

Dynamic aspects of the QCD critical point

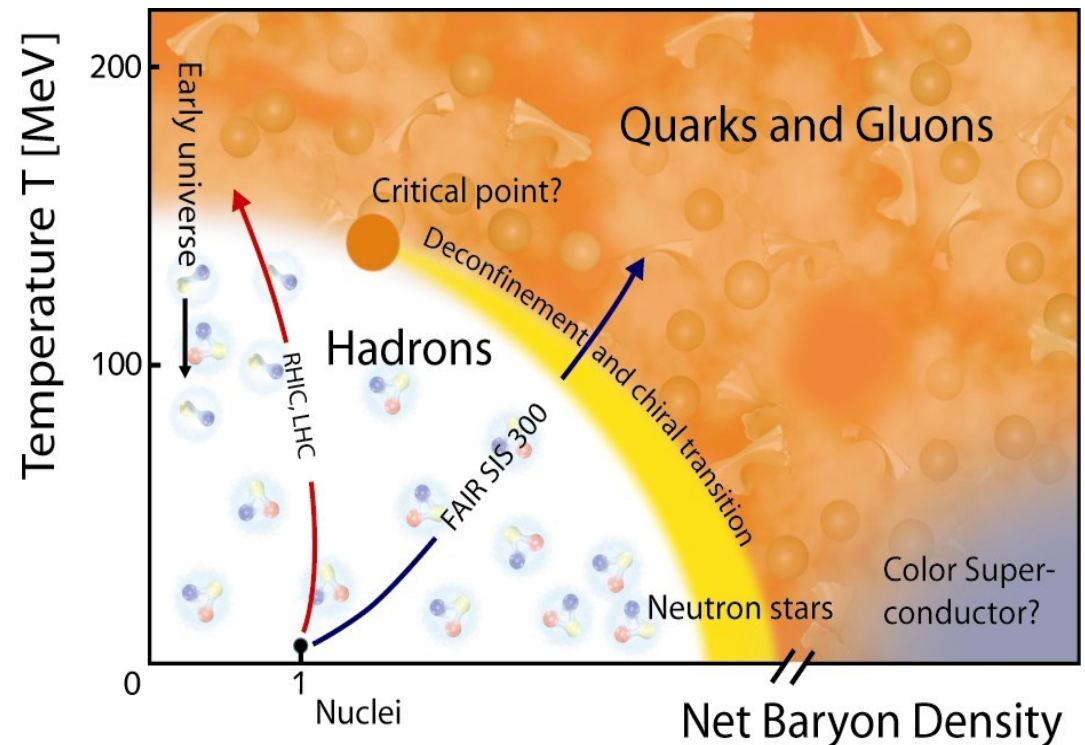
Hiro Fujii

HF, PRD67, 094018
HF-M.Ohtani, PRD70, 014016
and work in progress

Introduction

- LQCD + Models suggested the QCD-CP in T- μ plane
 - cross-over along T axis
 - 1st order along μ axis

Stephanov-Rajagopal-Shuryak
PRL 81, 4816 (98); PRD60,114028(98)



Introduction

- LQCD + Models suggested the QCD-CP in T- μ
 - cross-over along T axis
 - 1st order along μ axis
- Still difficult to conclude with LQCD+Models today
 - about the location (even the existence?)
- For exp'tal confirmation, we need good signatures
- Important to clarify the character of the QCD-CP

Stephanov-Rajagopal-Shuryak
PRL 81, 4816 (98), PRD60,114028(98)

For review, Stephanov
Fodor-Katz, Ejiri et al, de Forcrand-Philipsen

How to use chiral models

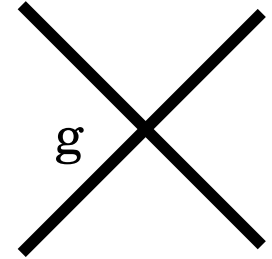
Stephanov et al., Hatta-Ikeda, HF-Ohtani, Sasaki et al.

- **Generic features** of the CP can be demonstrated in simple models
 - singular behavior of fluctuations, ..., etc.
- This talk: **spectral properties near the QCD-CP**
 - within a simple, but consistent model
 - message: 'pole' mass of σ is non-zero @QCD-CP
- Take **NJL model**, although t -dep Ginzburg-Landau theory may be used

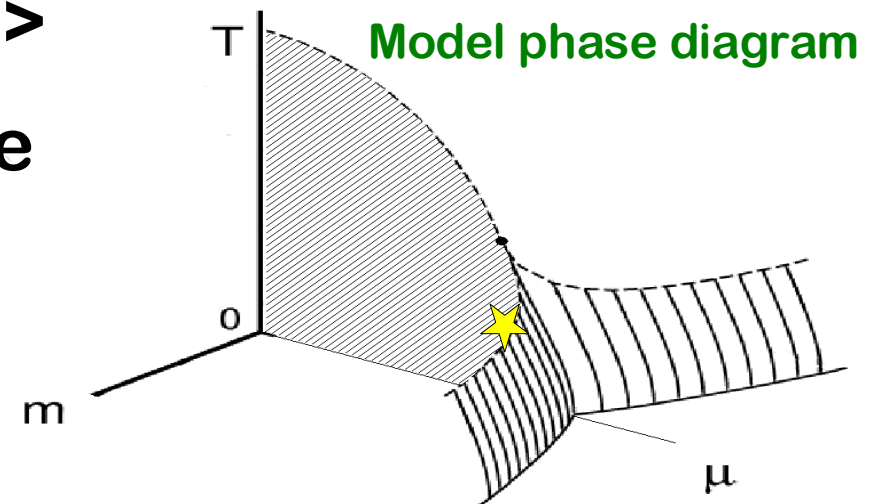
HF-Ohtani, PRD70, 014016

NJL model

- $$\mathcal{L} = \bar{q}(i\partial - m)q + g[(\bar{q}q)^2 + (\bar{q}i\gamma_5\tau^a q)^2]$$



- the simplest quark dynamics
- chiral symmetry breaking by $q^{\text{bar}}-q$ attraction
- no dynamic gluons, no confinement (nucleons)
- σ and π as flucsts of $\langle q^{\text{bar}} q \rangle$
- In reality, $m \neq 0$, of course

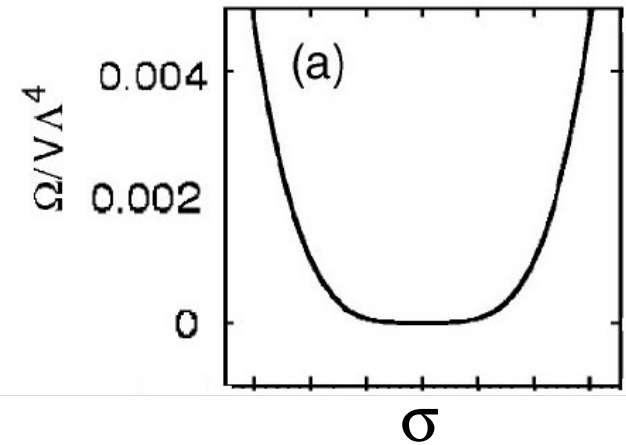


Chiral transition: warm-up

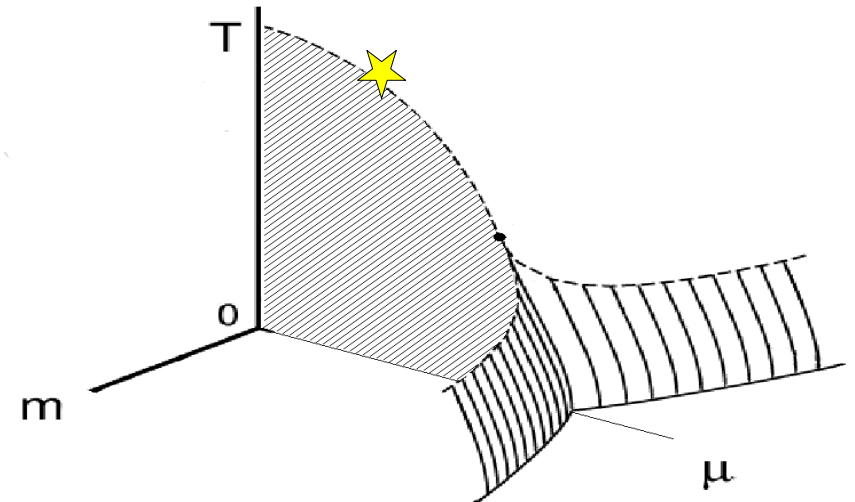
- **THE** order parameter $\langle q^{\text{bar}} q \rangle \sim \langle \sigma \rangle$
- Scalar density gap opens below CP
- chiral susceptibility: $\langle \sigma^2 \rangle - \langle \sigma \rangle^2$

$$- \chi_{mm} = - \frac{\partial^2 P}{\partial m^2} \sim \int d^3 r \frac{e^{-M_\sigma r}}{r} \sim \frac{1}{M_\sigma^2} \rightarrow \infty$$

