synchrotron radiation, power density is 1 kW/meter