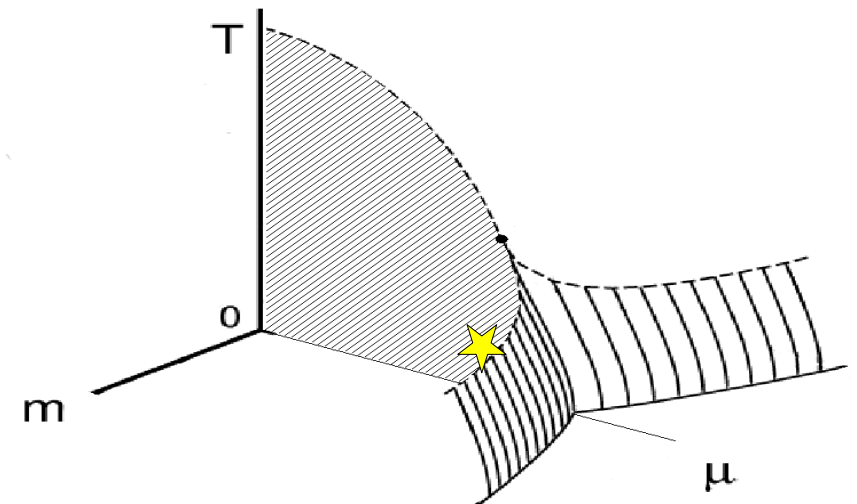
- long-range fluctuation @CP



- $\chi_{\mu\mu}$, **C** finite (in MF)
 - Since no gap for n_B & s



CP with nonzero m



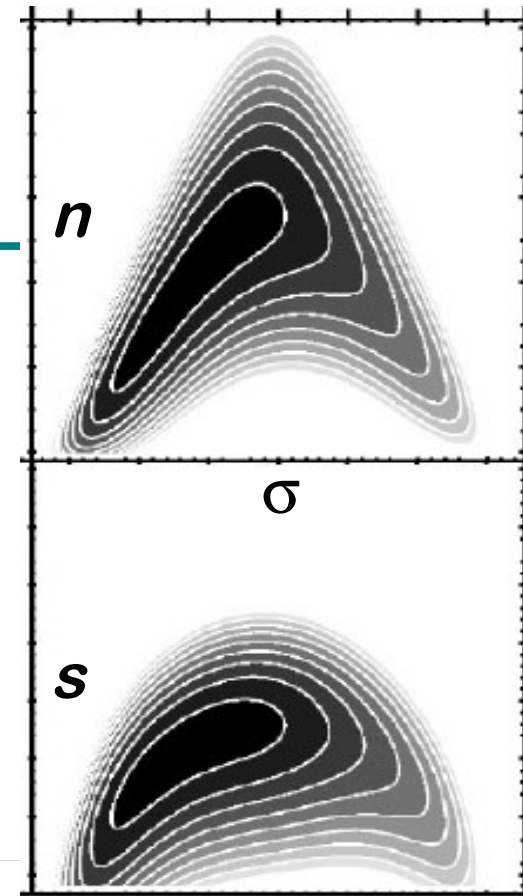
CP with nonzero m

- Across co-existence surface all densities have gaps (Clapeyron-Clausius)

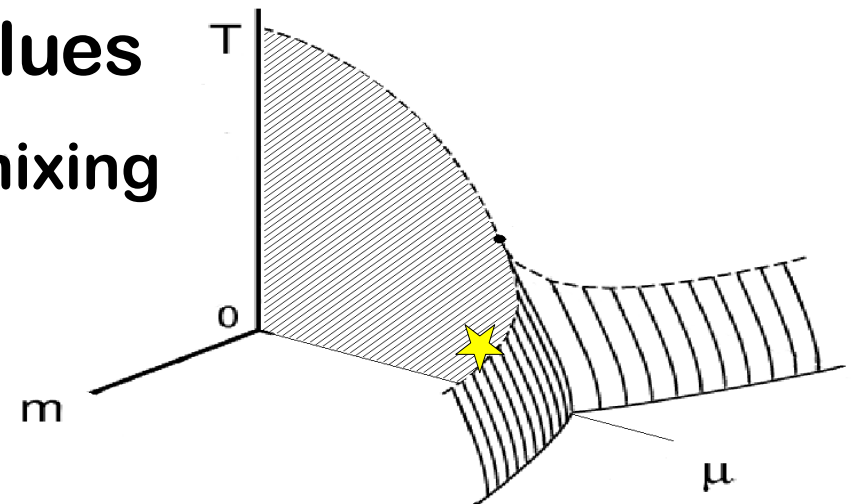
$$\frac{dT}{d\mu} = -\frac{\Delta n}{\Delta s}, \quad \frac{dT}{dm} = -\frac{\Delta \sigma}{\Delta s}$$

- Flat potential wrt σ , n_B , and/or s @CP
- order parameter = deviations of σ , n_B , and/or s from equilibrium values
 - Nonzero m generates linear mixing

- χ_{mm} , $\chi_{\mu\mu}$ and C all diverge

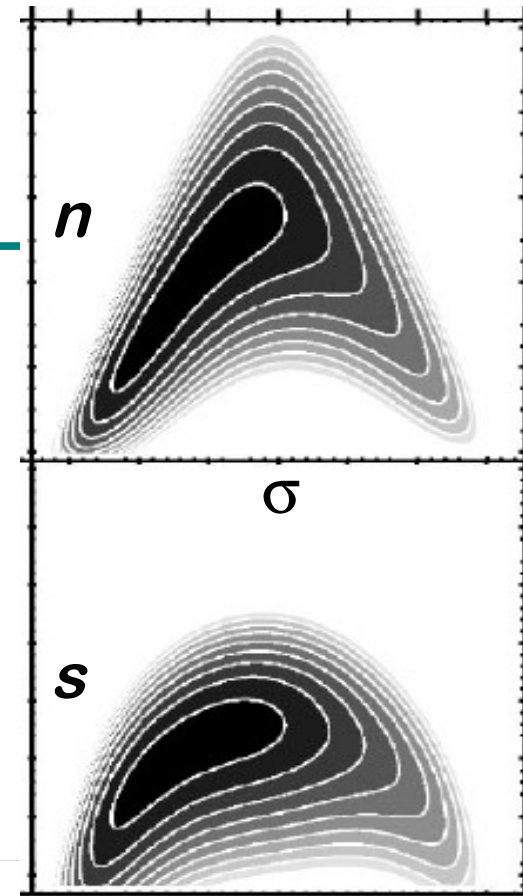


Contour plot of free energy

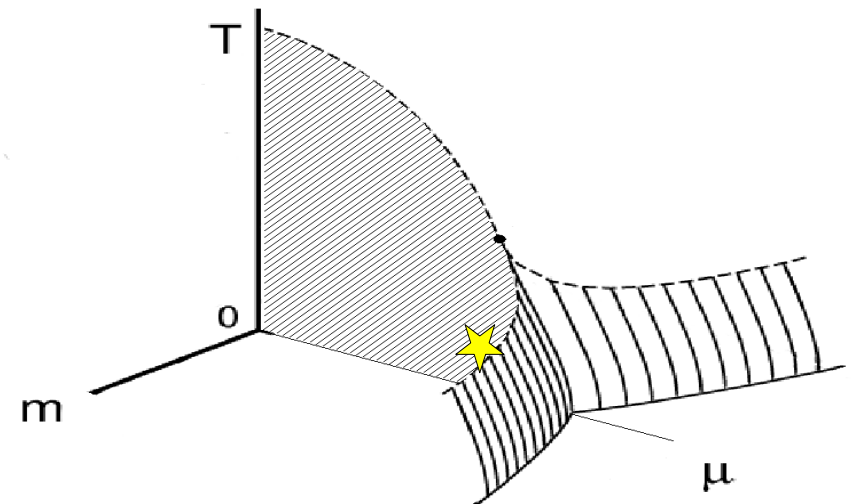


CP with nonzero m

- χ_{mm} , $\chi_{\mu\mu}$ and C all diverge
 - $\chi \sim 1/M_\sigma^2$
 - massless σ ? ω ? or what?



Contour plot of free energy



χ 's & spectral change

χ 's & spectral change

- $q^{\text{bar}}-q$ attraction induces the CP
- Susceptibilities, which diverge at CP, are expressed in terms of mode spectrum

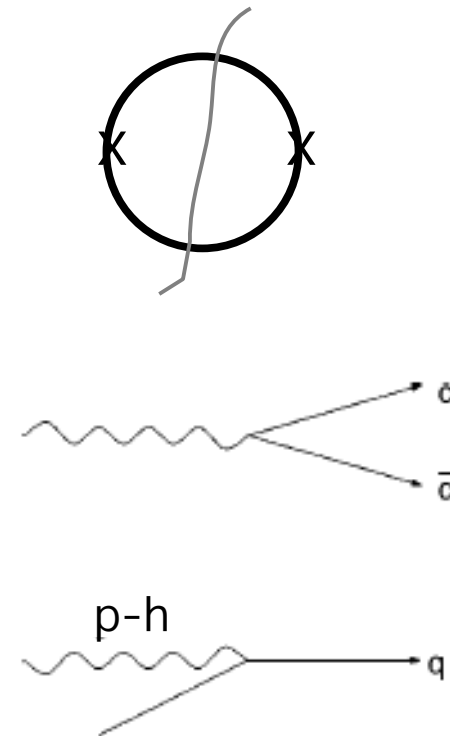
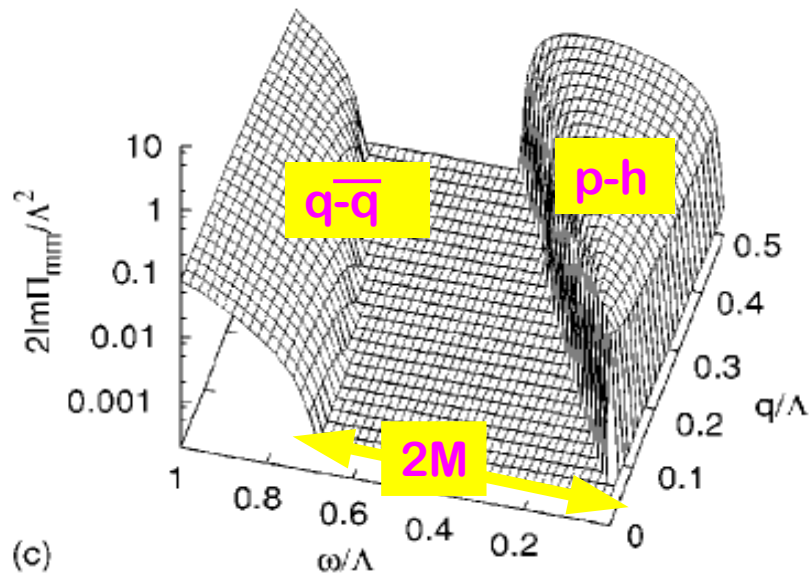
$$\chi_{ab}(q) = \int \frac{d\omega}{2\pi} \frac{\rho_{ab}(\omega, q)}{\omega}$$

- What mode softens in $\rho_{ab}(\omega, q)$ near CP ?
- RPA calc in NJL:

$$\chi_{ab}(iq_4, \mathbf{q}) = \Pi_{ab}(iq_4, \mathbf{q}) + \Pi_{am}(iq_4, \mathbf{q}) \times \frac{1}{1 - 2g\Pi_{mm}(iq_4, \mathbf{q})} 2g\Pi_{mb}(iq_4, \mathbf{q})$$

Free spectral fn. $Im\Pi$

- Spectral fn of massive quark gas
 - spectrum in time- and space-like regions
 - space-like (p-h) mode is characteristic in a medium



$O(4)$ case: dropping mass

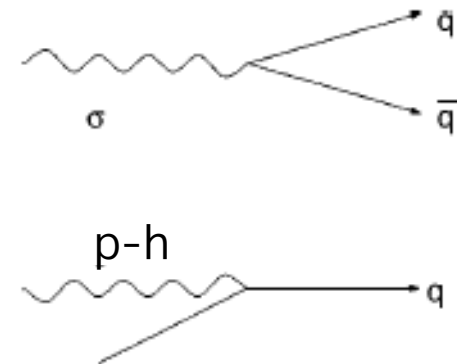
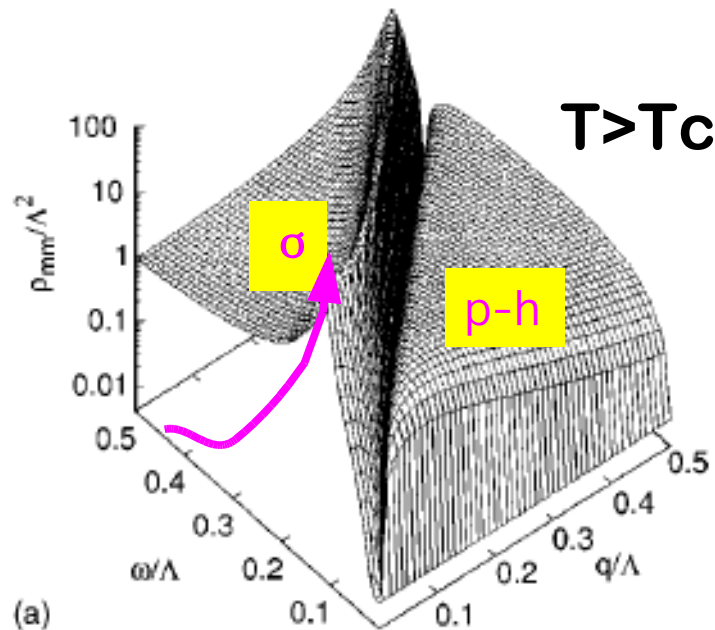
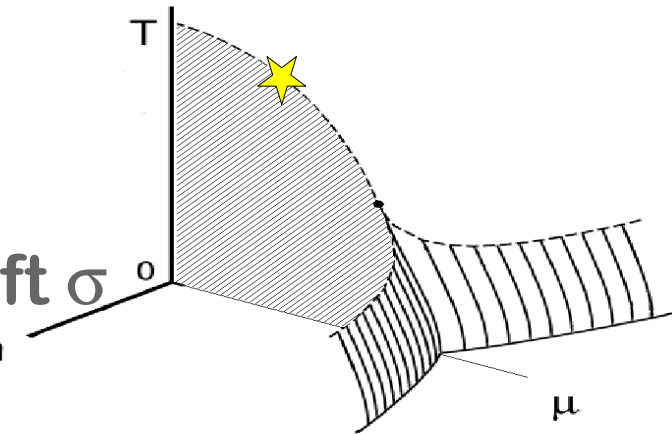
- **Common sense** near $O(4)$ CP

- dynamic σ -mode softens near CP

- thermally-massive quarks give soft σ

Hidaka-Kitazawa m

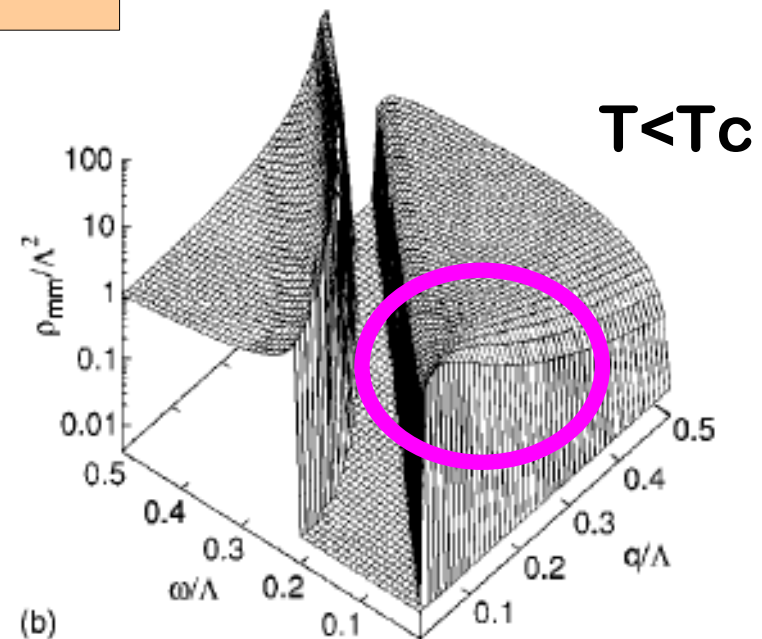
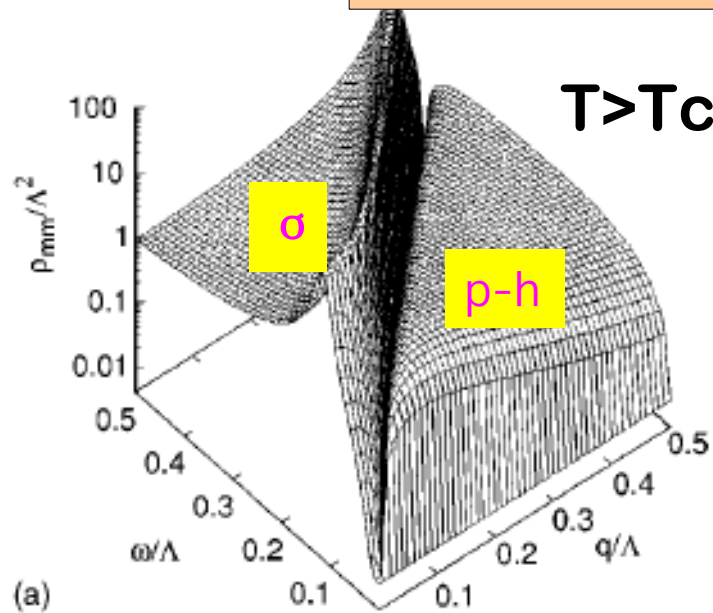
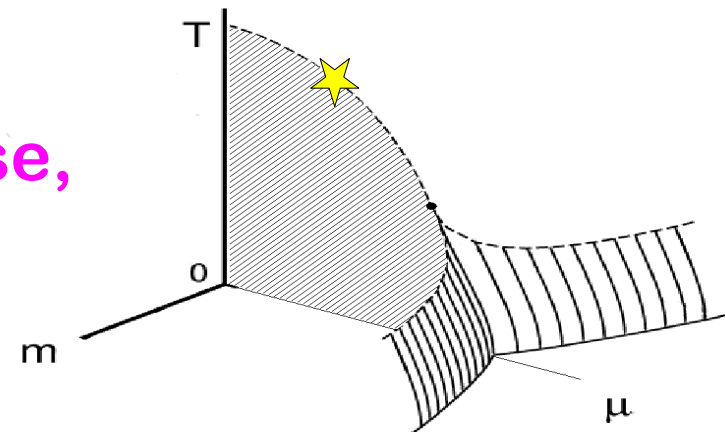
- no p-h mode contributes to χ_{mm}



$O(4)$ case: dropping mass & ...

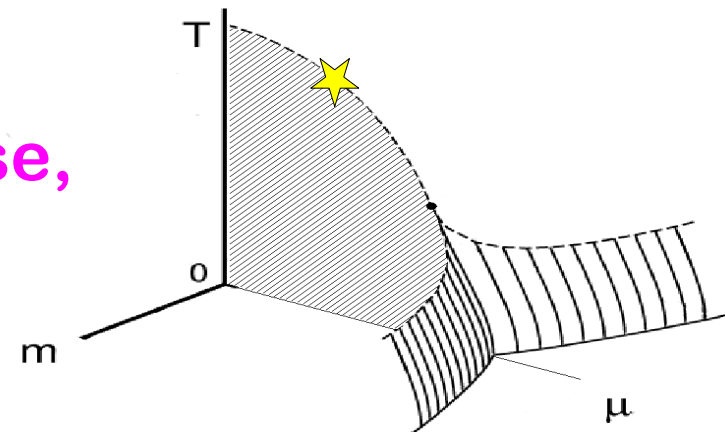
- A little surprise near $O(4)$ CP
 - if approached from broken phase, non-zero M mixes p-h in χ_{mm}

$$\chi_{ab}(q) = \int \frac{d\omega}{2\pi} \frac{\rho_{ab}(\omega, q)}{\omega}$$

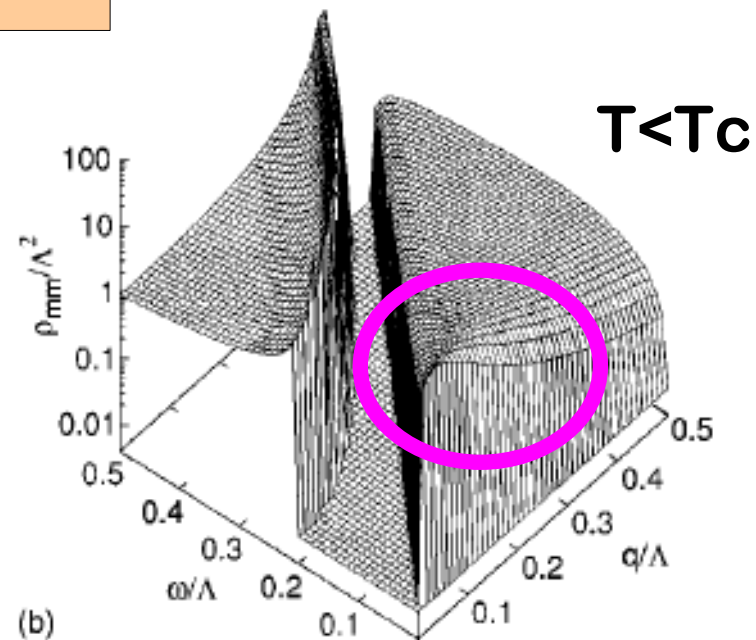
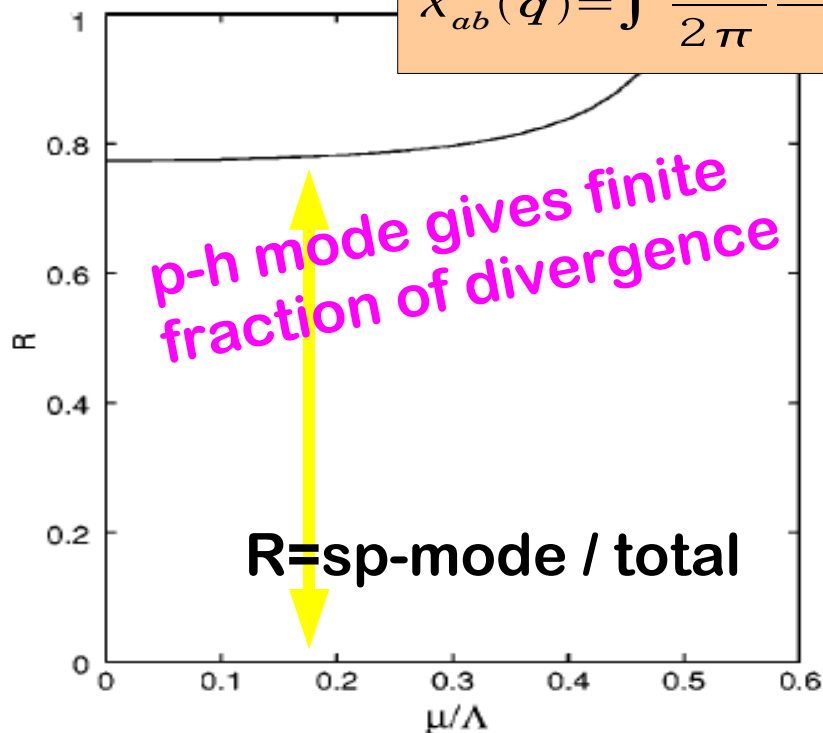


$O(4)$ case: dropping mass & ...

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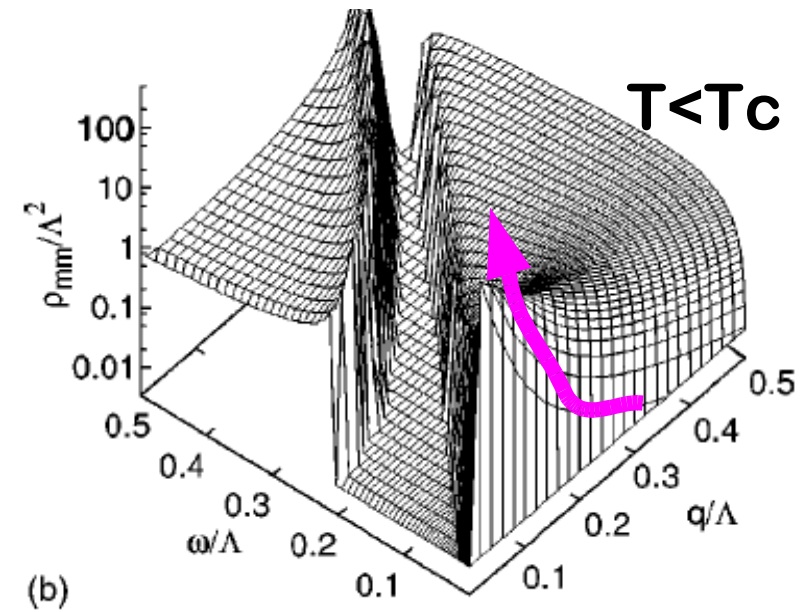
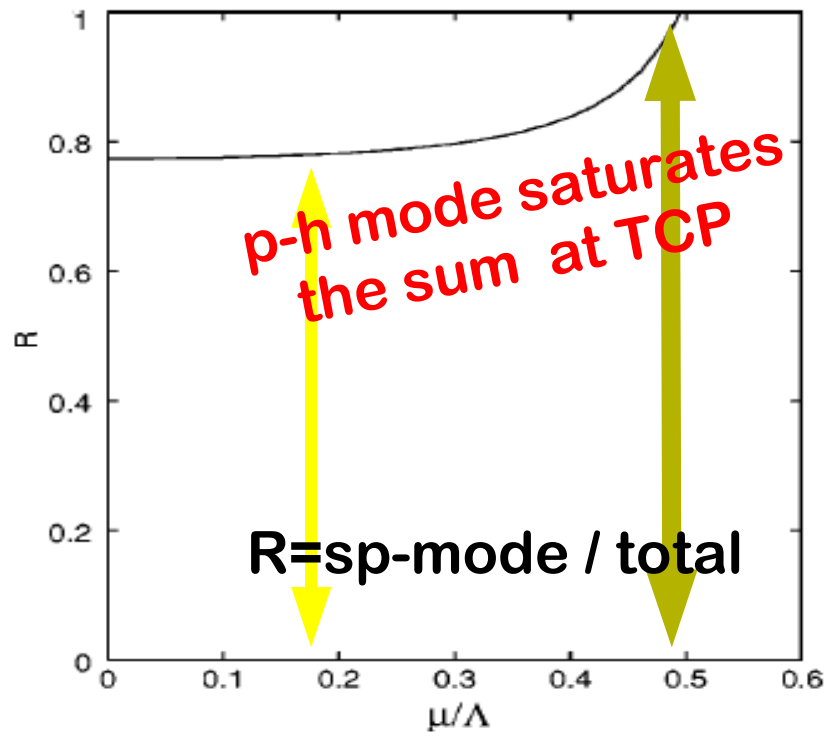
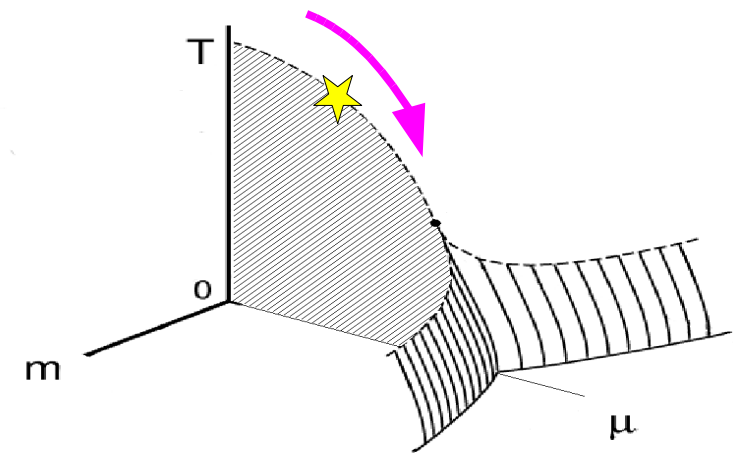


Toward TCP: what happens?

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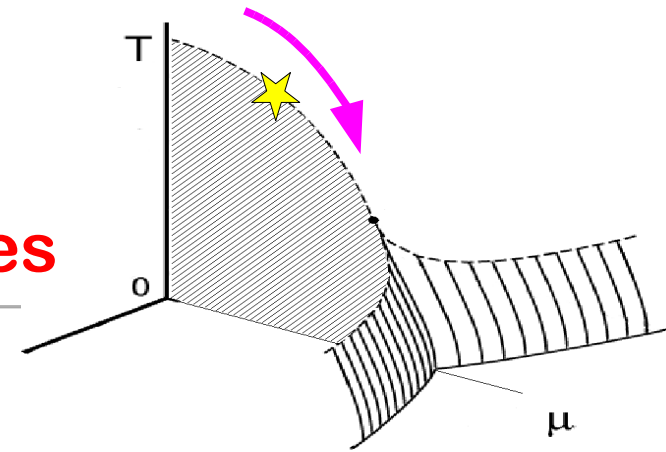
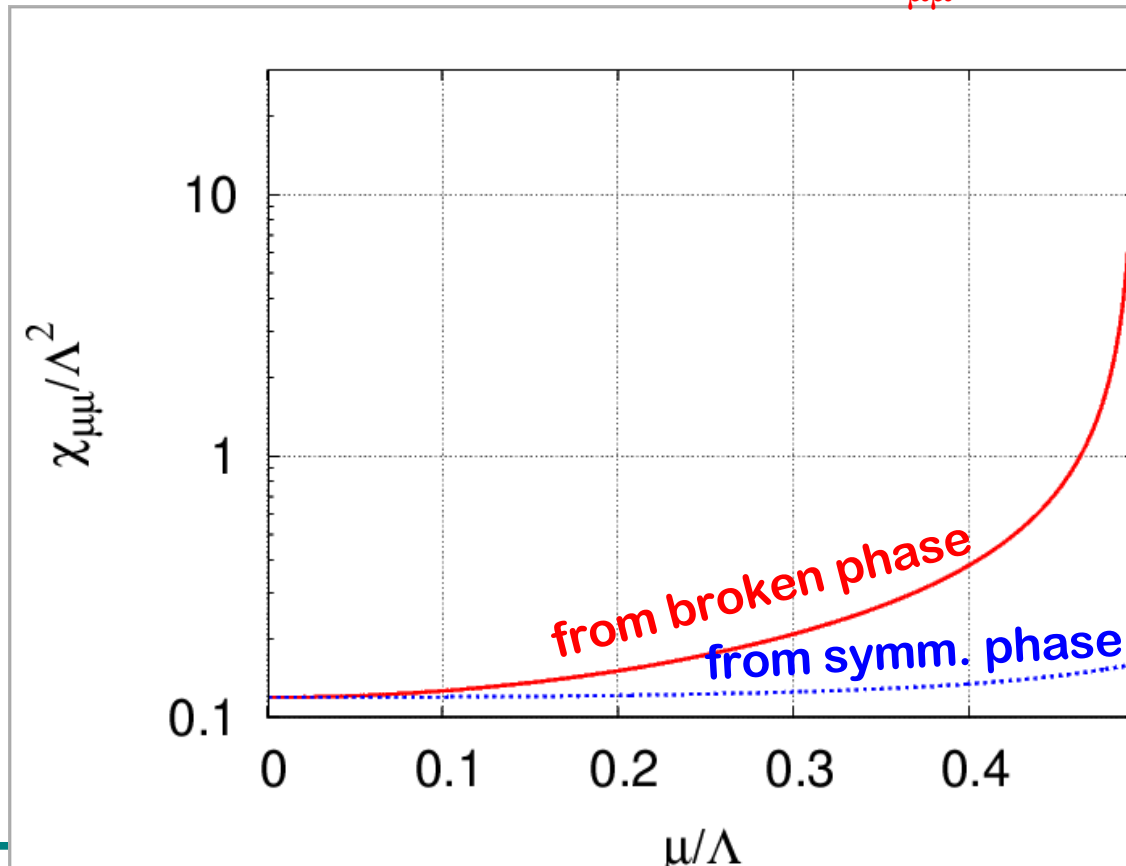
- In addition to soft sigma,
 - **p-h mode enhances and saturates the divergence !**

$$\chi_{ab}(q) = \int \frac{d\omega}{2\pi} \frac{\rho_{ab}(\omega, q)}{\omega}$$



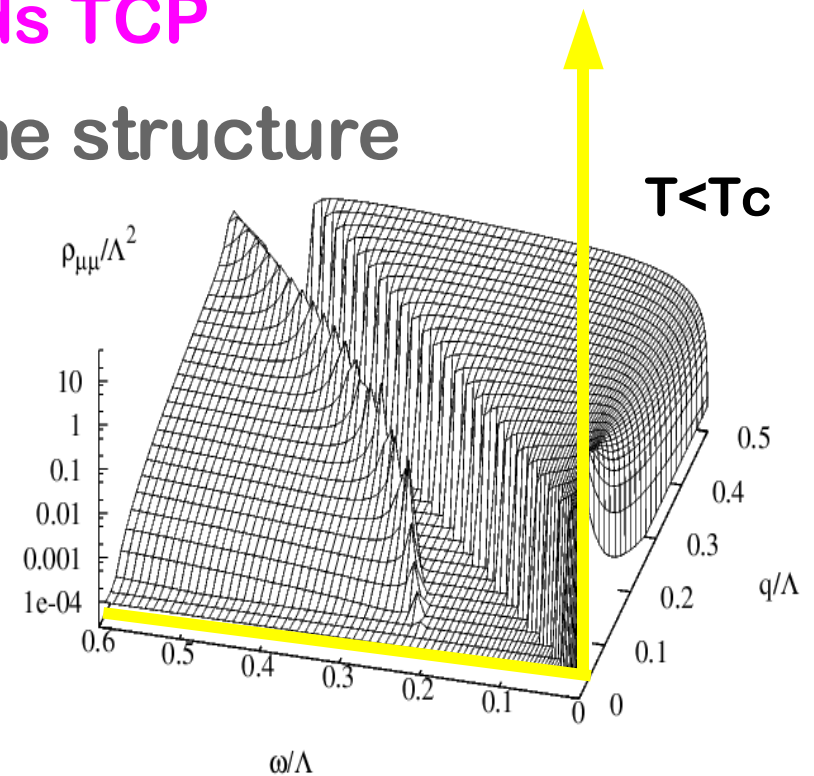
Toward TCP: what happens?

- In addition to soft sigma,
 - p-h mode enhances !
 - quark number suscept. $\chi_{\mu\mu}$ diverges

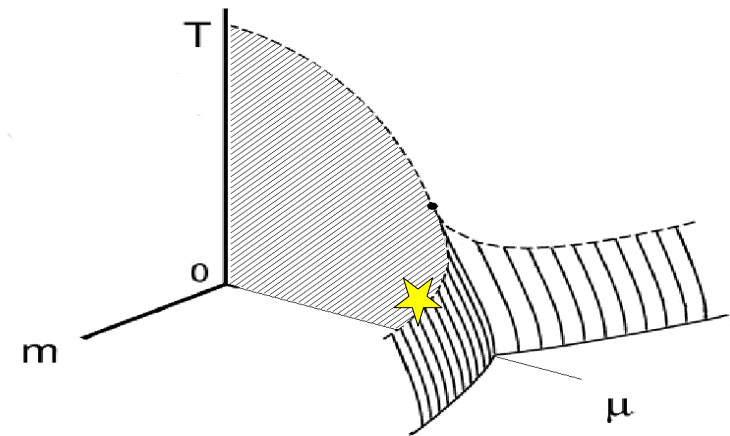


Why is p-h mode enhanced?

- $\chi_{\mu\mu}$ = fluctuation of conserved density
 - at $q=0$ (total charge) it cannot fluctuate
 - **spectral strength only at $\omega=0$:** $\lim_{q \rightarrow 0} \rho_{\mu\mu}(\omega, q) \propto \omega \delta(\omega)$
 - **the strength diverges towards TCP**
 - specific heat C has the same structure

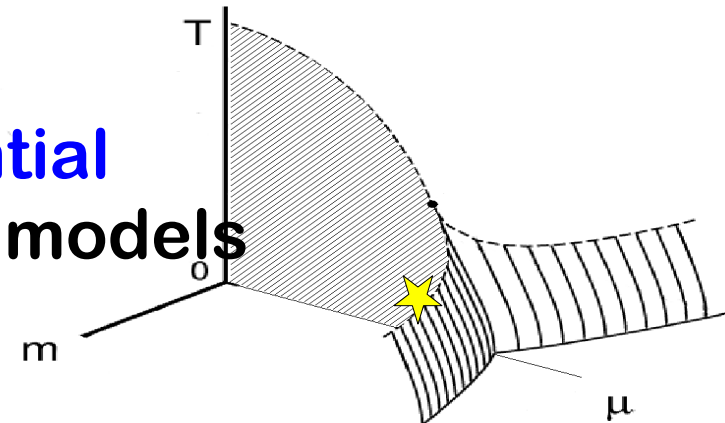


CP with nonzero m



CP with nonzero m

- mass of sigma?
 - Scavenius et al computed **potential curvature** & **pole masses** in two models



PRC64, 045202

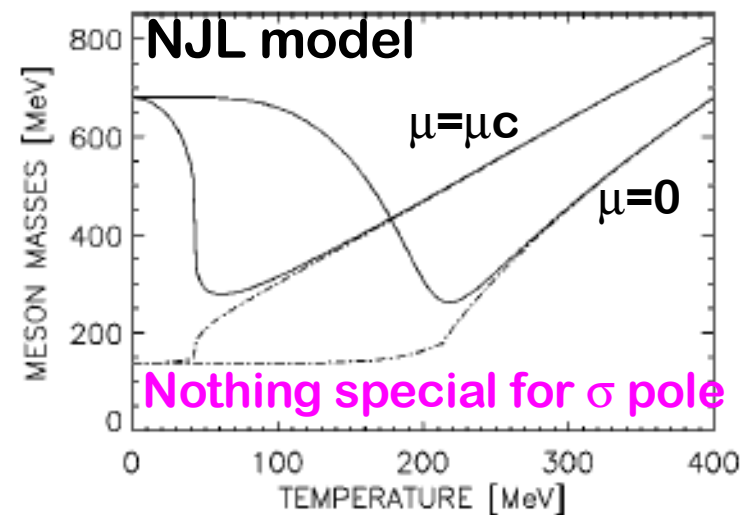
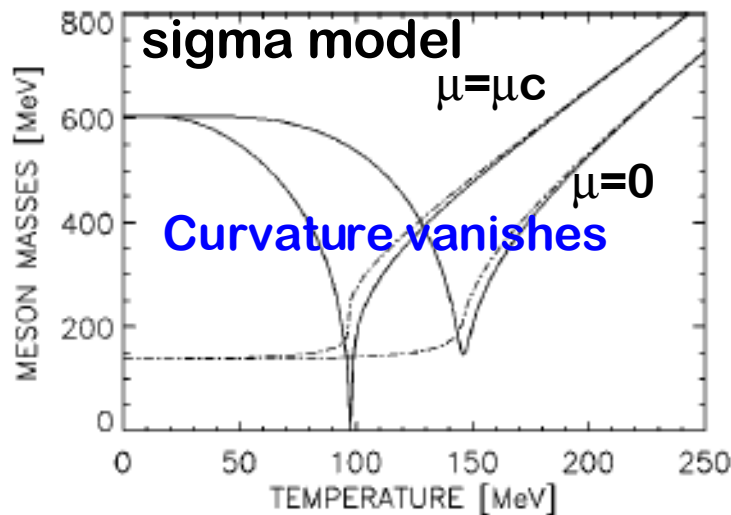
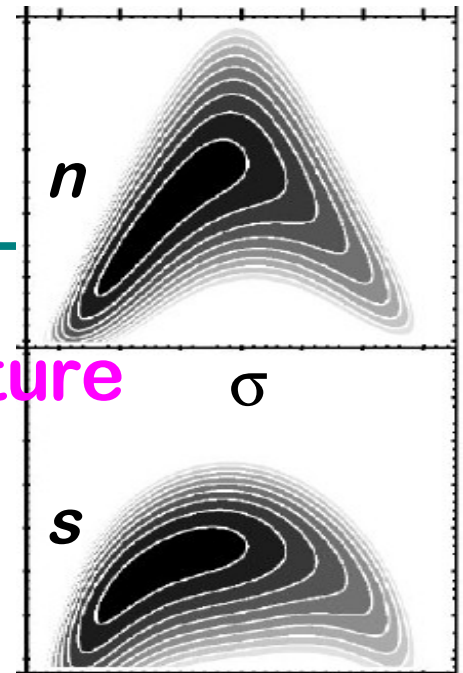


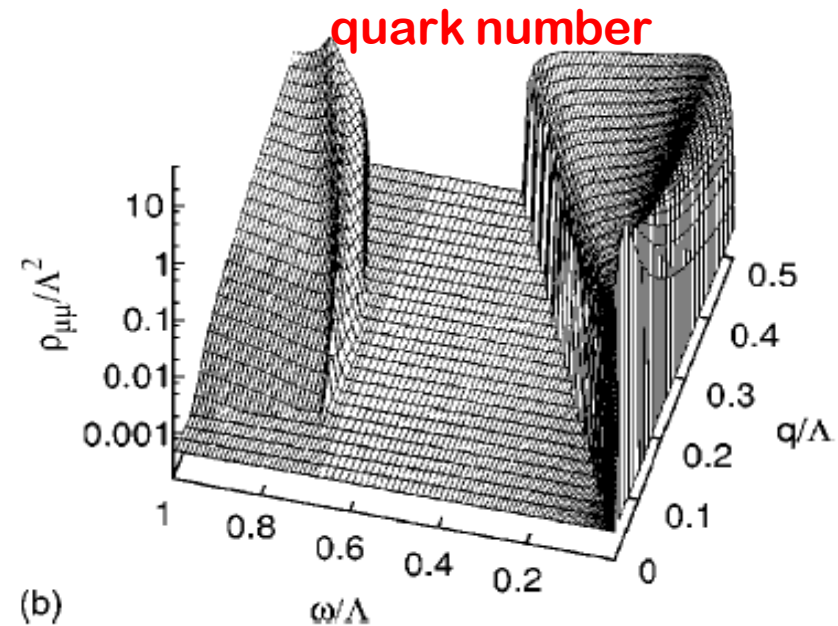
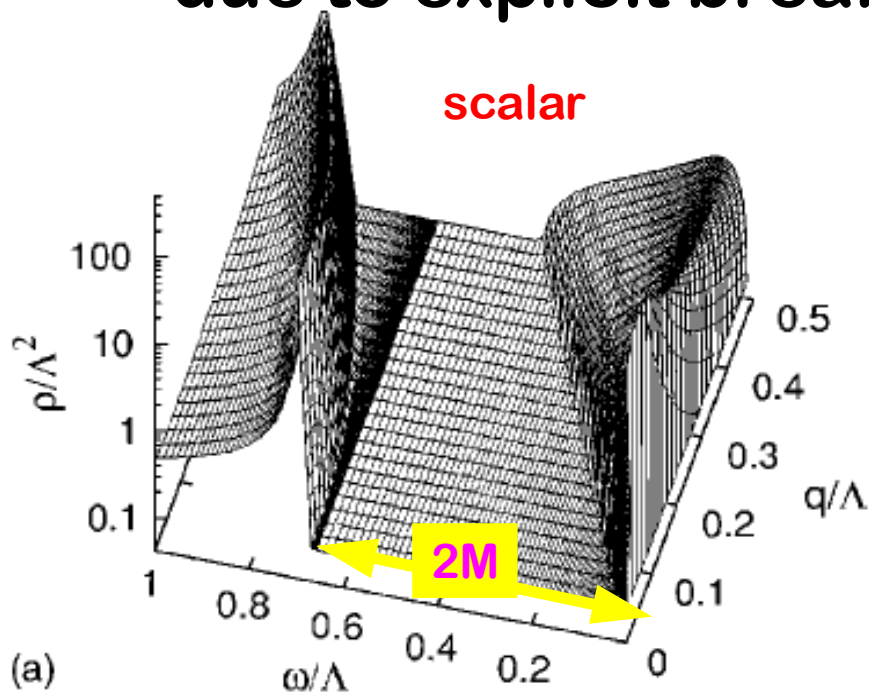
FIG. 6. The sigma mass (solid line) and pion mass (dashed line) in the sigma model (left) and NJL model (right) as functions of temperature for $\mu = 0$ (right pair) and for $\mu = \mu_c$ (left pair).

CP with nonzero m

- p-h mode is responsible for flat curvature & all divergent χ 's
 - no dropping mass at QCD-CP !
 - due to explicit breaking m



Contour plot of free energy



Dynamic Universality

- Static universality of QCD-CP
 - 3D Z(2) Ising
- **Dynamics is constrained by conservation laws**
 - possible slow modes near QCD-CP
 - σ (π massive due to $m_q \neq 0$)
 - $T^{0\mu}$, $n \dots 5$ densities
 - decoupling of σ due to finite m
 - slow modes = sound (2), shear (2), heat (1) = liq.gas
 - relativity vs dynamic universality ?

Hints to phenomenology

- **critical softening in p-h mode**
 - decoupling of σ and ω mesons
 - no direct access to dileptons, $\pi\pi$
 - open question on 'strong NN attraction' near CP
- **within critical region of gas-liquid CP**
 - shear viscosity (small exponent) $\eta \sim \xi^{(1/19)\epsilon}$
 - thermal conductivity $\lambda \sim \xi^{(18/19)\epsilon}$
- **finite size / time effects always in HIC**

Berdnikov-Rajagopal

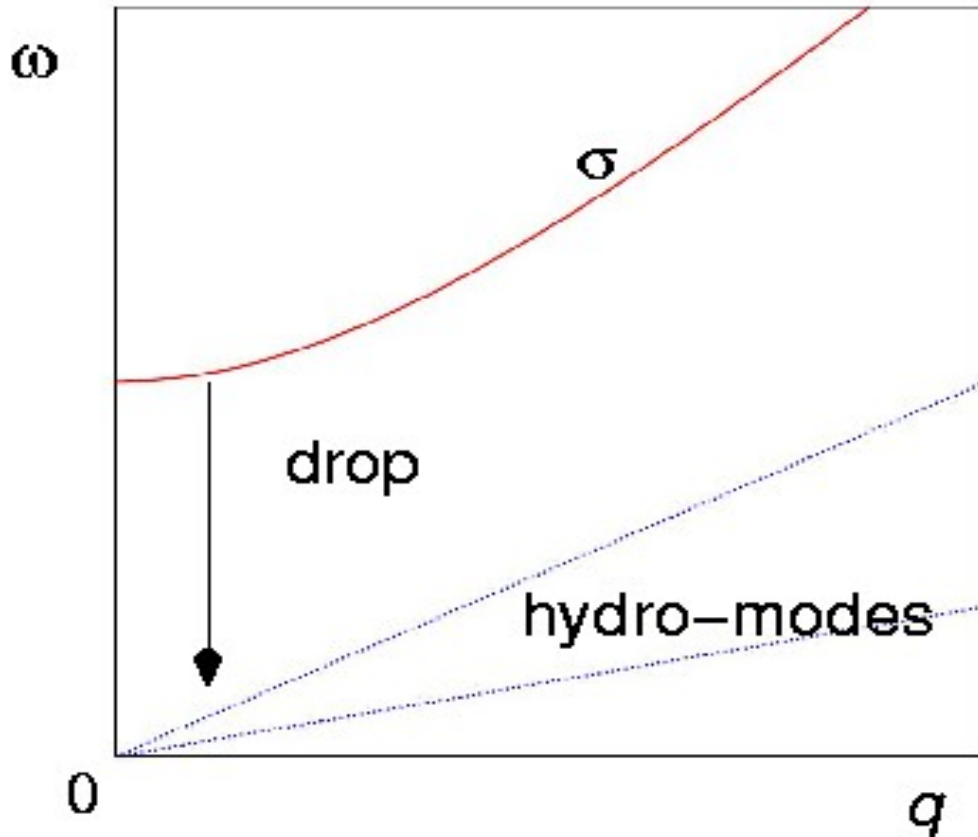
Summary

- **Simple chiral models demonstrate:**
 - σ softens near chiral CP (as everybody knows)
 - p-h also contributes χ_{mm} if approached from broken phase
 - Towards TCP, p-h mode is enhanced to saturate the divergence
- **Critical mode of QCD-CP has p-h type spectrum**
 - no massless σ , no direct dilepton, $\pi\pi$ decay
- **Dynamic Universality is the same as liq.-gas**
 - η & λ should become large in the critical region

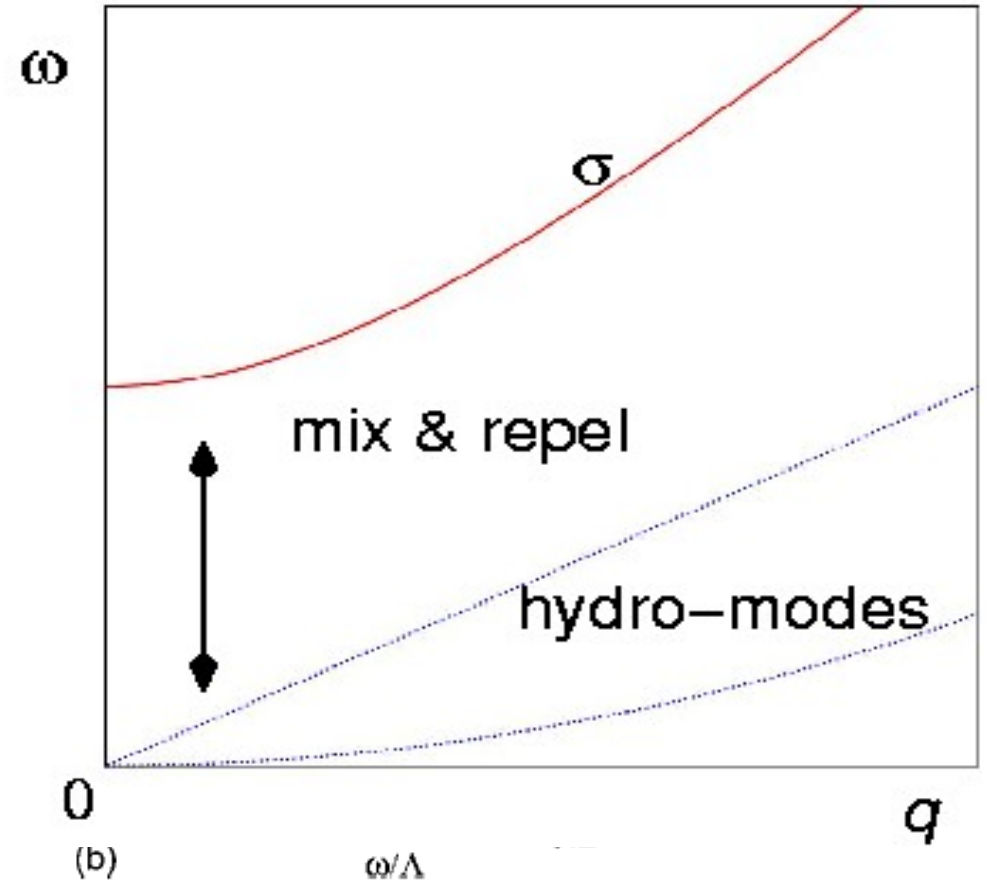
Why is p - h mode enhanced?

- Schematic sketch

w/ exact chiral symmetry

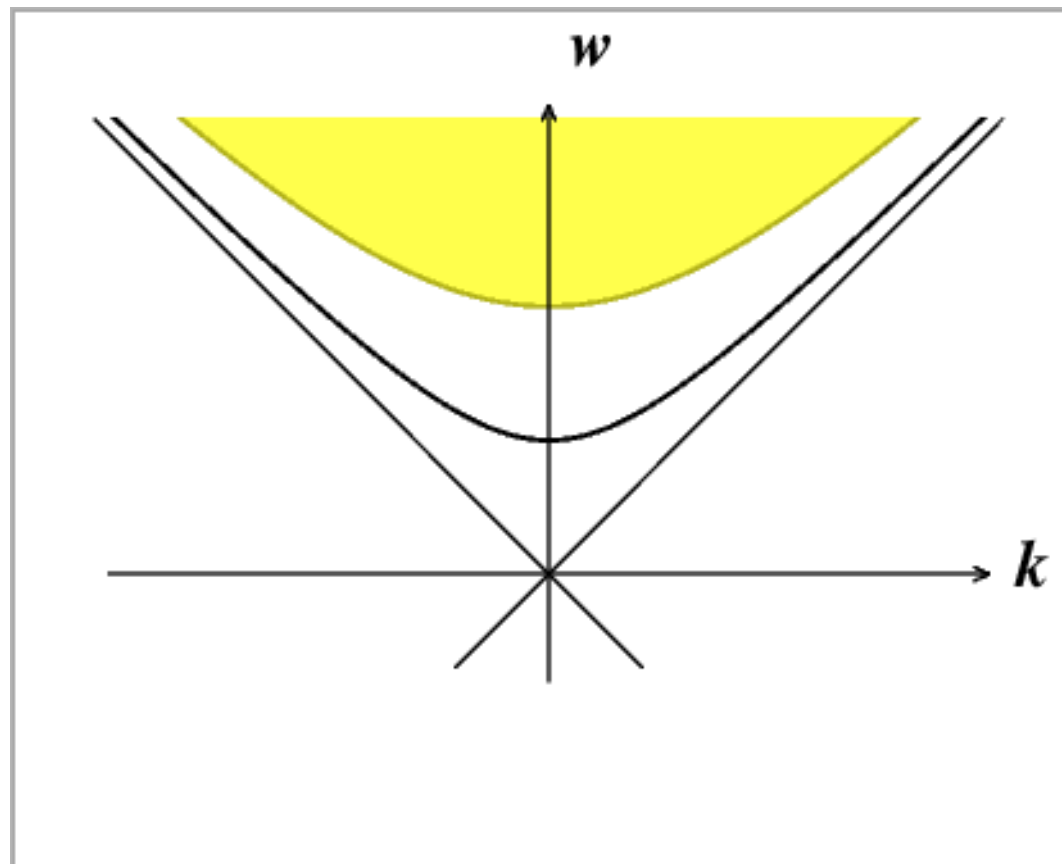


w/ non-zero quark mass



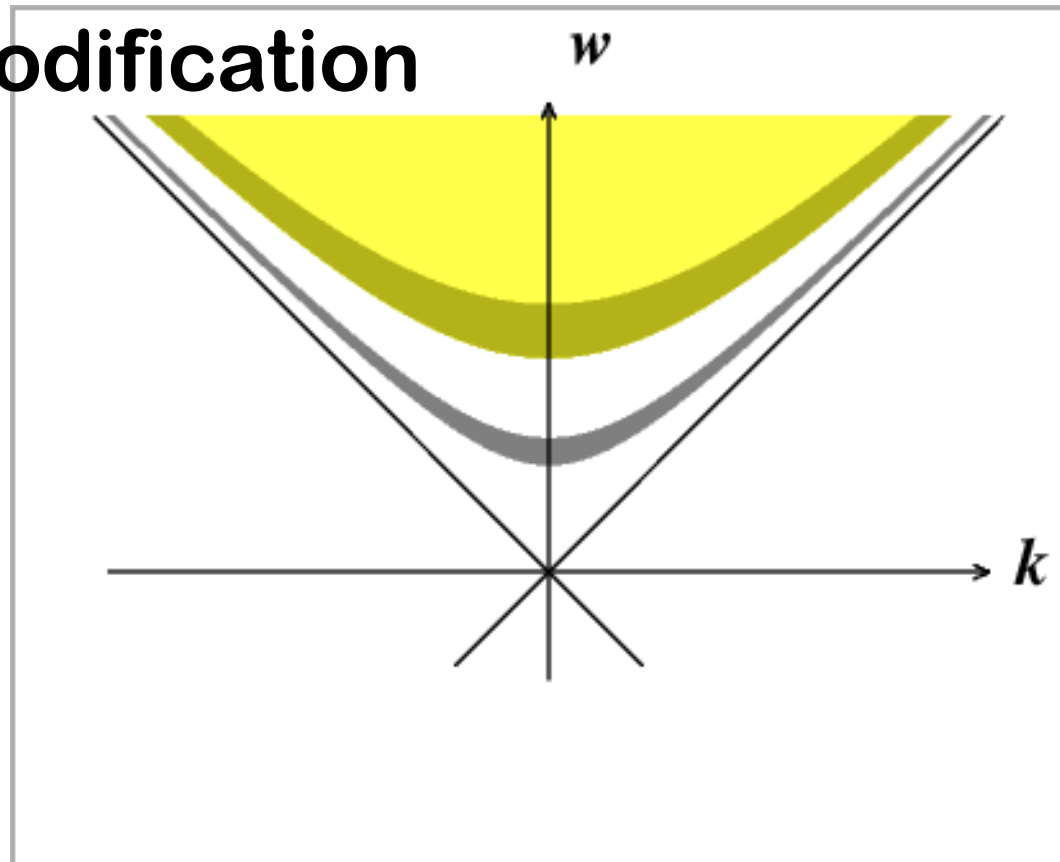
Region of spectrum

- vacuum



Region of spectrum

- mass modification



Region of spectrum

- medium fluctuation

